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# Neon isotopic composition of the mantle constrained by single vesicle analyses

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## ABSTRACT

The origin of volatiles on Earth is still a matter of debate. Noble gases are an efficient geochemical tool to constrain Earth formation processes due to their inertness. Several studies have focused on the neon isotopic composition of the lower mantle because the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio is thought to reflect that of Earth's primordial components. Two models to explain the origin of light noble gases on Earth have been proposed: either solar wind implantation onto the Earth's solid precursors or dissolution into the mantle of a primordial atmosphere captured from solar nebula gas. In order to test these two models, we analyzed the noble gas compositions (helium, neon and argon) of two submarine oceanic island basalt glasses from Fernandina volcano (Galápagos archipelago), which have among the most primitive/unradiogenic terrestrial helium and neon isotopic compositions. Several sample pieces are studied both by step-crushing and by laser ablation analyses of single vesicles. Results of step-crushing are consistent with those of laser ablation analyses, but the latter results provide new insights into the origin of atmospheric contamination. The single-vesicle laser-ablation measurements overlap with the step crushing results, but have systematically higher  $^{40}\text{Ar}/^{36}\text{Ar}$ , and  $^3\text{He}/^{36}\text{Ar}$ , suggesting less atmospheric contamination using this method. The single vesicle data therefore suggest that atmospheric contamination is introduced by exposure to the modern atmosphere, after sample collection.  $^3\text{He}/^4\text{He}$  values are about 23 times the atmospheric ratio (R/Ra) for the two Fernandina (Galápagos) samples, in agreement with previous studies. We obtain  $^{20}\text{Ne}/^{22}\text{Ne}$  and  $^{40}\text{Ar}/^{36}\text{Ar}$  isotopic ratios as high as 12.91 and 9400, respectively, for the mantle source of the Galápagos hotspot. The new data show that step-crushing and laser ablation analyses are complementary methods that should be used together to derive the noble gas ratios in uncontaminated samples. The results of neon compositions are consistent with previous hotspot studies and support the model of solar wind implantation associated with sputtering to explain helium and neon origins on Earth.

## 1. Introduction

The origin of volatiles on Earth is still a matter of intense research, but is fundamental to understanding Earth and atmosphere formation processes. Noble gases are important tools to address this problem due to their chemical inertness. In particular, neon gives precious information on that question. Neon has three iso-

topes of masses 20, 21 and 22, which can all be produced by nucleogenic reactions (e.g.  $^{18}\text{O}(\alpha, n)^{21}\text{Ne}$ ). However, the production rates of isotopes  $^{20}\text{Ne}$  and  $^{22}\text{Ne}$  are negligible in the terrestrial mantle (Yatsevich and Honda, 1997), so the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of the mantle reflects the primitive neon composition. Many studies have put forward that the mantle  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio is higher than 12 and is therefore “solar-like” (Ballentine et al., 2005; Honda et al., 1993; Kurz et al., 2009; Moreira et al., 1998; Mukhopadhyay, 2012; Raquin and Moreira, 2009; Sarda et al., 1988; Trielloff et al., 2000). By comparison, the atmospheric  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio is  $9.78 \pm 0.03$  (Sano et al., 2013) and the solar wind ratio, measured on the Gen-

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esis targets, is  $13.78 \pm 0.03$  ( $1\sigma$ ) according to Heber et al. (2009) or  $14.001 \pm 0.042$  ( $1\sigma$ ) according to Pepin et al. (2012). The solar wind is isotopically fractionated compared to the Sun, which, by calculation, has a  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of 13.34 using an inefficient Coulomb drag model (Heber et al., 2012).

It is essential to precisely measure the neon composition of the lower mantle. Indeed, the lower mantle must reflect a less degassed reservoir than the upper mantle and still contains primordial noble gas components, in particular primordial helium ( $^3\text{He}$ ) and neon (Allègre et al., 1983; Kaneoka and Takaoka, 1980; Kurz et al., 2009; Valbracht et al., 1997). In contrast, the upper mantle is considered to be more degassed of its primordial noble gases and should be more sensitive to subduction of atmospheric noble gases. Only a few hotspots are well-suited to determine the lower mantle noble gas composition, namely Hawaii, Galápagos and Iceland, because fresh submarine glass samples, erupted at great depths, are required to obtain accurate neon analyses. Submarine glasses are among the best samples for recording the quenched magma (mantle) noble gas composition directly after eruption.

Several studies have focused on the lower mantle composition and have led to two opposing models for the origin of noble gases on Earth. Trielloff et al. (2000) found that oceanic island basalts (OIB) from Hawaii and Iceland showed a mean  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of  $12.49 \pm 0.06$  ( $1\sigma$ ), very close to that of the Neon B component observed in gas-rich meteorites and on the lunar soil (Black, 1972). They argued that since there were no extant terrestrial measurements higher than  $\sim 12.5$ , Neon B was the likely terrestrial end-member, assumed to be derived from solar wind implantation (Black, 1972). Hence, Trielloff et al. (2000) suggested that planetesimals that accreted to form the Earth had suffered solar wind implantation and were then characterized by the Neon B component. However, Yokochi and Marty (2004) measured a  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of  $13.04 \pm 0.2$  ( $1\sigma$ ) in plume-related rocks from the Kola peninsula (Russia), which led them to propose that this ratio of 13 would represent a lower limit for the lower mantle neon composition and that terrestrial neon would come from the dissolution of solar nebula gas into the mantle. The latter scenario, also advocated by Mukhopadhyay (2012), implies that the solar nebula gas remained in the accretion disk long enough to permit the capture of a primordial atmosphere, an important temporal constraint for the early Earth. Raquin and Moreira (2009) developed a model that would explain implanted  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio, starting from solar wind values of  $\sim 13.8$ , by isotopic fractionation during implantation and erosion on the grain surfaces. In this model, lighter isotopes are preferentially lost from accreting and eroding grains due to their lower energy, shorter implantation depth, and preferential loss by erosion. Raquin and Moreira (2009) were able to derive a steady state  $^{20}\text{Ne}/^{22}\text{Ne}$  of  $\sim 12.5$  for accreting particles, using reasonable values of exposure time and erosion rates. However, the timing and geometries (particle sizes) of accretion are poorly known, so higher steady state values (such as 12.7) are possible (Moreira, 2013). In all accretion models, whether gases are introduced by implantation on solid grains or by equilibration with nebular gas, the neon isotopic composition of the Earth's interior is one of very few constraints.

One of the major problems in noble gas geochemistry is the ubiquity of an air-like noble gas component in mantle-derived samples. Step-crushing experiments always consist of mixing between an air-like component and a mantle component as shown by many studies (e.g., Trielloff et al., 2000 and Yokochi and Marty, 2004). Ballentine and Barfod (2000) suggested that this air-like component is due to atmospheric contamination such that air fills the many cracks and microfractures within samples once they are brought to the surface. Conversely, Sarda (2004) suggested that this air-like component could correspond to atmospheric noble gas re-

cycling into the mantle and that the larger vesicles would carry this recycled component. Several studies (e.g., Holland and Ballentine, 2006, Parai and Mukhopadhyay, 2015 and Tucker et al., 2012) have now suggested that both processes (shallow level contamination and recycling through subduction) may occur for heavy noble gases (Ar, Kr, Xe) whereas light noble gases (He, Ne) would not be significantly recycled into the mantle (Holland and Ballentine, 2006). Burnard et al. (1997) and Burnard (1999) introduced laser ablation analyses of single vesicles and never found vesicles filled with air. Raquin et al. (2008) showed that laser ablation analyses can (partly) remove this shallow level atmospheric contaminant and therefore is a promising method to determine the true mantle isotopic composition.

In this study, we focus on OIB samples from one of the most primitive hotspots for helium, the Galápagos hotspot (Kurz et al., 2009; Kurz and Geist, 1999) in order to determine their helium, neon and argon compositions and to constrain which of the two models for the origin of noble gases on Earth is best supported by the data. In addition to step-crushing experiments, we conduct laser ablation analyses on single vesicles. Since OIB samples are not very vesicular, pieces of samples for laser ablation are characterized with X-ray microtomography, a powerful and non-destructive technique, which allows precise location of vesicles in samples.

## 2. Samples and methods

### 2.1. Sample selection

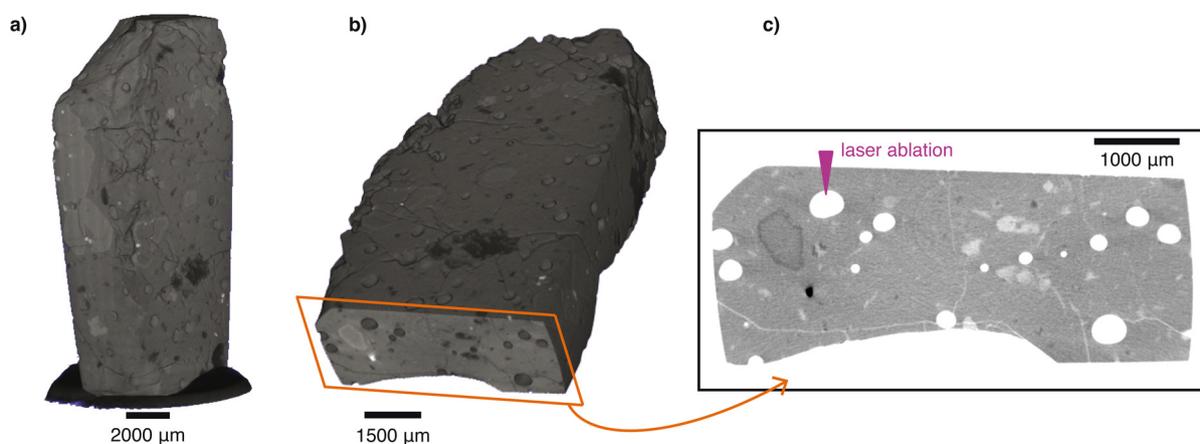
Two samples from Fernandina volcano (Galápagos hotspot) collected during the AHA-NEMO2 cruise are studied, samples AHA-NEMO2-D22A and AHA-NEMO2-D22B from a single dredge on the submarine western flank. Details about sample location can be found in Geist et al. (2006); helium and neon data for D22B and a number of other Fernandina samples can be found in Kurz et al. (2009).

These two samples were selected for their thick basaltic glassy pillow lava margins, which allowed the use of intact large (centimeter sized) glass chunks for the laser ablation and step crushing measurements. They are assumed to have recorded the quenched magma composition directly after eruption, without having suffered intense post-eruptive processes (such as slow cooling or crystallization). For such a study, it is best to choose samples collected at great water depth so that vesicle internal pressures are high when the magma is quenched (approximately equal to hydrostatic pressure (Colin et al., 2013)), increasing the likelihood of gas-rich vesicles. The two studied samples were dredged between 2210 and 2390 m deep.

### 2.2. X-ray microtomography

Before laser ablation analyses, pieces of samples were imaged via X-ray microtomography in order to locate the vesicles. X-ray computed microtomography is a powerful and non-destructive technique, which permits the reconstruction of the 3D volume of geological samples. The analyses were performed at the Institut des Sciences de la Terre at Orléans and at the Institute of Earth Sciences, University of Lausanne.

Sample pieces are first polished so as to make rectangular shapes of 5–8 mm high, 3–5 mm long and 2–5 mm wide. This ensures easy focusing of the laser beam. Details about X-ray microtomography principles and acquisition parameters can be found in Supplementary Information. Using the X-ray microtomography data, the internal 3D volume, and surfaces, of the sample were reconstructed with different softwares (CTvox, ImageJ, etc.) in order to locate vesicles. Fig. 1a shows one piece of the sample



**Fig. 1.** Principle for locating subsurface vesicles of sample AHA-NEMO2-D22B before laser ablation: a 3D volume (a) is obtained from all the slices. A virtual cut through the object is shown in (b) and the slice (c) represents one transverse section of the sample as shown in (b). On the slice (c), vesicles are in white and the glass is in gray, as well as a few crystals. Some cracks are also visible and link some vesicles to air (as the one in the lower right corner). On the slice (c) we can measure the vesicle depths before piercing them with the laser.

AHA-NEMO2-D22B obtained with Bruker's software CTvox. As explained in Supplementary Information, each slice corresponds to a transverse section of the sample; this is shown in Fig. 1b and Fig. 1c. In this way, we can identify vesicle positions and measure their diameters and depths before piercing them with the laser. This preliminary work is of huge importance since these OIB samples are not heavily vesiculated and also allows evaluation whether vesicles are intact or connected to the surface with cracks.

Sample vesicularities were estimated with ImageJ software (details about calculations are explained in Supplementary Information). Estimated vesicularities are low: on the order of 3–4% for the two Galápagos samples (refer to Table S2 in Supplementary Information). These low vesicularities demonstrate the importance of X-ray microtomography for single vesicle analyses.

### 2.3. Step-crushing and laser ablation analyses

All pieces of samples were cleaned in successive ethanol and acetone baths before the analyses. If some weathering marks were still detected with the binocular microscope, samples were further cleansed with oxalic acid (1%).

Pieces of samples for step-crushing analyses were loaded into crushers, which were baked under ultra-high vacuum between 100–150 °C for at least one night before starting the experiments in order to remove atmospheric noble gases adsorbed onto the sample surfaces and the crusher walls. The sample AHA-NEMO2-D22B was analyzed by step-crushing both at IPGP and at Woods Hole Oceanographic Institution (WHOI). The analyses at WHOI were conducted following the same procedure described by Kurz et al. (2009), but involved a much larger sample size and included argon measurements for each step. There is excellent overall agreement between the step crushing results in the two laboratories, as discussed further below.

For laser ablation analyses, pieces of samples were set in a laser cell which was also baked at 100 °C for at least two days. Since vesicles have been located via X-ray microtomography, we could pierce them with the laser. The laser used was an argon fluoride excimer laser of wavelength 193 nm (ATLEX-300i system). The laser was run at a frequency of 100 Hz with energy of 1.2 to 2 mJ per pulse of 3–7 ns. The laser beam reached a maximum depth of about 100–150 µm in 5–6 min with energy of 2 mJ per pulse. To pierce a vesicle, a one to six minutes laser shot was made. The pressure in the laser cell was measured with a MKS Baratron® manometer. If a vesicle was reached, there was a sudden pressure jump of 0.0050–0.0650 Torr depending on vesicle size. If no vesicle was pierced, the laser cell was pumped for one minute. Several

cycles of ablation/pumping were needed to reach some vesicles located more than 100–150 µm deep. In order to derive the quantity of CO<sub>2</sub> in the vesicles, the manometer was calibrated with an air standard. We then assumed that the major gas in the vesicles is CO<sub>2</sub> (Moore et al., 1977).

Once gases were released, either by crushing or by laser ablation, they were purified with a Ti-getter (which was first at 1073 K and then at ambient temperature) and a SAES getter before being trapped onto activated charcoal at about 14 K. Noble gases were then analyzed one after the other with the Noblesse multi-collector mass spectrometer (Nu instruments ©). Neon abundances and isotopic ratios have to be corrected *a posteriori* for interferences between <sup>20</sup>Ne and HF<sup>+</sup>, <sup>40</sup>Ar<sup>++</sup> and also between <sup>22</sup>Ne and CO<sub>2</sub><sup>++</sup>. To do that, a scan at mass 20, mainly composed of HF, was made before introducing the neon into the mass spectrometer, and <sup>40</sup>Ar<sup>+</sup> as well as CO<sub>2</sub><sup>+</sup> were also measured during neon cycles. Since the double-to-single charge ratios <sup>40</sup>Ar<sup>++</sup>/<sup>40</sup>Ar<sup>+</sup> and CO<sub>2</sub><sup>++</sup>/CO<sub>2</sub><sup>+</sup> were determined on the Noblesse mass spectrometer during the analyses (0.075 and 0.01 respectively), neon isotopic ratios can be derived.

Blanks are of huge importance, especially for laser ablation analyses because the gas quantity in a single vesicle is relatively small. Blank analyses were performed in the same way as for sample analyses. For example, if the laser shot to reach a vesicle lasted five minutes, we pierced the sample matrix during five minutes and then corrected the vesicle analysis with the corresponding matrix blank. The matrix blank compositions are indicated in Tables S3a and S3b in Supplementary Information. We started laser ablation as soon as blanks of the line (only with the laser cell and without the crushers) were about  $2.8 \times 10^{-10}$  cm<sup>3</sup> of <sup>4</sup>He, less than  $2.5 \times 10^{-13}$  cm<sup>3</sup> of <sup>22</sup>Ne and less than  $2.2 \times 10^{-12}$  cm<sup>3</sup> of <sup>36</sup>Ar. For step-crushing experiments, typical blanks were about  $5.0 \times 10^{-11}$  cm<sup>3</sup> of <sup>4</sup>He,  $1.3 \times 10^{-13}$  cm<sup>3</sup> of <sup>22</sup>Ne and less than  $4.0 \times 10^{-12}$  cm<sup>3</sup> of <sup>36</sup>Ar. Once the experiments were started, blanks were measured every day; they remained stable and even tend to decrease.

Two standards are used in the present study: an air standard using a pipette connected to a 1 l reservoir in which 0.410 cm<sup>3</sup> of air was expanded, and a homemade standard at the IPGP laboratory dedicated to He and Ne analyses in He–Ne poor samples. The latter was obtained by crushing 120 g of the MORB sample MD57 D9 from Indian Ocean, 7.947 g of sample PND6 from Pitcairn hotspot and 64.15 g of the MORB sample EW9309 33Dsg from South Atlantic. The standard was calibrated against air and the isotopic ratios are <sup>3</sup>He/<sup>4</sup>He = 7.73 Ra, <sup>20</sup>Ne/<sup>22</sup>Ne = 9.94,

**Table 1**  
Abundances of  $^4\text{He}$ ,  $^{22}\text{Ne}$  and  $^{36}\text{Ar}$  for the two Galápagos samples analyzed by step-crushing. The new results from WHOI are indicated. For sample AHA-NEMO2-D22A, 9 steps have been carried out and 11 for sample AHA-NEMO2-D22B. For each step, the number of strokes is indicated. Errors are 1 sigma uncertainties. n.a. stands for not analyzed.

Sample	Crush	Weigh (g)	$^4\text{He} \times 10^{-7}$ (cm <sup>3</sup> /g)	$^{22}\text{Ne} \times 10^{-12}$ (cm <sup>3</sup> /g)	$^{36}\text{Ar} \times 10^{-11}$ (cm <sup>3</sup> /g)
AHA-NEMO2-D22A	1x	0.577	3.10 ± 0.09	8.74 ± 0.26	13.6 ± 0.2
AHA-NEMO2-D22A	1x	0.577	4.02 ± 0.12	17.2 ± 0.52	7.61 ± 0.13
AHA-NEMO2-D22A	2x	0.577	6.44 ± 0.19	39.0 ± 1.2	n.a.
AHA-NEMO2-D22A	2x	0.577	2.82 ± 0.08	5.38 ± 0.16	1.51 ± 0.03
AHA-NEMO2-D22A	6x	0.577	3.33 ± 0.10	7.55 ± 0.23	21.2 ± 0.4
AHA-NEMO2-D22A	5x	0.577	2.64 ± 0.08	8.64 ± 0.26	12.7 ± 0.2
AHA-NEMO2-D22A	7x	0.577	2.40 ± 0.07	22.7 ± 0.68	58.5 ± 1.0
AHA-NEMO2-D22A	14x	0.577	1.41 ± 0.04	3.80 ± 0.12	2.67 ± 0.05
AHA-NEMO2-D22A	40x	0.577	1.87 ± 0.06	3.37 ± 0.11	0.95 ± 0.02
AHA-NEMO2-D22B	1x	1.1376	0.40 ± 0.01	1.20 ± 0.04	0.91 ± 0.03
AHA-NEMO2-D22B	5x	1.1376	4.46 ± 0.13	230.9 ± 6.9	482.3 ± 14.5
AHA-NEMO2-D22B	4x	1.1376	1.59 ± 0.05	5.17 ± 0.16	4.89 ± 0.15
AHA-NEMO2-D22B	4x	1.1376	1.63 ± 0.05	3.93 ± 0.12	2.98 ± 0.09
AHA-NEMO2-D22B	5x	1.1376	2.79 ± 0.08	28.2 ± 0.8	43.7 ± 1.31
AHA-NEMO2-D22B	2x	1.1376	1.60 ± 0.05	4.73 ± 0.14	5.22 ± 0.16
AHA-NEMO2-D22B	3x	1.1376	1.49 ± 0.04	5.16 ± 0.16	5.64 ± 0.17
AHA-NEMO2-D22B	6x	1.1376	1.54 ± 0.05	4.47 ± 0.14	6.25 ± 0.19
AHA-NEMO2-D22B	2x	1.1376	1.93 ± 0.06	3.48 ± 0.11	1.03 ± 0.03
AHA-NEMO2-D22B	6x	1.1376	1.28 ± 0.04	3.45 ± 0.11	1.04 ± 0.03
AHA-NEMO2-D22B	6x	1.1376	2.18 ± 0.07	6.21 ± 0.19	2.31 ± 0.07
AHA-NEMO2-D22B WHOI	1x	3.8756	0.38	5.31	9.29
AHA-NEMO2-D22B WHOI	1x	3.8756	0.53	4.07	5.78
AHA-NEMO2-D22B WHOI	1x	3.8756	0.41	1.00	0.46
AHA-NEMO2-D22B WHOI	1x	3.8756	1.31	5.65	5.29
AHA-NEMO2-D22B WHOI	1x	3.8756	0.43	1.01	0.48
AHA-NEMO2-D22B WHOI	1x	3.8756	0.25	1.61	1.47
AHA-NEMO2-D22B WHOI	2x	3.8756	0.60	1.40	0.73
AHA-NEMO2-D22B WHOI	2x	3.8756	0.30	5.00	7.39
AHA-NEMO2-D22B WHOI	3x	3.8756	0.42	15.3	28.8
AHA-NEMO2-D22B WHOI	3x	3.8756	1.38	27.1	47.2
AHA-NEMO2-D22B WHOI	3x	3.8756	0.93	2.82	1.90
AHA-NEMO2-D22B WHOI	3x	3.8756	0.72	27.4	45.7
AHA-NEMO2-D22B WHOI	3x	3.8756	0.35	0.80	0.27
AHA-NEMO2-D22B WHOI	4x	3.8756	1.10	3.85	3.14
AHA-NEMO2-D22B WHOI	4x	3.8756	1.03	8.74	8.61
AHA-NEMO2-D22B WHOI	4x	3.8756	0.50	4.06	5.15
AHA-NEMO2-D22B WHOI	4x	3.8756	0.74	1.87	2.27
AHA-NEMO2-D22B WHOI	4x	3.8756	0.81	1.92	1.11
AHA-NEMO2-D22B WHOI	4x	3.8756	0.57	1.21	0.32
AHA-NEMO2-D22B WHOI	4x	3.8756	0.92	4.21	4.04
AHA-NEMO2-D22B WHOI	4x	3.8756	0.36	2.85	3.87
AHA-NEMO2-D22B WHOI	5x	3.8756	0.34	1.07	0.76
AHA-NEMO2-D22B WHOI	5x	3.8756	0.51	1.41	0.88
AHA-NEMO2-D22B WHOI	6x	3.8756	1.08	2.42	0.64

$^{21}\text{Ne}/^{22}\text{Ne} = 0.0309$ . Each pipette of the standard corresponds to  $9.63 \times 10^{-8}$  cm<sup>3</sup> STP of  $^4\text{He}$  and  $3.56 \times 10^{-12}$  cm<sup>3</sup> STP of  $^{22}\text{Ne}$ . This standard was used for He and Ne measurements in vesicles and He analyses in crushing steps, whereas the air standard was used for Ne and Ar isotopic measurements in crushing steps.

### 3. Results

Noble gas abundances and isotopic compositions for step-crushing and laser ablation analyses are given in Tables 1 to 4.

#### 3.1. Helium and radiogenic argon

Helium measurements are presented relative to atmosphere (Table 2 and Table 4), where the sample  $^3\text{He}/^4\text{He}$  isotopic ratio (R) is reported relative to the atmospheric ratio (Ra) of  $1.384 \times 10^{-6}$  (Clarke et al., 1976). As shown in Table 4, vesicles of the same sample have homogeneous helium isotopic compositions, and the results of step-crushing (Table 2) are consistent with the laser ablation results for  $^3\text{He}/^4\text{He}$ . The sample AHA-NEMO2-D22A has a mean  $^3\text{He}/^4\text{He}$  of  $23.4 \pm 0.3$  Ra ( $1\sigma$ ) with step-crushing analyses

and a mean of  $22.6 \pm 0.2$  Ra ( $1\sigma$ ) with vesicles analyses. The sample AHA-NEMO2-D22B has a mean  $^3\text{He}/^4\text{He}$  between  $23.0 \pm 0.3$  Ra ( $1\sigma$ ) (WHOI) and  $23.2 \pm 0.4$  Ra ( $1\sigma$ ) (IPGP) with step-crushing analyses and of  $23.0 \pm 0.5$  Ra ( $1\sigma$ ) with laser ablation results. The mean  $^3\text{He}/^4\text{He}$  values are lower for vesicle analyses but are nearly the same within uncertainties as those derived through step-crushing analyses. Such values are close to those found by Kurz et al. (2009) ( $23.07 \pm 0.10$  for sample AHA-NEMO2-D22A and  $22.45 \pm 0.10$  for sample AHA-NEMO2-D22B).

The  $^4\text{He}/^{40}\text{Ar}^*$  ratio is a proxy of degassing processes ( $^{40}\text{Ar}^*$  is radiogenic  $^{40}\text{Ar}$  and is calculated by subtracting the atmosphere composition to the measured  $^{40}\text{Ar}$ ). Indeed,  $^4\text{He}$  and  $^{40}\text{Ar}^*$  are two radiogenic isotopes produced by the decay of  $^{235}\text{U}$ ,  $^{238}\text{U}$  and  $^{232}\text{Th}$  for helium and by the decay of  $^{40}\text{K}$  for argon, and the (U + Th)/K ratio is relatively homogeneous in the lower mantle (Arevalo et al., 2009). The  $^4\text{He}/^{40}\text{Ar}^*$  production ratio varies from 2 to 5 (Moreira and Kurz, 2013). Since helium is nine times more soluble in a tholeiitic melt than argon (Jambon et al., 1986), helium and argon would be fractionated if a magma was degassed during its ascent to the surface, leading to a higher  $^4\text{He}/^{40}\text{Ar}^*$  ratio than the production ratio. The two Galápagos samples have  $^4\text{He}/^{40}\text{Ar}^*$  ra-

**Table 2**

Isotopic compositions of the two Galápagos samples analyzed by step-crushing. The new results from WHOI are indicated. Ra is the air  $^3\text{He}/^4\text{He}$  ratio equals to  $1.384 \times 10^{-6}$ .  $^{40}\text{Ar}^*$  is radiogenic  $^{40}\text{Ar}$ . Errors are 1 sigma uncertainties. n.a. stands for not analyzed and n.d. for not determined. We consider a  $^{40}\text{Ar}/^{36}\text{Ar}$  ratio of 298.6 for the air as obtained by Lee et al. (2006) to determine the  $^{40}\text{Ar}/^{36}\text{Ar}$  and  $^4\text{He}/^{40}\text{Ar}^*$  ratios.

Sample	Crush	$^3\text{He}/^4\text{He}$ ( $\times$ Ra)	$^{20}\text{Ne}/^{22}\text{Ne}$	$^{21}\text{Ne}/^{22}\text{Ne}$	$^{40}\text{Ar}/^{36}\text{Ar}$	$^4\text{He}/^{40}\text{Ar}^*$
D22A	1x	23.57 $\pm$ 0.15	11.74 $\pm$ 0.06	0.0322 $\pm$ 0.0004	1077 $\pm$ 2	2.92 $\pm$ 0.10
D22A	1x	23.12 $\pm$ 0.13	11.18 $\pm$ 0.05	0.0312 $\pm$ 0.0002	2190 $\pm$ 5	2.79 $\pm$ 0.10
D22A	2x	23.56 $\pm$ 0.14	10.73 $\pm$ 0.05	0.0306 $\pm$ 0.0002	n.a.	n.a.
D22A	2x	23.45 $\pm$ 0.16	12.68 $\pm$ 0.08	0.0332 $\pm$ 0.0003	6638 $\pm$ 56	2.94 $\pm$ 0.11
D22A	6x	23.67 $\pm$ 0.15	12.18 $\pm$ 0.07	0.0330 $\pm$ 0.0003	824 $\pm$ 2	2.99 $\pm$ 0.10
D22A	5x	23.20 $\pm$ 0.13	11.61 $\pm$ 0.06	0.0318 $\pm$ 0.0004	975 $\pm$ 2	3.08 $\pm$ 0.11
D22A	7x	23.06 $\pm$ 0.16	10.42 $\pm$ 0.04	0.0300 $\pm$ 0.0002	437 $\pm$ 1	2.96 $\pm$ 0.10
D22A	14x	23.67 $\pm$ 0.18	11.87 $\pm$ 0.10	0.0316 $\pm$ 0.0006	2072 $\pm$ 10	2.99 $\pm$ 0.11
D22A	40x	23.08 $\pm$ 0.13	12.72 $\pm$ 0.11	0.0341 $\pm$ 0.0005	6702 $\pm$ 76	3.08 $\pm$ 0.12
D22B	1x	22.72 $\pm$ 0.14	11.80 $\pm$ 0.10	0.0314 $\pm$ 0.0008	1777 $\pm$ 24	3.00 $\pm$ 0.14
D22B	5x	23.44 $\pm$ 0.15	9.95 $\pm$ 0.04	0.0291 $\pm$ 0.0001	335 $\pm$ 1	2.52 $\pm$ 0.16
D22B	4x	22.12 $\pm$ 0.14	11.85 $\pm$ 0.05	0.0320 $\pm$ 0.0004	1544 $\pm$ 4	2.62 $\pm$ 0.11
D22B	4x	23.47 $\pm$ 0.15	12.13 $\pm$ 0.06	0.0324 $\pm$ 0.0004	2147 $\pm$ 8	2.96 $\pm$ 0.13
D22B	5x	23.35 $\pm$ 0.14	10.35 $\pm$ 0.04	0.0298 $\pm$ 0.0001	516 $\pm$ 1	2.94 $\pm$ 0.13
D22B	2x	23.40 $\pm$ 0.15	11.82 $\pm$ 0.05	0.0319 $\pm$ 0.0003	1332 $\pm$ 3	2.96 $\pm$ 0.13
D22B	3x	23.42 $\pm$ 0.15	11.41 $\pm$ 0.05	0.0314 $\pm$ 0.0003	1184 $\pm$ 3	2.98 $\pm$ 0.13
D22B	6x	23.51 $\pm$ 0.15	11.78 $\pm$ 0.05	0.0322 $\pm$ 0.0003	1155 $\pm$ 3	2.87 $\pm$ 0.12
D22B	2x	22.99 $\pm$ 0.12	12.91 $\pm$ 0.07	0.0336 $\pm$ 0.0003	6605 $\pm$ 57	2.96 $\pm$ 0.13
D22B	6x	22.85 $\pm$ 0.13	12.02 $\pm$ 0.06	0.0328 $\pm$ 0.0003	4541 $\pm$ 32	2.89 $\pm$ 0.13
D22B	6x	23.48 $\pm$ 0.14	11.92 $\pm$ 0.06	0.0323 $\pm$ 0.0003	3517 $\pm$ 13	2.92 $\pm$ 0.13
D22B WHOI	1x	23.33 $\pm$ 0.57	10.24 $\pm$ 0.04	0.0295 $\pm$ 0.0003	436 $\pm$ 6	2.87
D22B WHOI	1x	23.73 $\pm$ 0.58	10.70 $\pm$ 0.04	0.0302 $\pm$ 0.0004	601 $\pm$ 4	2.98
D22B WHOI	1x	23.43 $\pm$ 0.53	12.50 $\pm$ 0.05	0.0338 $\pm$ 0.0005	3408 $\pm$ 79	2.93
D22B WHOI	1x	23.33 $\pm$ 0.51	11.49 $\pm$ 0.04	0.0317 $\pm$ 0.0003	1165 $\pm$ 30	2.85
D22B WHOI	1x	23.05 $\pm$ 0.58	12.50 $\pm$ 0.05	0.0331 $\pm$ 0.0005	3406 $\pm$ 70	2.87
D22B WHOI	1x	23.09 $\pm$ 0.57	10.98 $\pm$ 0.04	0.0312 $\pm$ 0.0004	916 $\pm$ 11	2.78
D22B WHOI	2x	22.25 $\pm$ 0.49	12.49 $\pm$ 0.05	0.0338 $\pm$ 0.0005	3021 $\pm$ 59	2.98
D22B WHOI	2x	23.06 $\pm$ 0.49	10.28 $\pm$ 0.04	0.0295 $\pm$ 0.0003	437 $\pm$ 3	2.84
D22B WHOI	3x	22.57 $\pm$ 0.53	9.97 $\pm$ 0.04	0.0289 $\pm$ 0.0003	348 $\pm$ 3	2.79
D22B WHOI	3x	22.85 $\pm$ 0.47	10.19 $\pm$ 0.04	0.0293 $\pm$ 0.0003	396 $\pm$ 2	2.91
D22B WHOI	3x	22.86 $\pm$ 0.38	12.13 $\pm$ 0.05	0.0327 $\pm$ 0.0004	2031 $\pm$ 22	2.81
D22B WHOI	3x	22.65 $\pm$ 0.37	10.01 $\pm$ 0.04	0.0289 $\pm$ 0.0003	349 $\pm$ 2	2.95
D22B WHOI	3x	23.53 $\pm$ 0.51	12.59 $\pm$ 0.05	0.0338 $\pm$ 0.0005	4722 $\pm$ 281	2.91
D22B WHOI	4x	23.07 $\pm$ 0.37	11.82 $\pm$ 0.05	0.0322 $\pm$ 0.0004	1517 $\pm$ 165	2.86
D22B WHOI	4x	23.05 $\pm$ 0.39	10.81 $\pm$ 0.04	0.0303 $\pm$ 0.0003	704 $\pm$ 9	2.93
D22B WHOI	4x	23.01 $\pm$ 0.40	10.72 $\pm$ 0.04	0.0306 $\pm$ 0.0004	636 $\pm$ 4	2.83
D22B WHOI	4x	22.92 $\pm$ 0.37	12.34 $\pm$ 0.05	0.0330 $\pm$ 0.0004	1362 $\pm$ 13	3.03
D22B WHOI	4x	22.68 $\pm$ 0.44	12.58 $\pm$ 0.05	0.0336 $\pm$ 0.0004	2706 $\pm$ 36	3.02
D22B WHOI	4x	23.03 $\pm$ 0.41	12.79 $\pm$ 0.05	0.0341 $\pm$ 0.0005	6147 $\pm$ 137	3.01
D22B WHOI	4x	22.84 $\pm$ 0.32	11.33 $\pm$ 0.04	0.0313 $\pm$ 0.0003	1095 $\pm$ 24	2.85
D22B WHOI	4x	23.33 $\pm$ 0.34	10.64 $\pm$ 0.04	0.0304 $\pm$ 0.0004	620 $\pm$ 4	2.88
D22B WHOI	5x	23.42 $\pm$ 0.40	11.91 $\pm$ 0.05	0.0328 $\pm$ 0.0005	1865 $\pm$ 93	2.83
D22B WHOI	5x	22.59 $\pm$ 0.54	12.16 $\pm$ 0.05	0.0331 $\pm$ 0.0004	2239 $\pm$ 42	2.99
D22B WHOI	6x	23.07 $\pm$ 0.34	12.85 $\pm$ 0.05	0.0332 $\pm$ 0.0004	6026 $\pm$ 127	2.97

**Table 3**

$\text{CO}_2$  and noble gas abundances of the pierced vesicles for the two Galápagos samples. Volumes of vesicles are calculated with the 3D Object Counter plugin of ImageJ. The  $\text{CO}_2$  contents are calculated thanks to the pressure jump recorded on the manometer and calibrated with the air standard (refer to section 2.3). Errors are 1 sigma uncertainties. n.d. means not determined when gases from the vesicle are indistinguishable from the matrix blank.

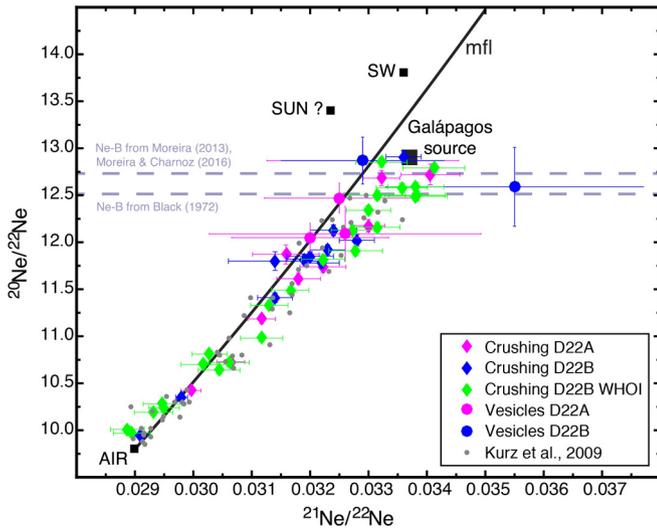
Sample	Volume $\times 10^6$ ( $\mu\text{m}^3$ )	$\text{CO}_2 \times 10^{-4}$ ( $\text{cm}^3$ )	$^4\text{He} \times 10^{-9}$ ( $\text{cm}^3$ )	$^{22}\text{Ne} \times 10^{-13}$ ( $\text{cm}^3$ )	$^{36}\text{Ar} \times 10^{-13}$ ( $\text{cm}^3$ )
<b>AHA-NEMO2-D22A vesicles</b>					
V1	5.0	4.78 $\pm$ 0.23	6.72 $\pm$ 0.21	1.49 $\pm$ 0.18	2.45 $\pm$ 1.33
V2	18.2	15.6 $\pm$ 0.7	17.9 $\pm$ 0.5	3.97 $\pm$ 0.21	8.23 $\pm$ 1.35
V3	16.7	13.8 $\pm$ 0.6	18.8 $\pm$ 0.6	4.72 $\pm$ 0.23	11.5 $\pm$ 1.4
V4	16.9	15.8 $\pm$ 0.7	21.1 $\pm$ 0.6	4.18 $\pm$ 0.22	12.0 $\pm$ 1.4
<b>AHA-NEMO2-D22B vesicles</b>					
V1	4.0	3.27 $\pm$ 0.16	4.76 $\pm$ 0.15	1.29 $\pm$ 0.18	2.64 $\pm$ 1.33
V2	2.9	2.06 $\pm$ 0.12	3.19 $\pm$ 0.11	0.84 $\pm$ 0.18	1.77 $\pm$ 1.33
V3	2.4	1.83 $\pm$ 0.11	2.86 $\pm$ 0.10	1.11 $\pm$ 0.18	1.09 $\pm$ 1.33
V4	1.9	1.47 $\pm$ 0.11	1.73 $\pm$ 0.07	n.d.	2.23 $\pm$ 1.33
V5	10.7	7.82 $\pm$ 0.36	10.4 $\pm$ 0.3	1.99 $\pm$ 0.19	6.58 $\pm$ 1.35
V6	25.2	20.6 $\pm$ 0.9	27.6 $\pm$ 0.8	5.31 $\pm$ 0.24	13.2 $\pm$ 1.4

tios similar to the production ratio (about 2.9 for step-crushing and laser ablation results indicated in Table 2 and Table 4), so we

assume that these two samples represent a nearly non-degassed magma.

**Table 4**  
Isotopic compositions of the pierced vesicles for the two Galápagos samples. Ra is the air  $^3\text{He}/^4\text{He}$  ratio equals to  $1.384 \times 10^{-6}$ .  $^{40}\text{Ar}^*$  is radiogenic  $^{40}\text{Ar}$ . Errors are 1 sigma uncertainties. n.d. means not determined when gases from the vesicle are indistinguishable from the matrix blank. We consider a  $^{40}\text{Ar}/^{36}\text{Ar}$  ratio of 298.6 for the air as obtained by Lee et al. (2006) to determine the  $^{40}\text{Ar}/^{36}\text{Ar}$  and  $^4\text{He}/^{40}\text{Ar}^*$  ratios.

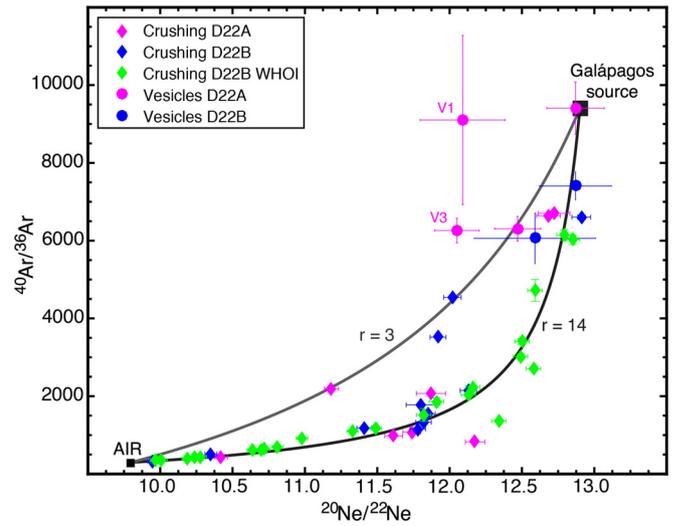
Sample	$\text{CO}_2/{}^3\text{He}$ $\times 10^9$	${}^3\text{He}/{}^4\text{He}$ ( $\times \text{Ra}$ )	${}^{20}\text{Ne}/{}^{22}\text{Ne}$	${}^{21}\text{Ne}/{}^{22}\text{Ne}$	${}^{40}\text{Ar}/{}^{36}\text{Ar}$	${}^4\text{He}/{}^{40}\text{Ar}^*$
<b>AHA-NEMO2-D22A vesicles</b>						
V1	$2.28 \pm 0.13$	$22.57 \pm 0.25$	$12.09 \pm 0.29$	$0.0326 \pm 0.0023$	$9102 \pm 2173$	$3.11 \pm 1.86$
V2	$2.81 \pm 0.16$	$22.42 \pm 0.17$	$12.87 \pm 0.20$	$0.0329 \pm 0.0016$	$9407 \pm 672$	$2.39 \pm 0.44$
V3	$2.35 \pm 0.13$	$22.54 \pm 0.17$	$12.05 \pm 0.15$	$0.0320 \pm 0.0013$	$6260 \pm 315$	$2.75 \pm 0.37$
V4	$2.37 \pm 0.13$	$22.87 \pm 0.16$	$12.47 \pm 0.16$	$0.0325 \pm 0.0013$	$6306 \pm 319$	$2.94 \pm 0.38$
<b>AHA-NEMO2-D22B vesicles</b>						
V1	$2.16 \pm 0.13$	$22.97 \pm 0.33$	$11.43 \pm 0.47$	$0.0295 \pm 0.0031$	$6564 \pm 1618$	$2.88 \pm 1.64$
V2	$2.03 \pm 0.14$	$22.97 \pm 0.43$	$12.15 \pm 0.71$	$0.0359 \pm 0.0053$	$6310 \pm 2536$	$3.01 \pm 2.60$
V3	$2.06 \pm 0.15$	$22.46 \pm 0.48$	$11.08 \pm 0.49$	$0.0333 \pm 0.0035$	$8682 \pm 5704$	n.d.
V4	$2.57 \pm 0.23$	$24.01 \pm 0.72$	n.d.	n.d.	$3283 \pm 911$	$2.59 \pm 1.74$
V5	$2.40 \pm 0.13$	$22.63 \pm 0.21$	$12.59 \pm 0.42$	$0.0355 \pm 0.0022$	$6061 \pm 653$	$2.75 \pm 0.65$
V6	$2.36 \pm 0.13$	$22.89 \pm 0.16$	$12.87 \pm 0.25$	$0.0329 \pm 0.0014$	$7408 \pm 363$	$2.93 \pm 0.35$



**Fig. 2.** The three-neon isotope diagram,  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  versus  ${}^{21}\text{Ne}/{}^{22}\text{Ne}$ , for analyses of step-crushing of sample AHA-NEMO2-D22A (pink diamonds) and sample AHA-NEMO2-D22B (blue diamonds). Data of single vesicles are also shown for sample AHA-NEMO2-D22A (mauve circles) and for the vesicles V5 and V6 of sample AHA-NEMO2-D22B (blue circles). Only vesicles for which the matrix blank contribution is less than 60% are taking into account. (Refer to Table 2 and Table 4.) For comparison, individual step-heating and crushing results (for which  ${}^{20}\text{Ne}$  is at least of  $1 \times 10^{-11}$  ccSTP) of Kurz et al. (2009) for Fernandina samples are indicated in gray (small points) without their uncertainties in order not to overload the figure. *mfl* stands for mass fractionation line and is obtained from a Rayleigh distillation law (Sarda et al., 1988). SW: Solar wind (Heber et al., 2009), SUN (Heber et al., 2012), Ne-B: Neon B, the value of 12.52 as defined by Black (1972) and the value of 12.73 proposed by Moreira (2013) and Moreira and Charnoz (2016). (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

### 3.2. Neon and argon isotopic ratios

The neon isotopic compositions of the two Galápagos samples are presented in Fig. 2 for the step-crushing and laser ablation analyses. Only vesicles for which the matrix blank contribution is less than 60% (and in most cases less than 30%) are considered to be reliable: this includes four vesicles of sample AHA-NEMO2-D22A and the vesicles V5 and V6 of sample AHA-NEMO2-D22B (Table 4). Step-crushing data show mixing between an atmospheric component and a mantle-derived component. The highest  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  ratios measured by step-crushing are  $12.68 \pm 0.08$ ,  $12.72 \pm 0.11$  and  $12.91 \pm 0.07$  ( $1\sigma$ ) at IPGP and  $12.79 \pm 0.05$  and  $12.85 \pm 0.05$  ( $1\sigma$ ) at WHOI. The data follow the same trend as those of Kurz et al. (2009) but these values are the highest ever determined for the Galápagos hotspot since Kurz et al. (2009)



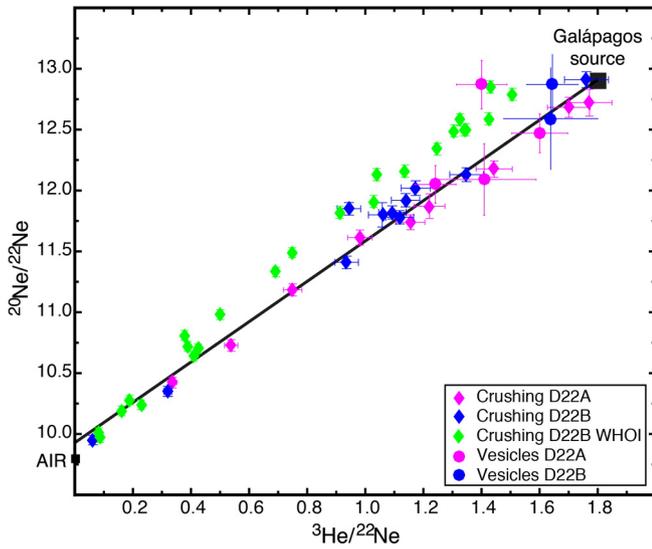
**Fig. 3.**  ${}^{40}\text{Ar}/{}^{36}\text{Ar}$  versus  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  for the same data as in Fig. 2, that is data of step-crushing for the two Galápagos samples, data of single vesicles for sample AHA-NEMO2-D22A and data of the vesicles V5 and V6 of sample AHA-NEMO2-D22B. Data of step-crushing show a mixture between an air-like component and a mantle-like component. Two mixing hyperbolas have been drawn in order to fit the data with  $r = 3$  and  $r = 14$ . The  $r$  parameter indicates the hyperbola curvature and is equal to  $({}^{36}\text{Ar}/{}^{22}\text{Ne})_{\text{AIR}} / ({}^{36}\text{Ar}/{}^{22}\text{Ne})_{\text{MANTLE}}$ . The hyperbolas have been derived considering an air-like endmember and a mantle endmember (with a  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  ratio of 12.9 and a  ${}^{40}\text{Ar}/{}^{36}\text{Ar}$  ratio of 9400).

found a  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  ratio as high as  $12.57 \pm 0.08$  ( $1\sigma$ ) for Fernandina glasses. These are among the highest terrestrial values, as discussed further below.

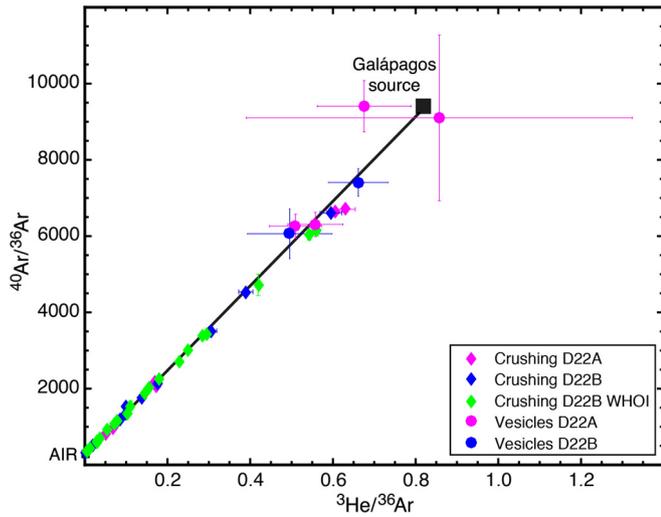
The isotopic compositions of single vesicles by laser ablation are very similar to the highest ratios obtained by step-crushing. Two vesicles have a  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  ratio of  $12.87 \pm 0.2$  ( $1\sigma$ ). Uncertainties on laser ablation data are larger than for step-crushing data because the blank corrections are more important for vesicles than for the crushing steps, and the quantity of neon introduced into the mass spectrometer is smaller.

As for the  ${}^{21}\text{Ne}/{}^{22}\text{Ne}$  ratio, the highest values obtained by step-crushing are between 0.0332 and 0.0341 whereas the  ${}^{21}\text{Ne}/{}^{22}\text{Ne}$  ratios of vesicles are 0.0325–0.0330. The extrapolated  ${}^{21}\text{Ne}/{}^{22}\text{Ne}$  ratio corresponding to a  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  ratio of 12.9 (the highest value obtained) is about 0.0337.

Fig. 3 shows that argon obtained by step-crushing consists of mixing between an air-like component and a mantle component. The highest  ${}^{40}\text{Ar}/{}^{36}\text{Ar}$  ratios are between  $6605 \pm 57$  and  $6702 \pm 76$  ( $1\sigma$ ) and are associated with the highest  ${}^{20}\text{Ne}/{}^{22}\text{Ne}$  ratios (Table 2). Moreover, the laser extracted vesicle data are sim-



**Fig. 4.**  $^{20}\text{Ne}/^{22}\text{Ne}$  versus  $^3\text{He}/^{22}\text{Ne}$  for the two Galápagos samples (same data as in Fig. 2 that is data of step-crushing for the two Galápagos samples, data of single vesicles for sample AHA-NEMO2-D22A and data of the vesicles V5 and V6 of sample AHA-NEMO2-D22B). Data of step-crushing show a mixture between an air-like component and a mantle component.



**Fig. 5.**  $^{40}\text{Ar}/^{36}\text{Ar}$  versus  $^3\text{He}/^{36}\text{Ar}$  for the two Galápagos samples (same data as in Fig. 2, that is data of step-crushing for the two Galápagos samples, data of single vesicles for sample AHA-NEMO2-D22A and data of the vesicles V5 and V6 of sample AHA-NEMO2-D22B). Data of step-crushing show a mixture between an air-like component and a mantle component.

ilar to the highest values obtained by step-crushing even if the  $^{40}\text{Ar}/^{36}\text{Ar}$  ratio presents larger variability, between  $6061 \pm 653$  and  $9407 \pm 672$  (Table 4).

### 3.3. Elemental ratios

Fig. 4 and Fig. 5 allow an estimate for the  $^3\text{He}/^{22}\text{Ne}$  and  $^3\text{He}/^{36}\text{Ar}$  elemental ratios of the Fernandina (Galápagos) source. Step-crushing data on these figures show mixing between an air-like component and a mantle component. The linear correlations observed allow extrapolation of the  $^3\text{He}/^{22}\text{Ne}$  and  $^3\text{He}/^{36}\text{Ar}$  ratios respectively for the  $^{20}\text{Ne}/^{22}\text{Ne}$  and  $^{40}\text{Ar}/^{36}\text{Ar}$  ratios determined by the analyses (section 3.2). Such a method is consistent only for samples having a  $^4\text{He}/^{40}\text{Ar}^*$  ratio similar to the production ratio (Moreira and Kurz, 2013) so that elemental fractionation between noble gases (e.g. by degassing) is very unlikely to have occurred.

Considering a  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of 12.91 (as suggested in Fig. 2) for the mantle source of the Galápagos hotspot, the extrapolation gives a  $^3\text{He}/^{22}\text{Ne}$  ratio of 1.80 (Fig. 4) with the IGP data and of 1.50 with the WHOI data. This relatively small difference might be due to the use of different glass chunks (sample heterogeneity). Kurz et al. (2009) determined a lower value of about 1 extrapolated to the maximum  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of 12.5 (at that time) for Fernandina samples. Trieloff et al. (2000) found a higher ratio of 2.5 and Mukhopadhyay (2012) a similar ratio of about 2 for the Icelandic hotspot. If we consider that the  $^{40}\text{Ar}/^{36}\text{Ar}$  ratio of the Fernandina source is 9400 (as suggested in Fig. 3), we can infer that the  $^3\text{He}/^{36}\text{Ar}$  ratio of the Fernandina source (and by implication, the lower mantle) is about 0.81 (Fig. 5). This is higher than the value of 0.35 found by Raquin and Moreira (2009) for Fernandina (Galápagos) samples and the values of 0.4 and 0.53 for Iceland samples found by Trieloff et al. (2000) and Mukhopadhyay (2012) respectively.

Hence, these results highlight the large contrast between MORB and OIB (Moreira and Kurz, 2013). The MORB source is characterized with a  $^3\text{He}/^{22}\text{Ne}$  ratio of at least 5 (Moreira et al., 1998; Tucker and Mukhopadhyay, 2014) and a  $^3\text{He}/^{36}\text{Ar}$  ratio of 0.45 (Moreira et al., 1998; Raquin et al., 2008).

## 4. Discussion

### 4.1. The air-like component in oceanic basalts

Step-crushing data are characterized by a mixing between a mantle-like component and an air-like component. The latter has been attributed either to atmospheric contamination (Ballentine and Barfod, 2000) or to atmospheric noble gas recycling into the mantle via subduction (Sarda, 2004) or to both processes for heavy noble gases (Holland and Ballentine, 2006; Parai and Mukhopadhyay, 2015; Tucker et al., 2012). However, none of the pierced vesicles, which span a large range of sizes (diameter ranges from 147  $\mu\text{m}$  to 333  $\mu\text{m}$ , volumes are indicated in Table 3) with high internal pressures, show an atmospheric composition. Several previous authors carried out laser ablation analyses and never found such a component in vesicles, in agreement with our results (Burnard, 1999; Burnard et al., 1997; Colin et al., 2013, 2015; Raquin et al., 2008). The observation that no vesicle, with a high internal pressure, has a pure atmospheric composition argues against the model of Sarda (2004). The latter model suggests that vesicles, in particular the big ones, could have an atmospheric composition because of air recycling through subduction into the mantle. Since the recycled material would be more fertile, this would melt first and generate the first vesicles with an atmospheric composition, which would have more time to grow. Note also that the laser ablation and step-crushing isotopic measurements overlap, but that the laser ablation  $^{40}\text{Ar}/^{36}\text{Ar}$  data are systematically higher as a group, demonstrating less atmospheric contamination (see Figs. 3 and 5).

Thus samples are likely contaminated by air once they are brought to the surface, as suggested by Ballentine and Barfod (2000). Air would progressively fill cracks and vesicles linked to these cracks. As seen in Fig. 1, there are many cracks, including microfractures that are hardly visually noticeable inside the glass, which connect some vesicles to one another or directly to the atmosphere. It is also worth noting that we did not manage to pierce some vesicles (no manometric signal), although located via X-ray microtomography. We suspect that these vesicles must already be open with very tiny cracks and filled with air, so no pressure jump was recorded on the manometer. Indeed if these vesicles are filled with air at atmospheric pressure, the manometer accuracy cannot allow the detection of a pressure jump.

We make a simple calculation in order to estimate the signal generated on the Noblesse mass spectrometer by a vesicle and a microfracture of given sizes if they were filled with air. First we consider a large vesicle of 500  $\mu\text{m}$  diameter, which is equivalent to a volume of 0.065  $\text{mm}^3$ , and corresponds to the largest size observed in our samples. If the vesicle is filled by air, then there is about  $2 \times 10^{-9}$   $\text{cm}^3$  STP of  $^{36}\text{Ar}$  in the vesicle since there is 0.934% argon in the atmosphere and the occurrence of  $^{36}\text{Ar}$  represents 0.3364% (Sano et al., 2013). By comparison, the crushing steps with neon and argon atmospheric ratios have a signal on the order of  $6 \times 10^{-10}$ – $5 \times 10^{-9}$   $\text{cm}^3$  STP for  $^{36}\text{Ar}$  (refer to Table 1). Making the same calculation for a microfracture of 2500  $\mu\text{m}$  long and 5  $\mu\text{m}$  in diameter (sizes which were observed by microtomography) would yield a signal of  $1.2 \times 10^{-12}$   $\text{cm}^3$  STP for  $^{36}\text{Ar}$ .

Hence, the atmospheric signature of some crushing steps may be due to one or several vesicles filled with air. Once sample glasses are collected, air may progressively fill the microfractures and the vesicles linked to these cracks. When the sample is placed under ultra-high vacuum for step-crushing analyses, some atmospheric gases at the sample surface are likely to be pumped out. However, the air filling the vesicles may not be completely removed because the cracks are very thin (resulting in very low conductance), unless perhaps the sample is pumped for a very long time.

#### 4.2. Neon and argon compositions of the Galápagos source

Since the results of step-crushing clearly indicate atmospheric contamination, it is difficult to use only these results to constrain the lower mantle composition. However, the vesicle compositions are consistently similar to the highest neon and argon isotopic ratios obtained by step-crushing. The fact that the Ne and Ar data of some vesicles, especially vesicles V1 and V3 of sample D22A (Fig. 3), are a bit scattered is likely due to very small gas quantities and to matrix effects related to laser ablation prior to vesicle penetration. Raquin et al. (2008) noticed that the signal of the matrix (which is ablated just before reaching a vesicle) can alter the signal of the vesicles, so the matrix blank correction may not be perfect for these vesicles. This would lower in particular the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of these vesicles (Fig. 2) since the matrix is imbued with atmospheric gas (Tables S3a and S3b in Supplementary Information). Indeed, it is difficult to apply a perfect matrix blank correction on the vesicle analyses because it is implied to reproduce exactly the same matrix ablation distances/geometry as for the vesicles. This is thus a form of atmospheric contamination, which may affect only the vesicles V1 and V3 of sample D22A. The crushing steps with the highest values must not therefore reflect atmospheric contamination, but instead correspond to mantle gases from several vesicles. Even if laser ablation analyses are associated with larger uncertainties, they are a very useful tool to support the fact that some crushing steps do not reflect any atmospheric contamination, and they allow us then to derive the Galápagos source composition (Figs. 2, 4, 5 in particular). Although the single-vesicle data are slightly scattered, they suggest that the upper limit for the Galápagos mantle source  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio is 12.9, which is different from solar wind and Sun's values.

Considering the new combination data of step-crushing and laser ablation, the Galápagos source  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio is likely to be 12.9, which is higher than the putative Neon B value of 12.5 suggested by Tieloff et al. (2000). The new Galapagos data are consistent with the actual measurements of Tieloff et al. (2000), Mukhopadhyay (2012) and Colin et al. (2015) who found a maximum  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of  $12.85 \pm 0.31$ ,  $12.88 \pm 0.06$  and  $12.73 \pm 0.04$  ( $1\sigma$ ), respectively, for Iceland glasses, another hotspot with primitive helium and neon. Such a range is also close, within un-

certainty, to the highest value of  $13.04 \pm 0.2$  ( $1\sigma$ ) obtained for the Kola peninsula (Russia) by Yokochi and Marty (2004).

Data can be fitted in Fig. 3 with hyperbolas considering mixing between an air-like endmember and a mantle endmember so as to better constrain the Galápagos source argon isotopic composition. At least two hyperbolas can fit the data, defined by different  $r$  parameters:  $r = 3$  and  $r = 14$ . The  $r$  parameter indicates the hyperbola curvature and is defined as the ratio between the air-like endmember ( $^{36}\text{Ar}/^{22}\text{Ne}$ )<sub>AIR</sub> and the mantle endmember ( $^{36}\text{Ar}/^{22}\text{Ne}$ )<sub>MANTLE</sub>. Thus considering a  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of 12.9 and a  $^{40}\text{Ar}/^{36}\text{Ar}$  ratio of 9400 for the Galápagos mantle source provides one of the best fits to the data. Hence, such values for the two isotopic ratios may be upper limits for this mantle source. The fact that several hyperbolas (at least two) fit the data in Fig. 3 may either be due to matrix effects (Raquin et al., 2008) or to the variation of the atmospheric component (its elemental composition) caused by ultra-high vacuum pumping.

We do not present the results for the  $^{38}\text{Ar}/^{36}\text{Ar}$  ratio because of analytical limitations. Blank corrections are quite important due to the very small signals. Further analyses will be required to resolve this analytical challenge so as to determine whether the lower mantle  $^{38}\text{Ar}/^{36}\text{Ar}$  ratio is atmospheric, as suggested by Raquin and Moreira (2009) and Tieloff et al. (2000), or not.

#### 4.3. Origin of helium, neon and argon on Earth

Values of approximately 12.9 for the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of the Ferdinandina mantle source, determined here, are higher than the Neon B values of 12.52 determined by Black (1972), and assumed to be an upper terrestrial mantle limit by Tieloff et al. (2000). However, it is close to the average of 12.8 determined for lunar soils by Eberhardt et al. (1972), and 12.73 calculated by Moreira (2013) and Moreira and Charnoz (2016). Raquin and Moreira (2009), Moreira (2013) and Moreira and Charnoz (2016) used a model of solar wind grain implantation coupled to isotope fractionation and loss by sputtering/erosion, in order to determine the steady state value of the implanted solar wind (Neon B). They found a  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of 12.53 using Heber et al. (2009)'s solar wind value and a  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of 12.73 with Pepin et al. (2012)'s solar wind value. Their model predicted values similar to the results obtained on gas-rich meteorites and on lunar soils (Black, 1972; Eberhardt et al., 1972), but depends upon sputtering/erosion rates and exposure times. If the exposure time is short, the implanted grains will have  $^{20}\text{Ne}/^{22}\text{Ne}$  close to that of the solar wind, but as exposure time increases, the  $^{20}\text{Ne}/^{22}\text{Ne}$  within the grains decreases due to preferential loss of  $^{20}\text{Ne}$ . Over time periods of  $\sim 10^5$  yrs, a steady state is reached, where there is a balance between implantation and erosion and the grain  $^{20}\text{Ne}/^{22}\text{Ne}$  is predicted to approach a single value. Moreira and Charnoz (2016) used a numerical model in order to determine the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of a population of grains submitted to solar wind irradiation at the beginning of solar system formation and for different exposure ages. They showed that the mean  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of a grain population could be higher than the Neon B steady state value and could for example reach 12.9 or even 13, if the steady state is not reached (exposure ages lower than a few kyr). It seems likely that all the solar nebula grains may have had different exposure times and some may not have reached the steady state. This may result in a higher  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio than the steady state value in the parent bodies of the Earth. On the contrary, Black (1972) determined a steady state value for Neon B on lunar soils, which have been exposed for hundreds of millions of years to solar wind implantation. Hence, the estimated Galápagos value of 12.9 could easily be explained with such a scenario, and the observed  $^{20}\text{Ne}/^{22}\text{Ne}$  terrestrial value places important constraints on the early Earth.

The fact that several studies found similar values for OIB (Colin et al., 2015; Mukhopadhyay, 2012; Trieloff et al., 2000) may indicate that this might be the upper limit for the lower mantle ratio. This is also consistent with the highest value of Yokochi and Marty (2004),  $13.04 \pm 0.2$  ( $1\sigma$ ) within uncertainties. Ballentine et al. (2005) also put forward that Neon B is likely the major neon component in the upper mantle.

Yokochi and Marty (2004) obtained a maximum value of  $13.04 \pm 0.2$  ( $1\sigma$ ) for the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio and suggested that the noble gas composition of the lower mantle is inherited from dissolution of solar nebula gas into a magma ocean. The  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of the Sun is estimated to be about 13.4 (Heber et al., 2012) and so Yokochi and Marty (2004) believed that a proto-atmosphere with a solar composition could have been captured early in Earth's history and would have been partly dissolved into the mantle. In the same way, Mukhopadhyay (2012) considered that the composition of the Iceland mantle source is solar-like. Up to now, no study has yet measured values close to 13.4 neither in MORBs nor in OIBs, which seems to question the solar gas dissolution origin for neon.

It can also be argued that the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio was around 13.4 from the dissolution of a proto-atmosphere into the mantle when the Earth formed and that recycling of atmospheric noble gases into the mantle through subduction over billions of years would have then lowered the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio to a value of 12.9. It is difficult to explain the terrestrial  $^{20}\text{Ne}/^{22}\text{Ne}$  with this scenario because it requires enough neon dissolution into the Earth's precursors (magma ocean) before the solar nebula was blown. However, the nebula is blown in a few Myr (e.g., Beckwith and Sargent, 1996 and Wyatt et al., 2003). It also implies, in a later stage, atmospheric neon recycling via subduction into the gas-rich lower mantle, i.e. to rapidly transport the dissolved neon into the interior. Finally, if only neon dissolution from solar nebula occurred, then one would expect to detect higher  $^{20}\text{Ne}/^{22}\text{Ne}$  ratios than 12.9 in the lower mantle since neon atmospheric recycling into the lower mantle is unlikely (Holland and Ballentine, 2006; Staudacher and Allègre, 1988) or may not have totally erased the solar signature. Marty (2012) has proposed that the depleted mantle consists of a mixture between a solar-like endmember and a Neon A endmember (this component is found in chondrites) to explain its  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio. This scenario implies recycling of the Neon A component into the upper mantle, since the latter component may have been brought in a later stage (e.g., O'Brien et al., 2014). However, even if subduction of heavy noble gases may occur (Holland and Ballentine, 2006; Kendrick et al., 2011; Mukhopadhyay, 2012; Parai and Mukhopadhyay, 2015; Parai et al., 2012; Tucker et al., 2012), subduction of light noble gases (helium and neon) into the upper mantle may be very difficult (Holland and Ballentine, 2006; Staudacher and Allègre, 1988).

However, it is important to note that the two models (solar wind implantation and dissolution of solar nebula gas) are not mutually exclusive. It is possible that a part of the terrestrial neon inventory derives from dissolution of solar nebula gas during the Earth's parent body formation. In this case, the Earth neon would correspond to a mixture between solar neon and solar wind implanted neon and only  $\sim 30\%$  of the Earth neon could come from solar nebula gas to account for the OIB  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio, considering a binary mixing between a solar endmember ( $^{20}\text{Ne}/^{22}\text{Ne} = 13.4$  (Heber et al., 2012)) and a Neon B endmember.

For all these reasons, the Earth (at least the lower mantle) is likely to have mainly acquired its primordial helium and neon from its solid grain precursors, which would have involved solar wind implantation associated with sputtering. The lower mantle neon isotopic composition would then reflect solar wind neon irradiation, which may not be at steady state, explaining the slightly higher  $^{20}\text{Ne}/^{22}\text{Ne}$ . This scenario does not require light noble gas recycling through subduction to account for the lower mantle com-

position. We note that there are still relatively few hotspot helium, neon and argon measurements, so this conclusion is based on existing data. If the  $^{38}\text{Ar}/^{36}\text{Ar}$  ratio is atmospheric in the mantle as suggested by Trieloff et al. (2000) and Raquin and Moreira (2009), then a more complicated scenario may be put forward for argon origin, unless argon was acquired in the same way as helium and neon.

## 5. Conclusion

We analyzed two oceanic island basalt glasses from Fernandina volcano (Galápagos) by step-crushing and laser ablation of single vesicles in order to derive their noble gas compositions. This study was carried out with the aim of determining which hypothesis for noble gas origin on Earth is more likely: solar wind implantation associated with sputtering (Raquin and Moreira, 2009; Trieloff et al., 2000) or dissolution of solar nebula gas into the mantle (Yokochi and Marty, 2004).

The data suggest that the mantle source of the Fernandina (Galápagos) samples has a  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of 12.9 and a  $^{40}\text{Ar}/^{36}\text{Ar}$  ratio of about 9400, consistent with studies on the Iceland hotspot. These two ratios are well constrained with new data from the Fernandina samples. The results of laser ablation analyses provide constraints on the atmospheric component of the step-crushing experiments, and the two methods are hence complementary. The laser ablation single-vesicle measurements have systematically higher  $^{40}\text{Ar}/^{36}\text{Ar}$  than the step crushing experiments, suggesting less atmospheric contamination, making this method promising for obtaining mantle values.

Values close to 12.9 for the  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio of the Galápagos hotspot source are the highest measured so far. They are similar, although slightly higher, to the Neon B composition obtained by Eberhardt et al. (1972) and Raquin and Moreira (2009), Moreira (2013) who respectively found an average of 12.8 and a calculated value of 12.73. The Galápagos source  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio (12.9) is slightly higher than the steady state value of the implanted solar wind coupled to sputtering. This could be explained in two ways. On the one hand, the grains in the accretion disk may not have all reached the steady state, resulting in a mean  $^{20}\text{Ne}/^{22}\text{Ne}$  ratio higher than the steady state, as proposed recently by Moreira and Charnoz (2016). On the other hand, a mixture of dissolved solar neon into planetary embryos having the size of Mars and implanted neon (Neon B) could also produce an isotopic ratio of  $\sim 12.9$ . In this case, only  $\sim 30\%$  of the neon derives from dissolution.

Therefore, even if a solar nebula gas contribution cannot be excluded, the new data are more supportive of the solar wind implantation/sputtering model for explaining the helium and neon origin on Earth.

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## Appendix A. Supplementary material

Supplementary material related to this article can be found online at <http://dx.doi.org/10.1016/j.epsl.2016.05.052>.

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