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LETTER TO THE EDITOR

No correlation of the Lyman continuum escape fraction with spectral hardness

R. Marques-Chaves¹, D. Schaerer^{1,2}, R. O. Amorín^{3,4}, H. Atek⁵, S. Borthakur⁶, J. Chisholm⁷, V. Fernández⁴, S. R. Flury⁸, M. Giavalisco⁸, A. Grazian⁹, M. J. Hayes¹⁰, T. M. Heckman¹¹, A. Henry¹², Y. I. Izotov¹³, A. E. Jaskot¹⁴, Z. Ji⁸, S. R. McCandliss¹¹, M. S. Oey¹⁵, G. Östlin¹⁰, S. Ravindranath¹², M. J. Rutkowski¹⁶, A. Saldana-Lopez¹, H. Teplitz¹⁷, T. X. Thuan¹⁸, A. Verhamme¹, B. Wang¹¹, G. Worseck¹⁹, and X. Xu¹¹

(Affiliations can be found after the references)

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ABSTRACT

The properties that govern the production and escape of hydrogen-ionizing photons (Lyman continuum, LyC; with energies >13.6 eV) in star-forming galaxies are still poorly understood, but they are key to identifying and characterizing the sources that reionized the Universe. Here we empirically explore the relation between the hardness of ionizing radiation and the LyC leakage in a large sample of low- z star-forming galaxies from the recent *Hubble* Space Telescope Low- z Lyman Continuum Survey. Using Sloan Digital Sky Survey stacks and deep X-shooter observations, we investigate the hardness of the ionizing spectra ($Q_{\text{He}^+}/Q_{\text{H}}$) between 54.4 eV (He⁺) and 13.6 eV (H) from the optical recombination lines He II 4686 Å and H β 4861 Å for galaxies with LyC escape fractions spanning a wide range, $f_{\text{esc}}(\text{LyC}) \approx 0\text{--}90\%$. We find that the observed intensity of He II/H β is primarily driven by variations in the metallicity, but is not correlated with LyC leakage. Both very strong ($\langle f_{\text{esc}}(\text{LyC}) \rangle \approx 0.5$) and nonleakers ($\langle f_{\text{esc}}(\text{LyC}) \rangle \approx 0$) present similar observed intensities of He II and H β at comparable metallicity, between ≈ 0.01 and ≈ 0.02 for $12 + \log(\text{O}/\text{H}) > 8.0$ and < 8.0 , respectively. Our results demonstrate that $Q_{\text{He}^+}/Q_{\text{H}}$ does not correlate with $f_{\text{esc}}(\text{LyC})$, which implies that strong LyC emitters do not show harder ionizing spectra than nonleakers at similar metallicity.

Key words. galaxies: starburst – galaxies: high-redshift – dark ages, reionization, first stars

1. Introduction

Numerous studies of star-forming galaxies at $z \sim 1\text{--}3$ have shown differences with low- z galaxies in the Sloan Digital Sky Survey (SDSS), based on shifts in the classical, optical BPT emission-line diagrams toward higher excitation in the nebular gas phase (see the review of Kewley et al. 2019, and references therein). Such shifts can be due to different physical effects causing harder ionizing spectra of their stellar populations (e.g., lower metallicities, younger ages), interstellar medium (ISM) properties (higher ionization parameter, different gas pressure and/or density), α -element enhancements, selection effects, or others (cf. Steidel et al. 2016; Kewley et al. 2019; Izotov et al. 2021a).

Possibly larger differences have been observed in rest-UV spectra of distant galaxies, where strong nebular emission from C IV $\lambda 1550$ and He II $\lambda 1640$ has been found (e.g., Stark et al. 2015; Mainali et al. 2017; Vanzella et al. 2021). These lines have also been detected in confirmed or suspected Lyman-continuum (LyC) emitters at $z \sim 2\text{--}4$ (see Vanzella et al. 2018, 2020; Naidu et al. 2022). These observations suggest hard ionizing spectra in particular because C IV and He II probe energies above 47.8 and 54.5 eV, respectively, and are thus sensitive to higher-energy radiation than the classical strong optical emission lines of [O III], [O II], [N II], and [S II] (e.g., Kewley et al. 2019; Feltre et al. 2016).

Nebular He II emission provides the best measure of the hardness of the ionizing radiation field because its recombination lines¹ are basically direct photons counters for energies >54 eV, whereas the forbidden metal lines depend on many parameters

¹ The strongest and most commonly detected lines are He II $\lambda 1640$ and He II $\lambda 4686$.

(ionization parameter and others). Naidu et al. (2022) stacked rest-UV spectra of Ly α emitters (LAEs) at $z \sim 2$, for which they estimated LyC escape fractions using indirect methods, finding narrow He II $\lambda 1640$, C IV $\lambda 1550$, and other lines in sources with high LyC escape, whereas low escape sources only show [C III] and [O III] emission. Based on this, they suggested a possible relation between the hardness of the ionizing spectra and LyC escape. He II is also seen in the stacks of the LyC candidates of Marchi et al. (2018) selected as Lyman break galaxies (LBGs), although the poor resolution makes it difficult to exclude a significant contribution from stellar emission. Moreover, nebular He II emission has not been reported in the LBG stacked LyC emitter spectra of Steidel et al. (2018). These findings call for a clarification of the possible link between LyC escape, nebular He II emission, and the hardness of the radiation between energies 13.6 and 54 eV.

At low redshift, Jaskot & Oey (2013) have examined how hard ionizing radiation can influence emission line diagnostics of the optical depth of LyC radiation in green pea galaxies, a class of strong emission line galaxies now known to contain LyC emitters. Schaerer et al. (2022) have recently discovered intense C IV and He II $\lambda 1640$ emission lines in three low- z strong LyC emitters, ($f_{\text{esc}} > 0.1$), with UV properties similar to the high- z galaxies mentioned above. They proposed that strong C IV $\lambda 1550$ emission indicates high LyC escape fractions. Furthermore, they also estimated that strong LyC leakers do not have harder ionizing spectra than nonleakers, and that the presence of strong C IV and He II in the spectra of LyC leakers could be primarily due to a high ionizing photon production (Schaerer et al. 2022).

Table 1. Summary of the results of SDSS stacks.

SDSS bins (1)	Denomination (2)	Definition (3)	N (4)	f_{esc} (5)	$12 + \log(\text{O}/\text{H})$ (6)	$\text{He II}/\text{H}\beta$ (7)	$Q_{\text{He}^+}/Q_{\text{H}}$ (8)
f_{esc}	Very strong leakers	$f_{\text{esc}} \geq 0.2$	8	$0.497^{+0.242}_{-0.140}$	$7.849^{+0.265}_{-0.165}$	0.016 ± 0.007	0.005 ± 0.002
	Strong leakers	$0.05 \leq f_{\text{esc}} < 0.2$	13	$0.117^{+0.053}_{-0.042}$	$8.107^{+0.296}_{-0.221}$	0.013 ± 0.006	0.007 ± 0.003
	Weak leakers	$f_{\text{esc}} < 0.05$	29	$0.025^{+0.014}_{-0.011}$	$8.115^{+0.221}_{-0.174}$	0.012 ± 0.004	0.007 ± 0.002
	Non-leakers	$S/\text{N}(\text{LyC}) < 2$	39	<0.011	$8.175^{+0.245}_{-0.186}$	0.015 ± 0.005	0.009 ± 0.003
f_{esc} and metallicity	$f_{\text{esc}} \geq 0.05$ and $12 + \log(\text{O}/\text{H}) < 8.0$		11	$0.346^{+0.322}_{-0.228}$	$7.811^{+0.161}_{-0.109}$	0.020 ± 0.005	0.008 ± 0.002
	$f_{\text{esc}} \geq 0.05$ and $12 + \log(\text{O}/\text{H}) = 8.0\text{--}8.3$		6	$0.203^{+0.159}_{-0.113}$	$8.095^{+0.110}_{-0.086}$	0.013 ± 0.006	0.006 ± 0.003
	$f_{\text{esc}} \geq 0.05$ and $12 + \log(\text{O}/\text{H}) > 8.3$		4	$0.120^{+0.066}_{-0.056}$	$8.423^{+0.102}_{-0.051}$	<0.011	<0.006
	$f_{\text{esc}} < 0.05$ and $12 + \log(\text{O}/\text{H}) < 8.0$		9	$0.030^{+0.013}_{-0.010}$	$7.895^{+0.060}_{-0.042}$	<0.016	<0.009
	$f_{\text{esc}} < 0.05$ and $12 + \log(\text{O}/\text{H}) = 8.0\text{--}8.3$		15	$0.022^{+0.015}_{-0.011}$	$8.153^{+0.100}_{-0.084}$	0.015 ± 0.004	0.008 ± 0.002
	$f_{\text{esc}} < 0.05$ and $12 + \log(\text{O}/\text{H}) > 8.3$		5	$0.023^{+0.015}_{-0.009}$	$8.395^{+0.122}_{-0.029}$	<0.023	<0.013
	$S/\text{N}(\text{LyC}) < 2$ and $12 + \log(\text{O}/\text{H}) < 8.0$		6	<0.027	$7.786^{+0.173}_{-0.120}$	0.030 ± 0.009	0.017 ± 0.005
	$S/\text{N}(\text{LyC}) < 2$ and $12 + \log(\text{O}/\text{H}) = 8.0\text{--}8.3$		22	<0.008	$8.175^{+0.091}_{-0.075}$	0.011 ± 0.004	0.006 ± 0.002
	$S/\text{N}(\text{LyC}) < 2$ and $12 + \log(\text{O}/\text{H}) > 8.3$		11	<0.009	$8.388^{+0.086}_{-0.057}$	<0.013	<0.007

Notes. Columns (1), (2), and (3): Groups, denomination, and definition of each bin. Column (4): Number of sources in each bin. Columns (5) and (6): mean and 68% confidence intervals (bootstrap resampling) of f_{esc} and $12 + \log(\text{O}/\text{H})$ for each bin. Column (7): Observed $\text{He II}/\text{H}\beta$ intensities. Column (8): Hardness of the ionizing spectra between 54.4 eV (He^+) and 13.6 eV (H) following Eq. (1) and assuming $f_{\text{esc}}^{\text{He}^+} \approx 0$. Upper limits refer to a 2σ limit.

On the other hand, Pérez-Montero et al. (2020) suggested that the observed $\text{He II } \lambda 4686$ emission in metal-poor low- z galaxies could be explained by significant photon leakage, without the need for additional sources of hard ionizing photons, while it is generally accepted that normal stellar populations cannot produce sufficient photons above 54 eV (see, e.g., Stasińska et al. 2015; Schaerer et al. 2019; Olivier et al. 2021; Simmonds et al. 2021). These results again call for a systematic approach to determine whether LyC emitting galaxies show a different He^+ -ionizing radiation than comparable galaxies with negligible Lyman continuum escape.

In this Letter we determine empirically and in a differential manner the hardness of the ionizing spectra between 54.4 and 13.6 eV (i.e., the ionization potential of He^+ and that of H) as a function of the LyC escape fraction for the first time. To achieve this goal, we use the data from 89 low-redshift star-forming galaxies for which Lyman continuum and nonionizing UV spectra are available from the recent *Hubble Space Telescope* (HST) Low- z Lyman Continuum Survey and other observations (see Sect. 2). Using optical spectra and stacking, we determine the hardness of ionizing radiation field from the optical recombination lines $\text{He II } \lambda 4686$ and $\text{H}\beta$ in a straightforward manner, in contrast to methods using forbidden optical lines and to methods relying on UV emission lines. Our first constraints on the hardness of LyC emitters and a control sample thus contribute to a better knowledge of the analogue sources of cosmic reionization, their ionizing radiation field and ISM.

2. Observations

2.1. Lyman continuum and UV observations

We used the recent LzLCS, which consists of 66 star-forming galaxies at $z \approx 0.2\text{--}0.4$ (Flury et al. 2022a,b; Saldana-Lopez et al. 2022) and 23 additional sources previously studied by Izotov et al. (2016b,c, 2018a, 2021b) and Wang et al. (2019). This unique dataset includes 37 LyC emitters detected with high signal-to-noise ratio ($S/\text{N} > 3$), with escape fractions spanning a wide range ($f_{\text{esc}} \sim 1\text{--}90\%$), and strong upper limits for the remaining sources. The galaxies span a wide range of

physical properties with stellar masses $\log(M_*/M_\odot) = 7.2\text{--}10.8$, star formation rates $\text{SFR} = (3\text{--}80) M_\odot \text{yr}^{-1}$, and nebular abundances $12 + \log(\text{O}/\text{H}) = 7.5\text{--}8.6$, derived using the direct method (Flury et al. 2022a).

2.2. SDSS optical spectra

Optical spectra are available for all sources from SDSS and BOSS surveys (Eisenstein et al. 2011). We retrieved optical spectra from the SDSS Science Archive Server. SDSS spectra have a spectral resolution $R \sim 2000$. The extinction curve of Fitzpatrick (1999) and the Galactic $E(B - V)$ values from the dust maps of Green et al. (2018) were used to correct the reddening effect of the Galaxy. Flux measurements were derived using Gaussian profiles and the PYTHON nonlinear least-squares function CURVE_FIT and corresponding uncertainties using a Monte Carlo method. The Cardelli et al. (1989) reddening law ($R_V = 3.1$) was adopted to correct for the internal extinction using the ratios of well-detected Balmer emission lines (following Izotov et al. 1994).

In the lines of interest in this work, $\text{H}\beta$ is detected with high significance in all galaxies, but not He II , yielding 3σ limits of $\text{He II}/\text{H}\beta \lesssim 0.04$ on average, which are fairly above the typical intensities $\text{He II}/\text{H}\beta \lesssim 0.02$ found in other compact star-forming galaxies (e.g., Izotov et al. 2016a), including LyC emitters (Guseva et al. 2020). Therefore, we performed a stacking analysis of the SDSS spectra to improve the $\text{He II}/\text{H}\beta$ limits.

We built two different groups of stacked spectra, one composed of bins in absolute LyC escape fraction, $f_{\text{esc}}(\text{LyC})$, and another one composed of bins in $f_{\text{esc}}(\text{LyC})$ and $12 + \log(\text{O}/\text{H})$. For the first group, we defined the four bins using f_{esc} determined from the UV analysis of Saldana-Lopez et al. (2022). For the second group, we built nine stack spectra considering three bins of $f_{\text{esc}}(\text{LyC})$ and three bins of metallicity. Table 1 contains the definition of the bins and the main derived properties from the stacked spectra.

For each bin, the SDSS spectra were deredshifted using the systemic redshifts from the observed wavelengths of bright optical lines and were resampled using a linear interpolation onto a common wavelength grid. Next, we normalized the spectra at

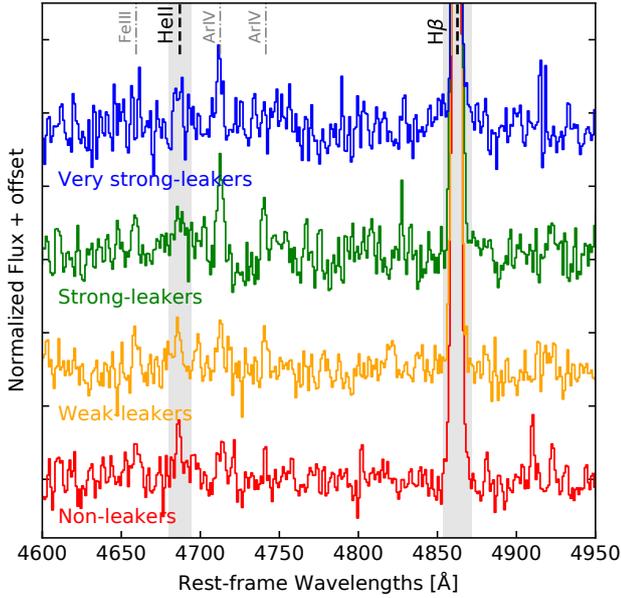


Fig. 1. Normalized stacked SDSS spectra in bins of f_{esc} : Nonleakers (red), weak leakers (orange), strong leakers (green), and very strong leakers (blue). Regions around He II 4686 Å and H β 4861 Å are highlighted in gray. He II is detected in the four stacks with $S/N \sim 2-3$. Other faint lines are also detected ([Fe III] 4658 Å, [Ar IV]+He I 4712 Å, and [Ar IV] 4740 Å), and their positions are marked with dashed gray lines.

$\lambda = 4750-4850$ Å, that is, relatively free of emission or absorption features. Finally, we stacked all spectra by averaging the flux in each spectral bin. We also tested other stacking methods using the median and the weight-average using the uncertainty spectra, but no significant differences in the observed intensity of He II/H β were found between these three methods. Figure 1 shows the four stacked spectra in bins of $f_{\text{esc}}(\text{LyC})$. He II is detected at least in three stacks, together with other faint lines, such as [Fe III] 4658 Å and the [Ar IV] doublet at 4711 Å and 4740 Å. We note the lack of significant broad emission in the stacks around 4650 Å and 5808 Å, which might be associated with a significant contribution of Wolf-Rayet stars (e.g., Brinchmann et al. 2008).

2.3. X-shooter/VLT observations

Spectroscopic observations of eight LzLCS sources were carried out during 2021 with the X-shooter instrument on the Very Large Telescope (VLT), with total integration times of 50–100 min. The X-shooter spectra were reduced in a standard manner using the ESO Reflex reduction pipeline (version 2.11.5; Freudling et al. 2013) to produce flux-calibrated spectra. The Galaxy and internal reddening were corrected using the same method as described in Sect. 2.2. We also used the measurements for five other LyC emitters observed with X-shooter by Guseva et al. (2020). In total, He II is detected in nine sources at $>2\sigma$, with intrinsic He II/H $\beta \approx 0.007-0.020$. For the remaining four sources, 2σ upper limits between ≈ 0.004 and ≈ 0.010 are inferred.

3. Results and discussion

3.1. He II/H β vs. escape of LyC photons

The main result of our study is presented in the left panel of Fig. 2, which shows the relation between the observed He II/H β

ratio and the observed LyC escape fraction f_{esc} . Overall, we find no significant dependence of the He II/H β intensity on f_{esc} from the SDSS stacks.

To quantify this, we computed Kendall’s- τ correlation coefficient between these two variables following Akritas & Siebert (1996), which allows the inclusion of upper limits. Using the results from SDSS stacks, we find $\tau(\text{SDSS}) = 0.00^{+0.33}_{-0.50}$ (p -value = 1, uncertainties derived using a Monte Carlo method), which is basically consistent with a null correlation within our uncertainties. On the other hand, using the results from the 13 individual sources with deep X-shooter spectra, we find a positive correlation with $\tau(\text{XS}) = 0.603^{+0.103}_{-0.128}$ (p -value = 4.14×10^{-3}). However, these results are statistically less meaningful than those obtained with SDSS stacks using the average properties of 89 sources that also include the X-shooter sources. In addition, the underlying effect of $12 + \log(\text{O}/\text{H})$ on He II/H β has not been considered so far.

The intensity of He II/H β of star-forming galaxies is known to vary with metallicity in that it increases with decreasing of $12 + \log(\text{O}/\text{H})$ (e.g., Brinchmann et al. 2008; Shirazi & Brinchmann 2012; Schaerer et al. 2019). As our sample spans a wide range of metallicities, $12 + \log(\text{O}/\text{H}) = 7.5-8.6$ (Flury et al. 2022a), this behavior should be considered in our analysis. Therefore, we used nine SDSS stacks grouped in three bins of f_{esc} and three bins of $12 + \log(\text{O}/\text{H})$ (see Table 1 for details).

The right panel of Fig. 2 shows the relation between He II/H β and $12 + \log(\text{O}/\text{H})$ for strong LyC leakers (solid circles) and non- or weak LyC leakers (empty and dashed circles) in the metallicity bins. For comparison, we also show the observed intensities of He II/H β and metallicities of a compilation of almost 900 star-forming galaxies from SDSS DR14 (pluses; Izotov et al. 2016a), for which the [O III] 4363 Å line is detected with an accuracy better than 4σ , allowing thus direct abundance determinations using the T_e method. The overall properties of the parent sample are discussed in Guseva et al. (2019).

Two main conclusions can be drawn from the right panel of Fig. 2. First, strong LyC leakers present roughly the same He II/H β intensities as non- or weak leakers at comparable metallicities. For instance, strong He II/H β intensities are found for both nonleakers and strong LyC leakers in the low-metallicity bin (He II/H $\beta = 0.030 \pm 0.009$ and 0.020 ± 0.005 , respectively), while sources with $12 + \log(\text{O}/\text{H}) = 8.0-8.3$ show He II/H $\beta \sim 0.012$, independently of $f_{\text{esc}}(\text{LyC})$. Moreover, for a specific range of $12 + \log(\text{O}/\text{H})$, the line intensities He II/H β inferred for strong leakers do not differ from those typically observed in other star-forming galaxies (Izotov et al. 2016a). Second, low-metallicity star-forming galaxies present stronger He II/H β than high-metallicity galaxies. This trend has been reported in other works of increasing the He II/H β intensity with decreasing metallicity. The same trend is seen in the individual sources with X-shooter observations, where we find an anticorrelation between He II/H β and metallicity of $\tau(\text{XS}) = -0.487^{+0.103}_{-0.103}$ (p -value = 2.04×10^{-2}). For example, the two X-shooter sources with the highest He II/H β ratios (≈ 0.02) present $12 + \log(\text{O}/\text{H}) < 8.0$, while the sources that are not detected in He II all have $12 + \log(\text{O}/\text{H}) > 8.0$. This behavior explains the positive correlation between He II/H β and f_{esc} seen in the individual sources with deep X-shooter spectra (Fig. 2, left).

3.2. No variation in hardness with LyC escape

The hardness of the ionizing radiation field between energies above 54.4 and 13.6 eV, described by $Q_{\text{He}^+}/Q_{\text{H}}$, is to first order

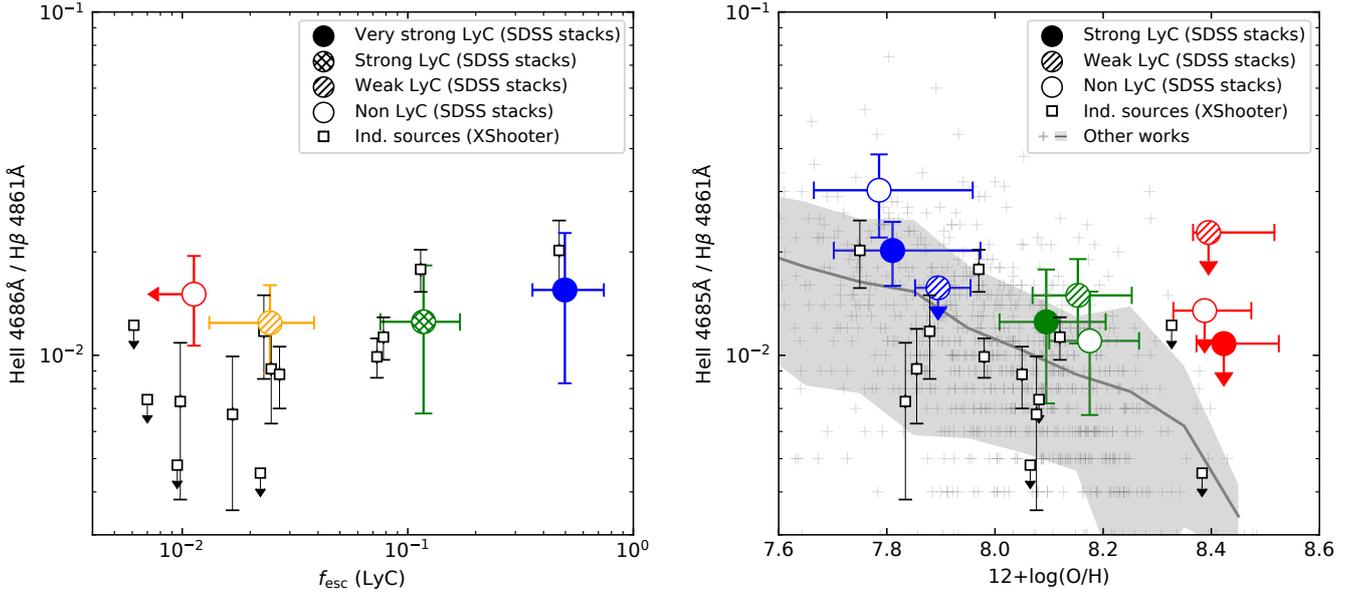


Fig. 2. Intensity of He II/H β as a function of $f_{\text{esc}}(\text{LyC})$ (*left*) and $12 + \log(\text{O}/\text{H})$ (*right*) for different SDSS stacks (circles). The *left panel* shows the four stacks in bins of $f_{\text{esc}}(\text{LyC})$ (very strong, strong, weak, and nonleakers in blue, green, yellow, and red, respectively), and the *right panel* shows the nine stacks in bins of $f_{\text{esc}}(\text{LyC})$ and $12 + \log(\text{O}/\text{H})$ ($12 + \log(\text{O}/\text{H}) < 8.0$, $8.0\text{--}8.3$, and >8.3 , in blue, green, and red, respectively). Details of the definition of stacks can be found in Table 1. x -values and horizontal error bars refer to the mean and 68% confidence intervals (bootstrap resampling) of $f_{\text{esc}}(\text{LyC})$ and $12 + \log(\text{O}/\text{H})$ measured for each bin (Table 1). Individual sources with X-shooter observations are also shown (empty squares; LzLCS sources and five LyC emitters were analyzed in Guseva et al. 2020). Upper limits for SDSS stacks and X-shooter observations refer to 2σ limits. The *right panel* also shows observations (pluses) of a compilation of almost 900 star-forming galaxies from Izotov et al. (2016a) (average and standard deviation marked with a solid line and region, respectively).

related to the relative recombination line intensities by

$$I(4686)/I(\text{H}\beta) = \frac{c_{4686} (1 - f_{\text{esc}}^{\text{He}^+}) \int_{54.4}^{\infty} (F_{\nu}/h\nu) d\nu}{c_{\text{H}\beta} (1 - f_{\text{esc}}) \int_{13.6}^{\infty} (F_{\nu}/h\nu) d\nu} \quad (1)$$

$$= 1.74 \frac{(1 - f_{\text{esc}}^{\text{He}^+}) Q_{\text{He}^+}}{(1 - f_{\text{esc}}) Q_{\text{H}}}, \quad (2)$$

where Q expresses the number of ionizing photons emitted above the corresponding ionization potential, and $c_{ji} = h\nu_{ji} \frac{\alpha_{ji}^{\text{eff}}}{\alpha_B}$ relates the recombination rate to the line intensity (Osterbrock & Ferland 2006). For c_{ji} we have adopted typical values of the electron temperature ($T_e = 10\text{ kK}$). The above expression also accounts for the escape of ionizing photons, which are a priori different for He $^+$ and H-ionizing photons. It is expected that $f_{\text{esc}}^{\text{He}^+} \ll f_{\text{esc}}$ because the doubly ionized He region is generally significantly smaller than the H II region, or in other words, He $^+$ -ionizing photons are absorbed much closer to the source than those of lower energy. Except possibly for extremely hard, power-law-like ionizing spectra, $f_{\text{esc}}^{\text{He}^+}$ is therefore expected to be very low, or at least lower than f_{esc} .

We have shown that the intensity of He II/H β does not depend on f_{esc} and is primarily driven by variations with metallicity. Therefore, Eq. (1) implies that the hardness $Q_{\text{He}^+}/Q_{\text{H}}$ does not correlate with the LyC escape, since, if anything, He II/H β should increase with increasing f_{esc} even for constant $Q_{\text{He}^+}/Q_{\text{H}}$ (see Eq. (1)). Table 1 provides the inferred values of $Q_{\text{He}^+}/Q_{\text{H}}$ for all the stacks using Eq. (1) and assuming $f_{\text{esc}}^{\text{He}^+} \approx 0$. No variation is found in $Q_{\text{He}^+}/Q_{\text{H}}$ between strong and nonleakers within the uncertainties.

3.3. Implications

Our results rule out LyC leakage to explain the origin of nebular He II emission put forward by Pérez-Montero et al. (2020). Their galaxies present roughly the same He II/H β intensities at comparable metallicity as those studied in our work (see their Fig. 1). According to these authors, the observed intensities of He II could be explained by density-bounded H II regions with very highly ionizing photon leaking, with a mean $f_{\text{esc}}(\text{LyC}) \approx 0.74$ for the entire sample. Clearly, the results shown in Fig. 2 contradict this scenario, where the same intensities of He II/H β are found in sources with both very high and very low $f_{\text{esc}}(\text{LyC})$, that is, without any dependence on f_{esc} . Furthermore, the invoked mean $f_{\text{esc}} \approx 0.74$ is significantly higher than measured in comparable low- z galaxies (Flury et al. 2022a,b), and for the strongest known low- z leaker (J1243+4646 with $f_{\text{esc}} \approx 0.73$ reported by Izotov et al. 2018b, or $f_{\text{esc}} \approx 0.89$ inferred in Saldana-Lopez et al. 2022), we measure He II/H $\beta < 0.02$ (2σ), which is not exceptionally high compared to many other star-forming galaxies with similar metallicity ($12 + \log(\text{O}/\text{H}) \approx 7.90$, cf. Fig. 2, right panel). Other mechanisms or sources are needed to explain the origin of nebular He II (see, e.g., Olivier et al. 2021; Simmonds et al. 2021, and references therein).

At $z \sim 2$, Naidu et al. (2022) recently identified two groups of LAEs, one probably showing strong LyC escape, and the other low LyC escape fractions. Comparing the stacked spectra of these two groups, they found differences in their observed rest-frame UV spectra, with the strong leaker candidates showing high-ionization lines of C IV 1550 Å and He II 1640 Å, and the other group showing only lower-ionization lines (e.g., C III 1909 Å). From this finding, Naidu et al. (2022) suggested that strong LyC leakers could have harder ionizing spectra. On the

other hand, if the galaxies studied in our work are comparable to high- z LAEs, our results imply that these observed differences are not related to LyC escape.

Generally speaking, the absence of hardness variations with f_{esc} shows that this property of the global radiation field does not determine the conditions for the ionizing photon escape. In contrast, other studies have found significant correlations between different physical properties and f_{esc} , which might indicate such physical processes. This includes highly concentrated star formation (high SFR surface densities), a high ionization parameter, an inhomogeneous ISM and dust distribution, and low amounts of dust (Verhamme et al. 2017; Gazagnes et al. 2018; Cen 2020; Flury et al. 2022b; Saldana-Lopez et al. 2022). Our finding does not exclude that radiative processes contribute to determining LyC escape, but it indicates that the hardness of the radiation field (over the energy range measured here) is not fundamental. For example, this suggests that LyC escape is not related to low-luminosity active galactic nuclei (AGNs). However, our differential study does not explain the origin of nebular He II emission, which is known to require sources of ionizing photons above 54 eV in amounts that are not predicted for normal stellar populations (see, e.g., Shirazi & Brinchmann 2012; Stasińska et al. 2015; Schaerer et al. 2019).

4. Summary

We have empirically investigated the hardness of the ionizing spectra between 54.4 and 13.6 eV ($Q_{\text{He}^+}/Q_{\text{H}}$, i.e., the ionization potential of He^+ and that of H) as a function of the LyC escape fraction, $f_{\text{esc}}(\text{LyC})$, of a large sample of star-forming galaxies at low redshift for which LyC are available from HST observations. Optical recombination lines of He II 4686 Å and H β 4861 Å from SDSS and X-shooter spectra were used to determine $Q_{\text{He}^+}/Q_{\text{H}}$ and its dependence on $f_{\text{esc}}(\text{LyC})$. The underlying effect of metallicity was also considered.

We built stacked spectra in bins of $f_{\text{esc}}(\text{LyC})$ and metallicity, allowing us to study the behavior of He II/H β across a wide range of $f_{\text{esc}}(\text{LyC}) \simeq 0\text{--}0.9$ and $12 + \log(\text{O}/\text{H}) = 7.5\text{--}8.6$. We find that the intensity of He II/H β does not depend on $f_{\text{esc}}(\text{LyC})$. Very strong leakers ($\langle f_{\text{esc}}(\text{LyC}) \rangle \simeq 0.5$) and non-leakers ($\langle f_{\text{esc}}(\text{LyC}) \rangle \simeq 0$) have similar intensities of He II/H β on average, about He II/H $\beta \simeq 0.01\text{--}0.02$. Instead, we find that He II/H β is primarily driven by variations in metallicity, where He II/H β increases with decreasing $12 + \log(\text{O}/\text{H})$, in particular, for $12 + \log(\text{O}/\text{H}) < 8.0$, as known from previous studies (e.g., Shirazi & Brinchmann 2012; Schaerer et al. 2019). At comparable metallicities, strong LyC leakers present roughly the same He II/H β intensities as non- or weak leakers and as in many other normal star-forming galaxies, where nebular He II is detected.

In short, our results demonstrate that $Q_{\text{He}^+}/Q_{\text{H}}$ does not correlate with $f_{\text{esc}}(\text{LyC})$. This implies that strong LyC emitters do not show harder ionizing spectra than nonleakers at similar metallicity. Future studies will address other hardness or softness indicators of the radiation field and will more broadly examine the nebular properties of the galaxies from the LzLCs.

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¹ Department of Astronomy, University of Geneva, 51 Chemin Pegasi, 1290 Versoix, Switzerland

e-mail: rui.marquescoelho@unige.ch

² CNRS, IRAP, 14 Avenue E. Belin, 31400 Toulouse, France

³ Instituto de Investigación Multidisciplinar en Ciencia y Tecnología, Universidad de La Serena, Raúl Bitrán 1305, La Serena, Chile

⁴ Departamento de Física y Astronomía, Universidad de La Serena, Avda. Juan Cisternas 1200, La Serena, Chile

⁵ Institut d’astrophysique de Paris, CNRS UMR7095, Sorbonne Université, 98bis Boulevard Arago, 75014 Paris, France

⁶ School of Earth & Space Exploration, Arizona State University, Tempe, AZ 85287, USA

⁷ Department of Astronomy, The University of Texas at Austin, 2515 Speedway, Stop C1400, Austin, TX 78712-1205, USA

⁸ Department of Astronomy, University of Massachusetts, Amherst, MA 01003, USA

⁹ INAF – Osservatorio Astronomico di Padova, Vicolo dell’ Osservatorio, 5, 35122 Padova, Italy

¹⁰ Department of Astronomy, Oskar Klein Centre, Stockholm University, 106 91 Stockholm, Sweden

¹¹ Department of Physics and Astronomy, Johns Hopkins University, 3400 North Charles Street, Baltimore, MD 21218, USA

¹² Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA

¹³ Bogolyubov Institute for Theoretical Physics, National Academy of Sciences of Ukraine, 14b Metrolohichna str., Kyiv 03143, Ukraine

¹⁴ Astronomy Department, Williams College, Williamstown, MA 01267, USA

¹⁵ University of Michigan, Department of Astronomy, 323 West Hall, 1085 S. University Ave, Ann Arbor, MI 48109, USA

¹⁶ Department of Physics and Astronomy, Minnesota State University, Mankato, MN 56001, USA

¹⁷ Infrared Processing and Analysis Center, California Institute of Technology, Pasadena, CA 91125, USA

¹⁸ Astronomy Department, University of Virginia, PO Box 400325, Charlottesville, VA 22904-4325, USA

¹⁹ Institut für Physik und Astronomie, Universität Potsdam, Karl-Liebknecht-Str. 24/25, 14476 Potsdam, Germany