



HAL
open science

Luminosity-dependent unification of active galactic nuclei and the X-ray Baldwin effect

C. Ricci, S. Paltani, H. Awaki, P. -O. Petrucci, y. Ueda, M. Brightman

► **To cite this version:**

C. Ricci, S. Paltani, H. Awaki, P. -O. Petrucci, y. Ueda, et al.. Luminosity-dependent unification of active galactic nuclei and the X-ray Baldwin effect. *Astronomy and Astrophysics - A&A*, EDP Sciences, 2013, 553, 10.1051/0004-6361/201220324 . insu-03617337

HAL Id: insu-03617337

<https://hal-insu.archives-ouvertes.fr/insu-03617337>

Submitted on 24 Mar 2022

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution| 4.0 International License

Luminosity-dependent unification of active galactic nuclei and the X-ray Baldwin effect

C. Ricci^{1,2,3}, S. Paltani^{1,2}, H. Awaki⁴, P.-O. Petrucci⁵, Y. Ueda³, and M. Brightman⁶

¹ ISDC Data Centre for Astrophysics, Université de Genève, ch. d'Ecogia 16, 1290 Versoix, Switzerland

² Observatoire de Genève, Université de Genève, 51 Ch. des Maillettes, 1290 Versoix, Switzerland

³ Department of Astronomy, Kyoto University, Oiwake-cho, Sakyo-ku, 606-8502 Kyoto, Japan
e-mail: ricci@kusastro.kyoto-u.ac.jp

⁴ Department of Physics, Ehime University, 790-8577 Matsuyama, Japan

⁵ UJF-Grenoble 1/CNRS-INSU, Institut de Planétologie et d'Astrophysique de Grenoble (IPAG) UMR 5274, 38041 Grenoble, France

⁶ Max-Planck-Institut für extraterrestrische Physik, Giessenbachstrasse 1, 85748 Garching bei München, Germany

Received 3 September 2012 / Accepted 31 January 2013

ABSTRACT

The existence of an anti-correlation between the equivalent width (EW) of the narrow core of the iron $K\alpha$ line and the luminosity of the continuum (i.e., the X-ray Baldwin effect) in type I active galactic nuclei has been confirmed in recent years by several studies carried out with *XMM-Newton*, *Chandra* and *Suzaku*. However, no general consensus on the origin of this trend has been reached so far. Several works have proposed the decrease of the covering factor of the molecular torus with the luminosity (in the framework of the luminosity-dependent unification models) as a possible explanation for the X-ray Baldwin effect. Using the fraction of obscured sources measured by recent X-ray and infrared (IR) surveys as a proxy of the half-opening angle of the torus and recent Monte Carlo simulations of the X-ray radiation reprocessed by a structure with a spherical-toroidal geometry, we test the hypothesis that the X-ray Baldwin effect is related to the decrease of the half-opening angle of the torus with the luminosity. Simulating the spectra of an unabsorbed population with a luminosity-dependent covering factor of the torus as predicted by recent X-ray surveys, we find that this mechanism is able to explain the observed X-ray Baldwin effect. Fitting the simulated data with a log-linear $L_{2-10\text{keV}}$ -EW relation, we found that in the Seyfert regime ($L_{2-10\text{keV}} \leq 10^{44.2} \text{ erg s}^{-1}$) luminosity-dependent unification produces a slope consistent with the observations for average values of the equatorial column densities of the torus of $\log N_{\text{H}}^{\text{T}} \gtrsim 23.1$, and can reproduce both the slope and the intercept for $\log N_{\text{H}}^{\text{T}} \approx 23.2$. Lower values of N_{H}^{T} are obtained assuming the decrease of the covering factor of the torus with the luminosity extrapolated from IR observations ($22.9 \lesssim \log N_{\text{H}}^{\text{T}} \lesssim 23$). In the quasar regime ($L_{2-10\text{keV}} > 10^{44.2} \text{ erg s}^{-1}$), a decrease of the covering factor of the torus with the luminosity slower than that observed in the Seyfert regime (as found by recent hard X-ray surveys) is able to reproduce the observations for $23.2 \leq \log N_{\text{H}}^{\text{T}} \leq 24.2$.

Key words. galaxies: Seyfert – quasars: general – galaxies: active – X-rays: galaxies

1. Introduction

There are two main observational signatures of reprocessed radiation in the X-ray spectra of active galactic nuclei (AGN): a Compton hump peaking around 30 keV and a fluorescent iron $K\alpha$ line. While the Compton hump is produced only if the reprocessing material is Compton thick (CT, $N_{\text{H}} \gtrsim 10^{24} \text{ cm}^{-2}$), the iron $K\alpha$ line is also created in Compton-thin material (e.g., Matt et al. 2003). The iron $K\alpha$ line is often observed as the superposition of two different components: a broad component and a narrow one. The broad component has a full width at half-maximum (FWHM) of $\gtrsim 30\,000 \text{ km s}^{-1}$ and is thought to be created close to the black hole in the accretion disk (e.g., Fabian et al. 2000), or to be related to the presence of features created by partially covering warm absorbers in the line of sight (e.g., Turner & Miller 2009; Miyakawa et al. 2012). While the broad component is observed in only ~ 35 – 45% of bright nearby AGN (e.g., de La Calle Pérez et al. 2010), the narrow (FWHM $\sim 2000 \text{ km s}^{-1}$, Shu et al. 2010) core of the iron line has been found to be almost ubiquitous (e.g., Nandra et al. 2007; Singh et al. 2011). This component peaks at 6.4 keV (e.g., Yaqoob & Padmanabhan 2004), which points to the line being produced in

cold neutral material. This material has often been identified as circumnuclear matter located at several thousand gravitational radii from the supermassive black hole, and it is likely related to the putative molecular torus (e.g., Nandra 2006), although a contribution of the outer part of the disk (e.g., Petrucci et al. 2002) or of the broad-line region (BLR, e.g., Bianchi et al. 2008) cannot be excluded. In a recent study, Shu et al. (2011) found that the weighted mean of the ratio between the FWHM of the narrow Fe $K\alpha$ line and that of optical lines produced in the BLR is ≈ 0.6 . This implies that the size of the iron $K\alpha$ -emitting region is on average about three times that of the BLR, indicating that most of the narrow iron $K\alpha$ emission is produced in the putative molecular torus.

One of the most interesting characteristics of the narrow component of the iron $K\alpha$ line is the inverse correlation between its equivalent width (EW) and the X-ray luminosity (e.g., Iwasawa & Taniguchi 1993). The existence of an anti-correlation between the EW of a line and the luminosity of the AGN continuum was found for the first time in the UV by Baldwin (1977) for the CIV $\lambda 1549$ line and dubbed the Baldwin effect. A similar trend was later found for several other emission lines, such as Ly α , [C III] $\lambda 1908$, Si IV $\lambda 1396$, Mg II $\lambda 2798$

Table 1. Summary of the most recent studies (along with the original work of Iwasawa & Taniguchi 1993) of the X-ray Baldwin effect.

(1)	(2)	(3)	(4)	(5)
Reference	A	B	Sample	Observatory/instrument
<i>Fits per observation</i>				
Shu et al. (2012)	$1.58 \pm 0.03^*$	-0.18 ± 0.03	32 (RQ)	Chandra/HEG
Shu et al. (2010)	$1.58 \pm 0.03^*$	-0.22 ± 0.03	33 (RQ+RL)	Chandra/HEG
Bianchi et al. (2007)	1.73 ± 0.03	-0.17 ± 0.03	157 (RQ)	XMM-Newton/EPIC
Jiang et al. (2006)	–	-0.20 ± 0.04	101 (RL+RQ)	XMM-Newton/EPIC + Chandra/HEG
Jiang et al. (2006)	–	-0.10 ± 0.05	75 (RQ)	XMM-Newton/EPIC + Chandra/HEG
Jiménez-Bailón et al. (2005)	–	-0.06 ± 0.20	38 (RQ)	XMM-Newton/EPIC
Zhou & Wang (2005)	–	-0.19 ± 0.04	66 (RQ+RL)	XMM-Newton/EPIC
Page et al. (2004)	–	-0.17 ± 0.08	53 (RQ+RL)	XMM-Newton/EPIC
Iwasawa & Taniguchi (1993)	–	-0.20 ± 0.03	37 (RQ+RL)	Ginga
<i>Fits per source</i>				
Shu et al. (2012)	$1.64 \pm 0.03^*$	-0.11 ± 0.03	32 (RQ)	Chandra/HEG
Shu et al. (2010)	$1.63 \pm 0.04^*$	-0.13 ± 0.04	33 (RQ+RL)	Chandra/HEG

Notes. The model commonly used to fit the data is $\log EW = A + B \log L_{X,44}$, where $L_{X,44}$ is the 2–10 keV luminosity in units of 10^{44} erg s $^{-1}$. The table is divided into works for which the fit was done *per observation* (i.e., using all the available observations for every source of the sample) and *per source* (i.e., averaging the values of EW and L_X of different observations of the same source). The table lists (1) the reference; (2) the value of the intercept; (3) the value of the slope; (4) the number and type of objects in the sample; (5) the observatory and instrument used. Dashes are given when the value is not available. ^(*) These values of the intercepts were corrected in order to be consistent with Eq. (5), because those reported in Shu et al. (2010, 2012) are relative to the 2–10 keV luminosity in units of 10^{43} erg s $^{-1}$ ($L_{X,43}$).

(Dietrich et al. 2002), UV iron emission lines (Green et al. 2001), mid-infrared (mid-IR) lines such as [AR III] $\lambda 8.99 \mu\text{m}$, [S IV] $\lambda 10.51 \mu\text{m}$ and [Ne II] $\lambda 12.81 \mu\text{m}$ (Hönig et al. 2008), and forbidden lines such as [O II] $\lambda 3727$ and [Ne V] $\lambda 3426$ (Croom et al. 2002). The slope of the Baldwin effect for most of these lines has been shown to be steeper for the lines originating from higher ionization species (e.g., Dietrich et al. 2002). The origin of the Baldwin effect is still unknown, although several possible explanations have been put forward, and might be different for lines originating in different regions of the AGN. For the lines produced in the BLR, the origin might be related to the lower ionization and photoelectric heating in the BLR gas of more luminous objects (e.g., Netzer et al. 1992).

Using *Ginga* observations of 37 AGN, Iwasawa & Taniguchi (1993) found in the X-rays the first evidence of an anti-correlation between the EW of the iron $K\alpha$ line and the 2–10 keV luminosity ($EW \propto L_X^{-0.20 \pm 0.04}$). This trend is usually called the *X-ray Baldwin effect* or the *Iwasawa-Taniguchi effect*. Using ASCA observations, Nandra et al. (1997) confirmed the existence of such an anti-correlation and argued that most of the effect could be explained by variations of the broad component of the iron $K\alpha$ line with the luminosity. The advent of *XMM-Newton*, *Chandra*, and *Suzaku* made clear however that most of the observed X-ray Baldwin effect is due to the narrow core of the iron $K\alpha$ line (e.g., Page et al. 2004; Shu et al. 2010; Fukazawa et al. 2011). The significance of the effect was questioned by Jiménez-Bailón et al. (2005), who discussed the possible importance of contamination from radio-loud (RL) AGN, which have on average larger X-ray luminosities and smaller signatures of reprocessed radiation (e.g., Reeves & Turner 2000). However, Grandi et al. (2006), studying *BeppoSAX* observations of RL AGN, also found evidence of an X-ray Baldwin effect. This, together with the study of a large sample of radio-quiet (RQ) AGN carried out by Bianchi et al. (2007) using *XMM-Newton* data, showed that the X-ray Baldwin effect is not a mere artifact. Jiang et al. (2006) suggested that the X-ray Baldwin effect might be due to the delay between the variability of the AGN primary energy source and that of the reprocessing material located farther away. Using *Chandra/HEG*

observations, Shu et al. (2010, 2012) confirmed the importance of variability, showing that averaging the values of EW obtained over multiple observations of individual sources would significantly attenuate the anti-correlation from $EW \propto L_X^{-0.18 \pm 0.03}$ to $EW \propto L_X^{-0.11 \pm 0.03}$. *Chandra/HEG* observations (Shu et al. 2010) have also shown that the normalization of the X-ray Baldwin effect is lower (i.e., the average EW of the narrow Fe $K\alpha$ component is smaller) than that obtained by previous works performed using *XMM-Newton*. In Table 1 we report the values of the slope and the intercept obtained from the most recent *XMM-Newton/EPIC* and *Chandra/HEG* works, along with the original *Ginga* work of Iwasawa & Taniguchi (1993).

An intriguing possibility is that the X-ray Baldwin effect is related to the decrease of the covering factor of the torus with luminosity (e.g., Page et al. 2004; Zhou & Wang 2005) in the frame of the luminosity-dependent unification schemes. The EW of the iron $K\alpha$ line is in fact expected to be proportional to the covering factor of the torus (e.g., Krolik et al. 1994; Ikeda et al. 2009), so that a covering factor of the torus decreasing with the luminosity might in principle be able to explain the X-ray Baldwin effect. Luminosity-dependent unification models have been proposed to explain the decrease of the fraction of obscured objects (f_{obs}) with the increase of the AGN output power. The first suggestion of the existence of a relation between f_{obs} and the luminosity was put forward about 30 years ago (Lawrence & Elvis 1982). Since then the idea that the covering factor of the obscuring material decreases with luminosity has been gaining observational evidence from radio (e.g., Grimes et al. 2004), infrared (e.g., Treister et al. 2008; Mor et al. 2009; Gandhi et al. 2009), optical (e.g., Simpson 2005), and X-ray (e.g., Ueda et al. 2003; Beckmann et al. 2009; Ueda et al. 2011) studies of AGN. Although luminosity-dependent unification models have long been suspected to play a major role in the X-ray Baldwin effect, no quantitative estimation of this effect has been performed so far.

In this work we quantify for the first time the influence of the decrease of the covering factor of the torus with the luminosity on the EW of the iron $K\alpha$ line. Using the Monte Carlo spectral simulations of a torus with a variable half-opening angle

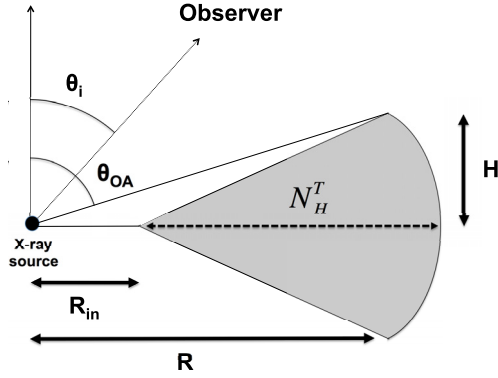


Fig. 1. Schematic representation of the torus geometry considered. The angle θ_i is the inclination of the line of sight with respect to the torus axis, while θ_{OA} and N_H^T are the half-opening angle and the equatorial column density of the molecular torus, respectively. H is the maximum height of the torus, while R_{in} and R are its inner and outer radius, respectively.

(θ_{OA}) (recently presented by Ikeda et al. 2009 and Brightman & Nandra 2011) together with the most recent and comprehensive observations in different energy bands of the decrease of f_{obs} with the luminosity, we show that this mechanism is able to explain the X-ray Baldwin effect. The paper is organized as follows. In Sect. 2 we present in detail our spectral simulations, and in Sect. 3 we compare the slopes and intercepts obtained by our simulations with those measured by high-sensitivity *Chandra*/HEG observations, as they provide the best energy resolution available to date. In Sect. 4 we discuss our results, and in Sect. 5 we present our conclusions.

2. Simulations

2.1. The relation between θ_{OA} and L_X

As a proxy of the relationship between the half-opening angle of the molecular torus θ_{OA} (see Fig. 1) and the intrinsic X-ray luminosity of AGN, we used the variation of the fraction of obscured sources f_{obs} with luminosity measured by the recent medium (2–10 keV) and hard (15–55 keV) X-ray surveys of Hasinger (2008) and Burlon et al. (2011). From a combination of surveys performed by HEAO-1, ASCA, *BeppoSAX*, *XMM-Newton*, and *Chandra* in the luminosity range $42 \leq \log L_X \leq 46$, Hasinger (2008) found in the 2–10 keV band

$$f_{obs} \simeq -0.226 \log L_X + 10.342, \quad (1)$$

where L_X is the luminosity in the 2–10 keV energy range. In the 15–55 keV band, using *Swift*/BAT to study AGN in the luminosity range $42 \leq \log L_X \leq 45$ and fitting the data with

$$f_{obs} \simeq R_{low} e^{-L_{HX}/L_C} + R_{high} (1 - e^{-L_{HX}/L_C}), \quad (2)$$

where L_{HX} is the luminosity in the 15–55 keV band, Burlon et al. (2011) obtained $R_{low} = 0.8$, $R_{high} = 0.2$, and $L_C = 10^{43.7} \text{ erg s}^{-1}$. We also used the results of the IR work of Maiolino et al. (2007), who found an anti-correlation between the ratio $\lambda L_\lambda(6.7 \mu\text{m})/\lambda L_\lambda(5100 \text{ \AA})$ and the [O III] $\lambda 5007 \text{ \AA}$ line luminosity. This trend was interpreted as an effect of the decrease of the covering factor of the circumnuclear dust as a function of luminosity, with f_{obs} varying as

$$f_{obs} \simeq \frac{1}{1 + L_{opt}^{0.414}}, \quad (3)$$

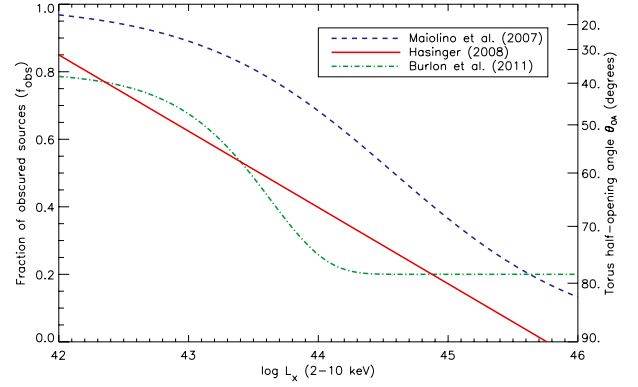


Fig. 2. Variation of the fraction of obscured sources (f_{obs}) and of the half-opening angle of the torus (θ_{OA}) with the 2–10 keV luminosity for three of the most recent medium X-ray (2–10 keV, Hasinger 2008), hard X-ray (15–55 keV, Burlon et al. 2011), and IR (Maiolino et al. 2007) studies (Eqs. (1)–(3)).

where $L_{opt} = (\lambda L_\lambda(5100 \text{ \AA}))/10^{45.63}$. Following Maiolino et al. (2007), it is possible to rewrite L_{opt} as a function of the X-ray luminosity: $L_{opt} = L_X^{1.39}/10^{61.97}$.

The fraction of obscured sources at a given luminosity can be easily related to the half-opening angle of the torus using

$$\theta_{OA} = \cos^{-1}(f_{obs}). \quad (4)$$

For Eq. (2) we converted the hard X-ray luminosity to the luminosity in the 2–10 keV band, assuming a photon index $\Gamma = 1.9$ (e.g., Beckmann et al. 2009). The three $\theta_{OA}-L_X$ relationships used in this work are shown in Fig. 2.

2.2. Spectral simulations and fitting

Ikeda et al. (2009) recently presented Monte Carlo simulations of the reprocessed X-ray emission of an AGN surrounded by a three-dimensional spherical-toroidal structure. The simulations were performed using the ray-tracing method, taking into account Compton down-scattering and absorption. They are stored in tables, so that they can be used to perform spectral fitting. The free parameters of this model are the half-opening angle of the torus θ_{OA} , the line-of-sight inclination angle θ_i , the torus equatorial column density N_H^T , and the photon index Γ of the primary continuum. We note that N_H^T should not be confused with the observed hydrogen column density N_H . In all objects, $N_H \leq N_H^T$, with $N_H = N_H^T$ only if $\theta_i = 90^\circ$. In Fig. 1 a schematic representation of the geometry considered is shown. The dependence of N_H on the inclination angle is given by Eq. (3) of Ikeda et al. (2009). In the model of Ikeda et al. (2009), the ratio R_{in}/R is fixed to 0.01 and the inclination angle of the observer θ_i can vary between 1° and 89° , N_H^T between 10^{22} and 10^{25} cm^{-2} , while θ_{OA} spans the range between 10° and 70° . Because of the assumed dependence of the half-opening angle on the luminosity, the interval of values of θ_{OA} allowed by the model limits the range of luminosities we can probe. In the following we will consider luminosities above $\log L_X = 42$, because below this value few AGN are detected and several works point towards a possible disappearance of the molecular torus (e.g., Elitzur & Shlosman 2006). The upper-limit luminosity we can reach depends on the $\theta_{OA}-L_X$ relationship used and, considering $\theta_{OA}^{max} = 70^\circ$, is $\log L_{max} = 44.2$ for Hasinger (2008), $\log L_{max} = 43.8$ for Burlon et al. (2011), and $\log L_{max} = 45.0$ for Maiolino et al. (2007). Due to its limitations in the values of θ_{OA} permitted, the model

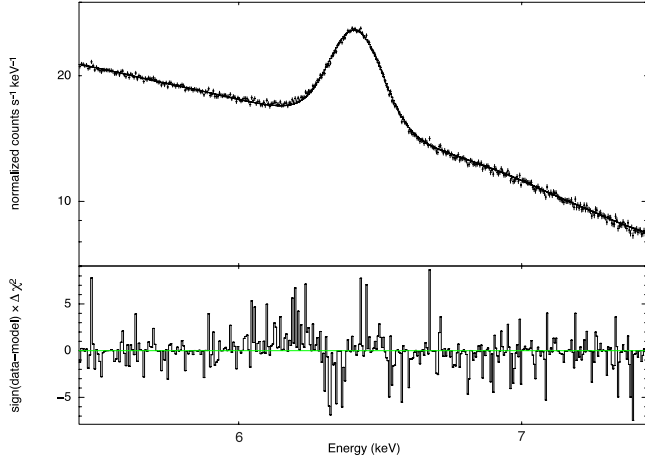


Fig. 3. *Top panel:* extract in the 5.5–7.5 keV region of a spectrum simulated using the model of Ikeda et al. (2009). The model has a torus with a half-opening angle of $\theta_{\text{OA}} = 46.2^\circ$ (equivalent to $\log L_X = 42.7$, according to the relationship of Hasinger 2008), an equatorial column density of $\log N_{\text{H}}^{\text{T}} = 23.8$, and an inclination angle of $\theta_i = 1^\circ$. The continuous line represents the fit to the simulated spectrum using for the continuum the same model we used for the simulations and a Gaussian line for the iron $K\alpha$ fluorescent line. *Bottom panel:* contribution to the chi-squared for the best fit to the simulated spectrum.

of Ikeda et al. (2009) restricts most of the simulations to the Seyfert regime ($\log L_X \leq \log L_X^{\text{Q}} = 44.2$).

To extend our study to the quasar regime ($L_X > L_X^{\text{Q}}$), we simulated the X-ray Baldwin effect using the spectral model of Brightman & Nandra (2011). This model considers the same geometry as Ikeda et al. (2009) but has $R_{\text{in}} = 0$, a θ_i -independent N_{H} , and it allows different values of θ_{OA} (26° to 84°). The higher maximum half-opening angle permitted by this model allows to reach higher maximum luminosities in the simulations for the relationships of Hasinger (2008; $\log L_{\text{max}} = 45.3$) and Maiolino et al. (2007; $\log L_{\text{max}} = 46.2$), while that of Burlon et al. (2011) is flat ($f_{\text{obs}} \simeq 0.2$) in the quasar regime. However, the lower boundary of θ_{OA} in the model of Brightman & Nandra (2011) limits the lower luminosity we can reach in the simulations, in particular $\log L_{\text{min}} \simeq 43$ for the relation of Maiolino et al. (2007). Thus, we used the model of Ikeda et al. (2009) to study the X-ray Baldwin effect in the Seyfert regime ($42 \leq \log L_X \leq 44.2$) and the model of Brightman & Nandra (2011) to probe the quasar regime ($\log L_X > 44.2$).

To estimate the influence of the decreasing covering factor of the torus on the EW of the iron line, we used the physical torus models to simulate a large number of spectra using the $\theta_{\text{OA}}-L_X$ relationships reported in Eqs. (1)–(3). Using the three different $\theta_{\text{OA}}-L_X$ relationships, we extrapolated the value of θ_{OA} for each luminosity bin ($\Delta \log L_X = 0.01$). We fixed the equatorial column density of the torus N_{H}^{T} to 26 different values, spanning $10^{22.5} \text{ cm}^{-2}$ and 10^{25} cm^{-2} , with a step of $\Delta \log N_{\text{H}}^{\text{T}} = 0.1$. To simulate an unabsorbed population, similarly to what is usually used to determine the X-ray Baldwin effect, we considered inclination angles θ_i between 1° and θ_{OA} for each value of L_X and N_{H}^{T} , with a binning of $\Delta \theta_i = 3^\circ$. In XSPEC 12.7.1 (Arnaud 1996), we simulated spectra for each of the three $\theta_{\text{OA}}-L_X$ relationships and for each bin of column density. For the continuum we used a power law with $\Gamma = 1.9$ (e.g., Beckmann et al. 2009). Our choice of the photon index does not significantly affect the simulations; by adding a scatter of $\Delta \Gamma = 0.3$, we found $EW-L_X$ trends consistent with those obtained without scatter. The metallicity was

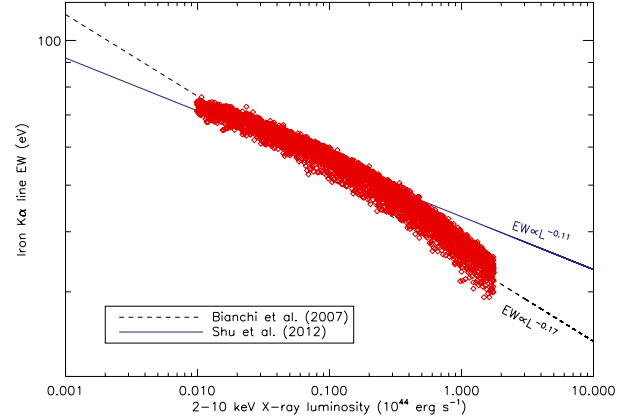


Fig. 4. Equivalent width of the iron $K\alpha$ line versus the X-ray luminosity obtained simulating a torus with an equatorial column density of $\log N_{\text{H}}^{\text{T}} = 23.1$ and a covering factor decreasing with the luminosity. The $\theta_{\text{OA}}-L_X$ relationship used here is that of Hasinger (2008; see Fig. 2). The relations of Bianchi et al. (2007, black dashed line) and Shu et al. (2012, blue line, obtained averaging different observations of the same source) are also shown for comparison (see also Table 1). The normalization of the relation of Bianchi et al. (2007) has been fixed to an arbitrary value for comparison.

set to the solar value in all our simulations, and the value of the normalization of the reflected component was fixed to that of the continuum.

We fitted the simulated spectra in the 0.3–10 keV band using the same model applied for the simulations, substituting the iron $K\alpha$ line component with a Gaussian line. We obtained a good fit for all the simulations, with a reduced chi-squared of $\chi_v^2 \lesssim 1.1$. We show in Fig. 3 an example of a typical fit. The model used for the continuum fitting does not significantly affect the results, and the EW of the iron $K\alpha$ line differs only $\sim 4\%$ on average when using an alternative model like pexrav (Magdziarz & Zdziarski 1995). We evaluated the EW of the iron $K\alpha$ line and studied its relationship with the luminosity. As an example, we show in Fig. 4 the simulated X-ray Baldwin effect obtained using the model of Ikeda et al. (2009) and the $\theta_{\text{OA}}-L_X$ relationship of Hasinger (2008) for $\log N_{\text{H}}^{\text{T}} = 23.1$, together with the fit to the X-ray Baldwin effect obtained by the recent works of Bianchi et al. (2007) and Shu et al. (2012). The spread in EW for a given luminosity is due to the range of values of θ_i that we considered: larger values of EW usually correspond to lower values of θ_i .

Most of the studies of the X-ray Baldwin effect performed in recent years have used a relationship of the type

$$\log EW = A + B \log L_{X,44} \quad (5)$$

to fit the $EW-L_X$ trend, where $L_{X,44}$ is the luminosity in units of $10^{44} \text{ erg s}^{-1}$. In order to compare the simulated $EW-L_X$ trend with that observed in unabsorbed populations of AGN, we fitted the simulated data, for each value of N_{H}^{T} , with Eq. (5) using the weighted least-square method, with weights of $w = \sin \theta_i$. This makes it possible to account for the non-uniform probability of randomly observing an AGN within a certain solid angle from the polar axis.

3. The X-ray Baldwin effect

Most of the studies on the X-ray Baldwin effect have found, using Eq. (5), a slope of $B \sim -0.2$ (e.g., Iwasawa & Taniguchi 1993; Page et al. 2004; Bianchi et al. 2007). However, in all

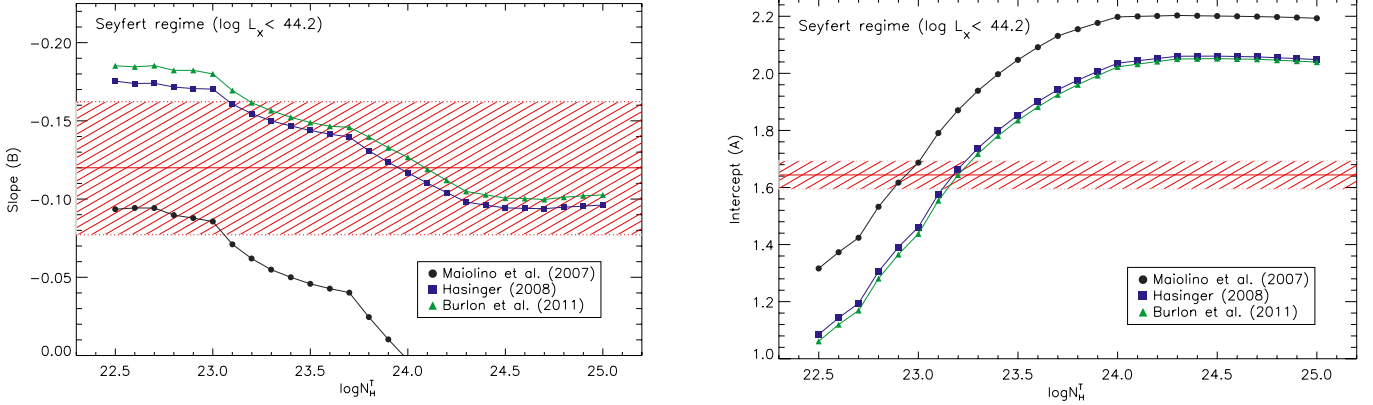


Fig. 5. *Left panel:* value of the slope (B) of the X-ray Baldwin effect obtained simulating the variation of the reprocessed X-ray radiation with the luminosity for tori with different values of the equatorial column density N_{H}^{T} (in the range $\log N_{\text{H}}^{\text{T}} = 22.5\text{--}25$) for the different $\theta_{\text{OA}}\text{--}L_{\text{X}}$ relationships shown in Fig. 2 (Eqs. (1)–(3)) in the Seyfert regime ($42 < \log L_{\text{X}} \leq 44.2$). The model used is that of Ikeda et al. (2009), and the simulated data were fitted with Eq. (5). The red line represents the value of the slope ($B = -0.12 \pm 0.04$) obtained by fitting in the Seyfert regime the data *per source* reported in Shu et al. (2010); the red shadowed area represents the 1σ contour of the slope. *Right panel:* same as left panel, but considering the intercept (A) of the X-ray Baldwin effect obtained by our simulations. The red line represents the value of the intercept ($A = 1.64 \pm 0.05$) obtained by the fit to the *Chandra*/HEG data in the Seyfert regime, while the red shadowed area represents its 1σ error.

these works the values of EW and L_{X} are obtained from individual observations of sources. As pointed out by Jiang et al. (2006) and confirmed by Shu et al. (2010, 2012), flux variability might play an important role in the observed X-ray Baldwin effect. Shu et al. (2012) recently found a clear anti-correlation between the EW of the iron $\text{K}\alpha$ line and the luminosity for individual sources that had several *Chandra*/HEG observations. Using a sample of 32 RQ AGN, they also found that the fit *per source* (i.e., averaging different observations of the same object) results in a significantly flatter slope ($B = -0.11 \pm 0.03$) than that done *per observation* (i.e., using all the available observations for every source of the sample; $B = -0.18 \pm 0.03$).

To compare the slope obtained by the simulations in the two different luminosity bands with real data, we fitted the data *per source* (33 AGN) of Shu et al. (2010; obtained fixing $\sigma = 1$ eV) with Eq. (5) in the Seyfert and quasar regime. Consistent with what was done by Shu et al. (2010) we did not use the three AGN for which only upper limits of EW were obtained. Using the 28 AGN in the Seyfert regime, we found that the best fit to the X-ray Baldwin effect is given by $A = 1.64 \pm 0.05$ and $B = -0.12 \pm 0.04$. *Chandra*/HEG observations are available for only five objects in the quasar regime, and at these luminosities we obtained $A = 1.5 \pm 0.3$ and $B = -0.16 \pm 0.22$.

3.1. Seyfert regime

Fitting the simulated data with Eq. (5) in the Seyfert regime, we found that the slope obtained becomes flatter for increasing values of the equatorial column density of the torus (Fig. 5, left panel). This is related to the fact that the iron $\text{K}\alpha$ EW is tightly connected to both N_{H}^{T} and θ_{OA} , so that EW is more strongly dependent on θ_{OA} (and thus on its variation) for low values of N_{H}^{T} , which results in a steeper slope. On the other hand, the dependence on θ_{OA} becomes weaker and the slope flatter when N_{H}^{T} increases. As shown in Fig. 5 (left panel), while the values of B obtained by the $\theta_{\text{OA}}\text{--}L_{\text{X}}$ relationships of Burlon et al. (2011) and Hasinger (2008) are similar along the whole range of N_{H}^{T} considered, flatter slopes are obtained for that of Maiolino et al. (2007). In particular for the latter relationship the correlation becomes positive for $\log N_{\text{H}}^{\text{T}} \gtrsim 24$.

By combining the observations of the X-ray Baldwin effect with our simulations, we can extrapolate the average value of the equatorial column density of the torus of the unobscured AGN in the *Chandra*/HEG sample of Shu et al. (2010). A slope consistent within 1σ with our fit to the X-ray Baldwin effect in the Seyfert regime is obtained for $N_{\text{H}}^{\text{T}} \geq 10^{23.1} \text{ cm}^{-2}$ for the $\theta_{\text{OA}}\text{--}L_{\text{X}}$ relationship of Hasinger (2008), for $N_{\text{H}}^{\text{T}} \gtrsim 10^{23.2} \text{ cm}^{-2}$ for that of Burlon et al. (2011), and for $N_{\text{H}}^{\text{T}} \lesssim 10^{23} \text{ cm}^{-2}$ for the relationship of Maiolino et al. (2007). Comparing the value of the intercept obtained by the simulations with that resulting from our fit to the X-ray Baldwin effect in the Seyfert regime, we found that only a narrow range of average equatorial column densities of the torus can reproduce the observations (right panel of Fig. 5). For the $\theta_{\text{OA}}\text{--}L_{\text{X}}$ relationship of Maiolino et al. (2007), we found $22.9 \lesssim \log N_{\text{H}}^{\text{T}} \lesssim 23$, while for those of Burlon et al. (2011) and Hasinger (2008) we obtained $\log N_{\text{H}}^{\text{T}} \approx 23.2$. Using both the values of A and B , it is possible to extract the average values of N_{H}^{T} that can explain the X-ray Baldwin effect. As can be seen from the two figures for both the $\theta_{\text{OA}}\text{--}L_{\text{X}}$ relationship of Hasinger (2008) and Burlon et al. (2011), the only value of column density consistent with both the observed intercept and slope is $\log N_{\text{H}}^{\text{T}} \approx 23.2$. For the IR $\theta_{\text{OA}}\text{--}L_{\text{X}}$ relationship of Maiolino et al. (2007) the values allowed are in the range $22.9 \lesssim \log N_{\text{H}}^{\text{T}} \lesssim 23$. A similar result is obtained studying the A/B chi-squared contour plot of our fit to the *Chandra*/HEG data. Using the model of Brightman & Nandra (2011) to simulate the X-ray Baldwin effect in the Seyfert regime, we obtained a range of N_{H}^{T} consistent with that found using the model of Ikeda et al. (2009) ($23.2 \lesssim \log N_{\text{H}}^{\text{T}} \lesssim 23.3$). The lower values of N_{H}^{T} needed to explain the X-ray Baldwin effect using the relationship of Maiolino et al. (2007) are due to the larger values of f_{obs} (and thus of θ_{OA}) predicted by Eq. (3). This is again due to the fact that EW depends on both N_{H}^{T} and θ_{OA} , so that increasing the latter, one would obtain lower values of the former.

3.2. Quasar regime

To study the behavior of B in the quasar regime, we fitted the simulated data using Eq. (5) for each value of N_{H}^{T} for luminosities $L_{\text{X}} > L_{\text{X}}^{\text{Q}}$. Our simulations show that EW decreases

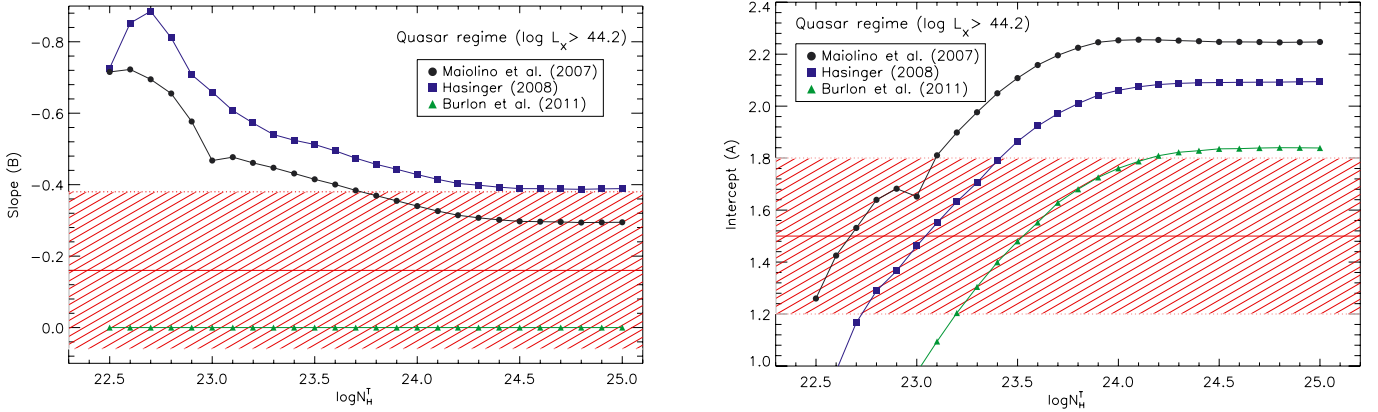


Fig. 6. *Left panel:* value of the slope (B) of the X-ray Baldwin effect obtained simulating the variation of the reprocessed X-ray radiation with the luminosity for tori with different values of the equatorial column density N_{H}^{T} for the different $\theta_{\text{OA}}-L_{\text{X}}$ relationships shown in Fig. 2 (Eqs. (1)–(3)) in the quasar regime ($\log L_{\text{X}} > 44.2$). The model used is that of [Brightman & Nandra \(2011\)](#), and the simulated data were fitted with Eq. (5). The red line represents the value of the slope ($B = -0.16 \pm 0.22$), which was obtained by fitting in the quasar regime the data *per source* reported in [Shu et al. \(2010\)](#); the red shadowed area represents the 1σ contour of the slope. *Right panel:* same as left panel, but considering the intercept (A) of the X-ray Baldwin effect obtained by our simulations. The red line represents the value of the intercept ($A = 1.5 \pm 0.3$) obtained by the fit to the *Chandra*/HEG data in the quasar regime, while the red shadowed area represents its 1σ error.

more steeply in the quasar regime than at lower luminosities for the $\theta_{\text{OA}}-L_{\text{X}}$ relationships of [Hasinger \(2008\)](#) and [Maiolino et al. \(2007\)](#), with values of the slope of $B \lesssim -0.3$ (left panel of Fig. 6). The slopes obtained by the simulations are consistent with those found by fitting *Chandra*/HEG data in the quasar regime for the $\theta_{\text{OA}}-L_{\text{X}}$ relationship of [Maiolino et al. \(2007\)](#) for $N_{\text{H}}^{\text{T}} \geq 10^{23.7} \text{ cm}^{-2}$. The values of the slope expected using the relationship of [Hasinger \(2008\)](#) are steeper than the observed value for the whole range of column densities considered, while the flattening of the relationship of [Burlon et al. \(2011\)](#) at high luminosities results in a slope of $B \sim 0$ along the whole range of N_{H}^{T} , consistent within 1σ with the observations. The intercepts obtained using the relationship of [Maiolino et al. \(2007\)](#) are consistent with the observed value for $22.5 \leq \log N_{\text{H}}^{\text{T}} \leq 23.1$ (right panel of Fig. 6), with no overlap with the values of N_{H}^{T} needed by the slope. We obtained intercepts that are consistent with the observations for $22.7 \leq \log N_{\text{H}}^{\text{T}} \leq 23.4$ and $23.2 \leq \log N_{\text{H}}^{\text{T}} \leq 24.2$ for the relation of [Hasinger \(2008\)](#) and [Burlon et al. \(2011\)](#), respectively. Thus only the hard X-ray $\theta_{\text{OA}}-L_{\text{X}}$ relationship of [Burlon et al. \(2011\)](#) is able to explain (for $23.2 \leq \log N_{\text{H}}^{\text{T}} \leq 24.2$) at the same time both the intercept and the slope of the X-ray Baldwin effect at high luminosities.

4. Discussion

Since its discovery about 20 years ago, the existence of the X-ray Baldwin effect has been confirmed by several works performed with the highest spectral resolution available in the X-ray band. So far, several explanations have been proposed. [Jiang et al. \(2006\)](#) argued that the observed anti-correlation could be related to the delay of the reprocessed radiation with respect to the primary continuum. The response of the circumnuclear material to the irradiated flux is not simultaneous, and one should always take this effect into account when performing studies of reprocessed features such as the iron $K\alpha$ line or the Compton hump. [Shu et al. \(2010, 2012\)](#) have shown that averaging the values of L_{X} and EW for all the observations of each source results in a significantly flattened anti-correlation. However, by itself variability fails to fully account for the observed correlation ([Shu et al. 2012](#)).

4.1. Explaining the X-ray Baldwin effect with a luminosity-dependent covering factor of the torus

A mechanism often invoked to explain the X-ray Baldwin effect is the decrease of the covering factor of the molecular torus with the luminosity (e.g., [Page et al. 2004](#); [Bianchi et al. 2007](#)). The decrease of the fraction of obscured sources with the luminosity has been reported by several works performed at different wavelengths in the last decade, although some discordant results have been presented (e.g., [Dwelly & Page 2006](#); [Lawrence & Elvis 2010](#)). In this work, we have shown that the covering factor–luminosity relationships obtained in the medium and hard X-ray band can explain well the X-ray Baldwin effect in the $10^{42}-10^{44.2} \text{ erg s}^{-1}$ luminosity range. In particular, our simulations show that it is possible to reproduce the slope of the X-ray Baldwin effect with luminosity-dependent unification for average equatorial column densities of the torus of $\log N_{\text{H}}^{\text{T}} \gtrsim 23.1$, and both the slope and the intercept for $\log N_{\text{H}}^{\text{T}} \approx 23.2$ (Fig. 5). In the same luminosity range, the $\theta_{\text{OA}}-L_{\text{X}}$ IR relationship of [Maiolino et al. \(2007\)](#) can explain the X-ray Baldwin effect for $22.9 \leq \log N_{\text{H}}^{\text{T}} \leq 23$.

In the quasar regime, we have shown that, while the medium X-ray $\theta_{\text{OA}}-L_{\text{X}}$ relationship of [Hasinger \(2008\)](#) cannot explain the observations (Fig. 6), the slope obtained by the IR $\theta_{\text{OA}}-L_{\text{X}}$ relationship is consistent with our fit to the *Chandra*/HEG data in the same luminosity band. However, as for the latter relationship, because the range of values of N_{H}^{T} required to explain the slope and the intercept do not overlap, it cannot be considered a likely explanation. The hard X-ray $\theta_{\text{OA}}-L_{\text{X}}$ relationship of [Burlon et al. \(2011\)](#) is flat above L_{X}^{Q} , and would thus produce a constant EW of the iron $K\alpha$ line and a slope of $B \sim 0$ for the whole range of N_{H}^{T} , consistent within 1σ with the value found using *Chandra*/HEG data. The intercept we found using this relation is also consistent with the observational value for a large range of average equatorial column densities of the torus ($23.2 \leq \log N_{\text{H}}^{\text{T}} \leq 24.2$). The fraction of obscured AGN is not well constrained at high luminosities, thus any trend with the luminosity between that of [Burlon et al. \(2011\)](#) and that of [Hasinger \(2008\)](#) would be able to reproduce the observed slope. Because the relation of [Hasinger \(2008\)](#) produces negative

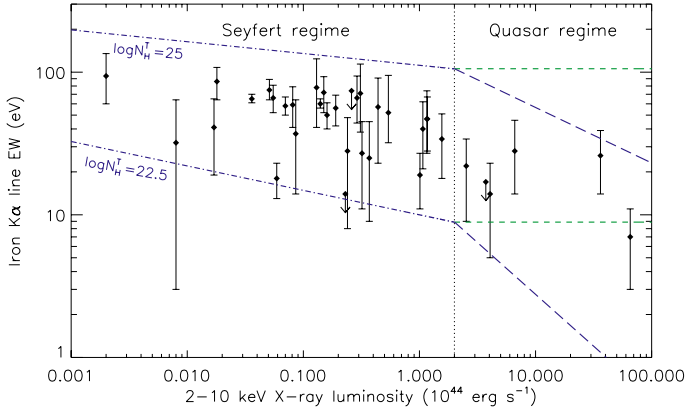


Fig. 7. Iron $K\alpha$ EW versus X-ray luminosities and predicted trends obtained for different values of the equatorial column density of the torus. The points are the values of the EW of the iron $K\alpha$ line reported by [Shu et al. \(2010\)](#) and obtained by averaging multiple *Chandra*/HEG observations of AGN. The two blue dash-dotted lines are the fits to our simulations of the X-ray Baldwin effect using the $\theta_{\text{OA}}-L_X$ relationship of [Hasinger \(2008\)](#) for the Seyfert regime ($\log EW = 1.01 - 0.17 \log L_{X,44}$ for $\log N_{\text{H}}^{\text{T}} = 22.5$, and $\log EW = 2.05 - 0.08 \log L_{X,44}$ for $\log N_{\text{H}}^{\text{T}} = 25$). The blue long-dashed lines represent the $EW-L_X$ relations obtained in the quasar regime ($\log EW = 1.17 - 0.73 \log L_{X,44}$ for $\log N_{\text{H}}^{\text{T}} = 22.5$ and $\log EW = 2.14 - 0.39 \log L_{X,44}$ for $\log N_{\text{H}}^{\text{T}} = 25$) using the relationship of [Hasinger \(2008\)](#), while the green dashed lines represent those attained using the relationship of [Burlon et al. \(2011\)](#). The intercepts obtained in the quasar regime have been modified to match those obtained at lower luminosities.

values of f_{obs} for $\log L_X \gtrsim 45.8$, a flattening of the decline is expected below this luminosity. From Fig. 7 (left panel) in [Hasinger \(2008\)](#), one can see that above $\log L_X \approx 45$ the value of f_{obs} appears to be constant, similarly to what has been found in the hard X-ray band. A flattening of f_{obs} in the quasar regime would also be expected when considering the large amount of accreting material needed to power the AGN at these luminosities. It must, however, be stressed that the X-ray Baldwin effect has not been well studied at high luminosities, and that the *Chandra*/HEG sample we used in this luminosity range is small, not allowing us to reach a firm conclusion on the variation of the iron $K\alpha$ EW with the luminosity. Possible evidence of a flattening of the X-ray Baldwin effect in the quasar regime has recently been found by [Krumpe et al. \(2010\)](#).

4.2. The differences between X-ray and IR half-opening angle-luminosity relationships

The main difference between the different $\theta_{\text{OA}}-L_X$ relationships that we used is that those obtained in the X-rays are based on direct observations of the absorbing material in the line of sight, while the relationship of [Maiolino et al. \(2007\)](#) is extrapolated from the ratio of the thermal infrared emission to the primary AGN continuum. The $\theta_{\text{OA}}-L_{\lambda}(5100 \text{ \AA})$ relationship obtained by [Maiolino et al. \(2007\)](#) was converted into $\theta_{\text{OA}}-L_X$ using the $L_{2\text{keV}}-L_{\lambda}(2500 \text{ \AA})$ relation obtained by [Steffen et al. \(2006\)](#). It was then converted into the $L_X-L_{\lambda}(5100 \text{ \AA})$ relation assuming the optical-UV spectral slope obtained by [Vanden Berk et al. \(2001\)](#). All this is likely to introduce some error in the $\theta_{\text{OA}}-L_X$ obtained by their IR work. It has been argued by [Maiolino et al. \(2007\)](#) that the larger normalization of the $\theta_{\text{OA}}-L_X$ relation they found is related to the fact that medium X-ray surveys such as that of [Hasinger \(2008\)](#) are likely to miss a certain fraction

of heavily obscured objects, which can instead be detected in the IR. This is because most of the X-ray emission is depleted for Compton-thick AGN at energies $\lesssim 10$ keV. However, hard X-ray surveys such as that of [Burlon et al. \(2011\)](#), which are much less biased by absorption, have found a normalization of the $\theta_{\text{OA}}-L_X$ relation consistent with that obtained at lower X-ray energies (see Fig. 1). Our results also show that the X-ray Baldwin effect can be explained by a Compton-thin torus and that larger values of f_{obs} would imply even lower average values of N_{H}^{T} . Consequently missing heavily obscured objects would not significantly affect our results. X-rays are also probably better suited to probe the material responsible for the iron $K\alpha$ line emission. The line can, in fact, be emitted by both gas and dust, and while the IR can probe only the latter, X-rays are able to infer the amount of both gas and dust. We thus conclude that for the purpose of our study the X-ray $\theta_{\text{OA}}-L_X$ relations are better suited than those extrapolated from IR observations.

It could be argued that the value of the hydrogen column density commonly used to determine the fraction of obscured objects (and thus as an indicator of the presence of the torus) is $N_{\text{H}} \sim 10^{22} \text{ cm}^{-2}$, while the torus is believed to have larger values of N_{H}^{T} , and that this obscuration might be related to the presence of dust lanes or molecular structures in the host galaxy (e.g., [Matt 2000](#)). However, [Bianchi et al. \(2009\)](#), using the results of [Della Ceca et al. \(2008\)](#), have shown that a similar decrease of f_{obs} with the luminosity is obtained when this threshold is set to a larger value of line-of-sight column density. They found that the fraction of CT AGN decreases with the luminosity as $f_{\text{CT}} \propto L^{-0.22}$, similarly to what is found for f_{obs} in the X-rays (e.g., [Hasinger 2008](#), see Eq. (1)). A similar result was obtained by [Fiore et al. \(2009\)](#) from a *Chandra* and *Spitzer* study of AGN in the COSMOS field. This implies that the contribution of dust lanes or galactic molecular structures to the observed N_{H} does not significantly affect the $f_{\text{obs}}-L_X$ relationships obtained in the X-rays.

4.3. The equatorial column density of the torus

The distribution of values of the equatorial column density of the torus in AGN is still poorly constrained. X-ray observation can in fact infer solely the obscuration in the line of sight, and only studies of the reprocessed X-ray emission performed using physical torus models like those of [Ikeda et al. \(2009\)](#), [Brightman & Nandra \(2011\)](#) and [Murphy & Yaqoob \(2009\)](#) can help to deduce the value of the equatorial column density of the torus. However, these kinds of studies are still very scarce (e.g., [Rivers et al. 2011](#); [Brightman & Ueda 2012](#)), besides being largely geometry dependent. Studies of AGN in the mid-IR band performed using the clumpy torus formalism of [Nenkova et al. \(2008\)](#) have shown that the number of clouds along the equator is $N_0 \sim 5-10$ ([Mor et al. 2009](#)). If, as reported by [Mor et al. \(2009\)](#), each cloud of the torus has an optical depth of $\tau_V \sim 30-100$ (i.e., $\log N_{\text{H}}^{\text{T}} \sim 22-23$), the equatorial column density of the torus is expected to be $\log N_{\text{H}}^{\text{T}} \sim 22.5-24$, in agreement with our results. However, the value of the intercept of the X-ray Baldwin effect obtained by the simulations is strongly dependent on our assumptions on the metallicity of the torus. This implies that our constraints of the average value of N_{H}^{T} are also tightly related to the choice of the metallicity: lower values of the metallicity would lead to larger values of N_{H}^{T} . To study this effect, we repeated our study in the Seyfert regime using half-solar metallicities for the reflection model. We found that in this scenario one needs values of the equatorial column density of the

torus about two times larger than for the solar-metallicity case ($23.5 \lesssim \log N_{\text{H}}^{\text{T}} \lesssim 23.7$) to explain the X-ray Baldwin effect. To have a Compton-thick torus, one would then need values of the metallicity of $Z \lesssim 0.2 Z_{\odot}$. A Compton-thick torus has often been invoked to explain the Compton hump observed in the spectrum of many unobscured AGN (e.g., Bianchi et al. 2004). It is still unclear, however, which fraction of the Compton hump is produced in the distant reflector and which in the accretion flow. From our study, we have found that in unobscured objects the X-ray Baldwin effect can be explained by a luminosity-dependent covering factor of the torus for an average value of the equatorial column density of $\log N_{\text{H}}^{\text{T}} \approx 23.2$. This value is lower than the line-of-sight N_{H} of many Seyfert 2s (e.g., Ricci et al. 2011; Burlon et al. 2011) and might be related either to the geometry we adopted or to the presence of objects with sub-solar metallicities. In particular, due to the constant reflection angle relative to a local normal in any point of the reflecting surface, the spherical-toroidal geometry produces larger values of EW (and thus larger values of the intercept) with respect to a toroidal structure (Murphy & Yaqoob 2009).

It is possible that there exists a wide spread of equatorial column densities of the torus in AGN. Thus one could envisage that this, together with the different values of θ_i , would introduce the scatter observed in the anti-correlation. In Fig. 7 we show the X-ray Baldwin effect obtained from our simulations using the relationship of Hasinger (2008) for $\log N_{\text{H}}^{\text{T}} = 22.5$ and $\log N_{\text{H}}^{\text{T}} = 25$ in both the Seyfert and the quasar regime, together with the time-averaged *Chandra*/HEG data of Shu et al. (2010). From the figure it is evident that all the data are well within the range expected from our simulations, both in the Seyfert and in the quasar regime. It is not known whether there exists a relation between the equatorial column density of the torus and the AGN luminosity. However, we have shown that the variation of the covering factor of the torus with the luminosity alone can fully explain the observed trend, so that no additional luminosity-dependent physical parameter is needed.

4.4. Luminosity-dependent unification of AGN

The relation of f_{obs} with the luminosity might be connected to the increase of the inner radius of the torus with the luminosity due to dust sublimation. Both near-IR reverberation (Suganuma et al. 2006) and mid-IR interferometric (Tristram & Schartmann 2011) studies have confirmed that the inner radius of the molecular torus increases with the luminosity as $R_{\text{in}} \propto L^{0.5}$. Considering the geometry of Fig. 1, the fraction of obscured objects (see Eq. (4)) would be related to the ratio of the height to the inner radius of the torus (H/R) by

$$f_{\text{obs}} \approx \frac{H}{R} \sqrt{\frac{1}{1 + (H/R)^2}}. \quad (6)$$

For values of $H/R \lesssim 1$, $f_{\text{obs}} \propto H/R$. In the original formulation of the *receding torus* model (Lawrence 1991), the height H was considered to be constant. Assuming that R has the same luminosity dependence as R_{in} , this would lead to $f_{\text{obs}} \propto L^{-0.5}$. This has been shown to be inconsistent with recent observations, which point towards a flatter slope ($f_{\text{obs}} \propto L^{-0.25}$, e.g., Hasinger 2008, see Eq. (1)) and imply $H \propto L^{0.25}$. Hönig & Beckert (2007) have shown that in the frame of the *radiation-limited clumpy dust torus* model, one would obtain $H \propto L^{0.25}$, which is in agreement with the observations.

The decrease of the covering factor with luminosity might also have important implications on the AGN dichotomy. It has

been shown that Seyfert 2s appear to have on average lower luminosities and lower Eddington ratios (e.g., Beckmann et al. 2009; Ricci et al. 2011) than Seyfert 1s and Seyfert 1.5s, which suggests that they have on average a torus with a larger covering factor. This idea is also supported both by the fact that the luminosity function of type II AGN has been found to peak at lower luminosities than that of type I AGN (Della Ceca et al. 2008; Burlon et al. 2011) and by the results obtained by the recent mid-IR work of Ramos Almeida et al. (2011, see also Elitzur 2012). A difference in the torus covering factor distribution between different types of AGN would also explain why the average hard X-ray spectrum of Compton-thin Seyfert 2s shows a larger reflection component than that of Seyfert 1s and Seyfert 1.5s (Ricci et al. 2011). A similar result was found by Brightman & Ueda (2012): studying high-redshift AGN in the *Chandra* Deep Field South, they found that more obscured objects appear to have tori with larger covering factors, although they did not find a clear luminosity dependence.

5. Summary and conclusions

In this work we have studied the hypothesis that the X-ray Baldwin effect is related to the decrease of the covering factor of the torus with the luminosity. We used the physical torus models of Ikeda et al. (2009) and Brightman & Nandra (2011) to account for the reprocessed X-ray radiation, and the values of the fraction of obscured sources obtained by recent surveys in the X-rays and in the IR as a proxy of the covering factor of the torus. Our simulations show that the variation of the covering factor of the torus with the luminosity can explain the X-ray Baldwin effect. In the Seyfert regime ($L_{\text{X}} \leq 10^{44.2} \text{ erg s}^{-1}$), the observed $EW-L_{\text{X}}$ trend can be exactly (both in slope and intercept) reproduced by an average value of the equatorial column density of the torus of $\log N_{\text{H}}^{\text{T}} \approx 23.2$, while a slope consistent with the observations is obtained for a larger range of column densities ($\log N_{\text{H}}^{\text{T}} \gtrsim 23.1$). At higher luminosities ($L_{\text{X}} > 10^{44.2} \text{ erg s}^{-1}$), the situation is less clear due to the small number of high-quality observations available. Moreover, it is not clear whether f_{obs} still decreases with the luminosity in the quasar regime, similarly to what is found in the Seyfert regime (as shown by Hasinger 2008), or whether it is constant (as found by Burlon et al. 2011). A flattening of the $\theta_{\text{OA}}-L_{\text{X}}$ relationship could explain the observations at high luminosities (for $23.2 \lesssim \log N_{\text{H}}^{\text{T}} \lesssim 24.2$), and it might naturally arise from the large amount of accreting mass needed to power these luminous quasars.

In the next few years ASTRO-H (Takahashi et al. 2010), with its high energy-resolution calorimeter SXS, will allow to constrain the origin of the narrow component of the iron $K\alpha$ line, being able to separate the narrow core coming from the torus from the flux emitted closer to the central engine even better than *Chandra*/HEG. ASTRO-H will also be able to probe the narrow iron $K\alpha$ line in the quasar regime, which will allow us to understand the behavior of the X-ray Baldwin effect at high luminosities.

Acknowledgements. We thank the anonymous referee for his/her comments that helped improve the paper. We thank XinWen Shu and Rivay Mor for providing us with useful details about their work, and Chin Shin Chang and Poshak Gandhi for their comments on the manuscript. C.R. thanks the Sherpa group and IPAG for their hospitality during his stay in Grenoble. C.R. is a Fellow of the Japan Society for the Promotion of Science (JSPS).

References

- Arnaud, K. A. 1996, in *Astronomical Data Analysis Software and Systems V*, eds. G. H. Jacoby, & J. Barnes, ASP Conf. Ser., 101, 17
- Baldwin, J. A. 1977, *ApJ*, 214, 679
- Beckmann, V., Soldi, S., Ricci, C., et al. 2009, *A&A*, 505, 417
- Bianchi, S., Matt, G., Balestra, I., Guainazzi, M., & Perola, G. C. 2004, *A&A*, 422, 65
- Bianchi, S., Guainazzi, M., Matt, G., & Fonseca Bonilla, N. 2007, *A&A*, 467, L19
- Bianchi, S., La Franca, F., Matt, G., et al. 2008, *MNRAS*, 389, L52
- Bianchi, S., Bonilla, N. F., Guainazzi, M., Matt, G., & Ponti, G. 2009, *A&A*, 501, 915
- Brightman, M., & Nandra, K. 2011, *MNRAS*, 413, 1206
- Brightman, M., & Ueda, Y. 2012, *MNRAS*, 2850
- Burlon, D., Ajello, M., Greiner, J., et al. 2011, *ApJ*, 728, 58
- Croom, S. M., Rhoads, K., Corbett, E. A., et al. 2002, *MNRAS*, 337, 275
- de La Calle Pérez, I., Longinotti, A. L., Guainazzi, M., et al. 2010, *A&A*, 524, A50
- Della Cecca, R., Caccianiga, A., Severgnini, P., et al. 2008, *A&A*, 487, 119
- Dietrich, M., Hamann, F., Shields, J. C., et al. 2002, *ApJ*, 581, 912
- Dwelly, T., & Page, M. J. 2006, *MNRAS*, 372, 1755
- Elitzur, M. 2012, *ApJ*, 747, L33
- Elitzur, M., & Shlosman, I. 2006, *ApJ*, 648, L101
- Fabian, A. C., Iwasawa, K., Reynolds, C. S., & Young, A. J. 2000, *PASP*, 112, 1145
- Fiore, F., Puccetti, S., Brusa, M., et al. 2009, *ApJ*, 693, 447
- Fukazawa, Y., Hiragi, K., Mizuno, M., et al. 2011, *ApJ*, 727, 19
- Gandhi, P., Horst, H., Smette, A., et al. 2009, *A&A*, 502, 457
- Grandi, P., Malaguti, G., & Focchini, M. 2006, *ApJ*, 642, 113
- Green, P. J., Forster, K., & Kuraszewicz, J. 2001, *ApJ*, 556, 727
- Grimes, J. A., Rawlings, S., & Willott, C. J. 2004, *MNRAS*, 349, 503
- Hasinger, G. 2008, *A&A*, 490, 905
- Hönig, S. F., & Beckert, T. 2007, *MNRAS*, 380, 1172
- Hönig, S. F., Smette, A., Beckert, T., et al. 2008, *A&A*, 485, L21
- Ikeda, S., Awaki, H., & Terashima, Y. 2009, *ApJ*, 692, 608
- Iwasawa, K., & Taniguchi, Y. 1993, *ApJ*, 413, L15
- Jiang, P., Wang, J. X., & Wang, T. G. 2006, *ApJ*, 644, 725
- Jiménez-Bailón, E., Piconcelli, E., Guainazzi, M., et al. 2005, *A&A*, 435, 449
- Krolik, J. H., Madau, P., & Zycki, P. T. 1994, *ApJ*, 420, L57
- Krumpe, M., Lamer, G., Markowitz, A., & Corral, A. 2010, *ApJ*, 725, 2444
- Lawrence, A. 1991, *MNRAS*, 252, 586
- Lawrence, A., & Elvis, M. 1982, *ApJ*, 256, 410
- Lawrence, A., & Elvis, M. 2010, *ApJ*, 714, 561
- Magdziarz, P., & Zdziarski, A. A. 1995, *MNRAS*, 273, 837
- Maiolino, R., Shemmer, O., Imanishi, M., et al. 2007, *A&A*, 468, 979
- Matt, G. 2000, *A&A*, 355, L31
- Matt, G., Guainazzi, M., & Maiolino, R. 2003, *MNRAS*, 342, 422
- Miyakawa, T., Ebisawa, K., & Inoue, H. 2012, *PASJ*, 64, 140
- Mor, R., Netzer, H., & Elitzur, M. 2009, *ApJ*, 705, 298
- Murphy, K. D., & Yaqoob, T. 2009, *MNRAS*, 397, 1549
- Nandra, K. 2006, *MNRAS*, 368, L62
- Nandra, K., George, I. M., Mushotzky, R. F., Turner, T. J., & Yaqoob, T. 1997, *ApJ*, 488, L91
- Nandra, K., O'Neill, P. M., George, I. M., & Reeves, J. N. 2007, *MNRAS*, 382, 194
- Neenkova, M., Sirocky, M. M., Ivezić, Ž., & Elitzur, M. 2008, *ApJ*, 685, 147
- Netzer, H., Laor, A., & Gondhalekar, P. M. 1992, *MNRAS*, 254, 15
- Page, K. L., O'Brien, P. T., Reeves, J. N., & Turner, M. J. L. 2004, *MNRAS*, 347, 316
- Petrucci, P. O., Henri, G., Maraschi, L., et al. 2002, *A&A*, 388, L5
- Ramos Almeida, C., Levenson, N. A., Alonso-Herrero, A., et al. 2011, *ApJ*, 731, 92
- Reeves, J. N., & Turner, M. J. L. 2000, *MNRAS*, 316, 234
- Ricci, C., Walter, R., Courvoisier, T. J.-L., & Paltani, S. 2011, *A&A*, 532, A102
- Rivers, E., Markowitz, A., & Rothschild, R. 2011, *ApJ*, 732, 36
- Shu, X. W., Yaqoob, T., & Wang, J. X. 2010, *ApJS*, 187, 581
- Shu, X. W., Yaqoob, T., & Wang, J. X. 2011, *ApJ*, 738, 147
- Shu, X. W., Wang, J. X., Yaqoob, T., Jiang, P., & Zhou, Y. Y. 2012, *ApJ*, 744, L21
- Simpson, C. 2005, *MNRAS*, 360, 565
- Singh, V., Shastri, P., & Risaliti, G. 2011, *A&A*, 532, A84
- Steffen, A. T., Strateva, I., Brandt, W. N., et al. 2006, *AJ*, 131, 2826
- Suganuma, M., Yoshii, Y., Kobayashi, Y., et al. 2006, *ApJ*, 639, 46
- Takahashi, T., Mitsuda, K., Kelley, R., et al. 2010, in *SPIE Conf. Ser.*, 7732
- Treister, E., Krolik, J. H., & Dullemond, C. 2008, *ApJ*, 679, 140
- Tristram, K. R. W., & Schartmann, M. 2011, *A&A*, 531, A99
- Turner, T. J., & Miller, L. 2009, *A&ARv*, 17, 47
- Ueda, Y., Akiyama, M., Ohta, K., & Miyaji, T. 2003, *ApJ*, 598, 886
- Ueda, Y., Hiroi, K., Isobe, N., et al. 2011, *PASJ*, 63, 937
- Vanden Berk, D. E., Richards, G. T., Bauer, A., et al. 2001, *AJ*, 122, 549
- Yaqoob, T., & Padmanabhan, U. 2004, *ApJ*, 604, 63
- Zhou, X.-L., & Wang, J.-M. 2005, *ApJ*, 618, L83