Reactive bromine chemistry in Mount Etna’s volcanic plume: the influence of total Br, high-temperature processing, aerosol loading and plume–air mixing

T.J. Roberts, R.S. Martin, L. Jourdain

To cite this version:

HAL Id: insu-01095138
https://hal-insu.archives-ouvertes.fr/insu-01095138
Submitted on 15 Dec 2014

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L’archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d’enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.
Reactive bromine chemistry in Mount Etna’s volcanic plume: the influence of total Br, high-temperature processing, aerosol loading and plume–air mixing

T. J. Roberts¹, R. S. Martin², and L. Jourdain¹

¹LPC2E, UMR 7328, CNRS-Université d’Orléans, 3A Avenue de la Recherche Scientifique, 45071 Orleans, CEDEX 2, France
²Department of Geography, University of Cambridge, Downing Place, CB2 3EN, UK

Correspondence to: T. J. Roberts (tjardaroberts@gmail.com)

Received: 4 December 2013 – Published in Atmos. Chem. Phys. Discuss.: 3 March 2014
Revised: 10 July 2014 – Accepted: 22 July 2014 – Published: 23 October 2014

Abstract. Volcanic emissions present a source of reactive halogens to the troposphere, through rapid plume chemistry that converts the emitted HBr to more reactive forms such as BrO. The nature of this process is poorly quantified, yet is of interest in order to understand volcanic impacts on the troposphere, and infer volcanic activity from volcanic gas measurements (i.e. BrO / SO₂ ratios). Recent observations from Etna report an initial increase and subsequent plateau or decline in BrO / SO₂ ratios downwind.

We present daytime PlumeChem model simulations that reproduce and explain the reported trend in BrO / SO₂ at Etna including the initial rise and subsequent plateau. Suites of model simulations also investigate the influences of volcanic aerosol loading, bromine emission, and plume–air mixing rate on the downwind plume chemistry. Emitted volcanic HBr is converted into reactive bromine by autocatalytic bromine chemistry cycles whose onset is accelerated by the model high-temperature initialisation. These rapid chemistry cycles also impact the reactive bromine speciation through inter-conversion of Br, Br₂, BrO, BrONO₂, BrCl, HOBr.

We predict a new evolution of Br speciation in the plume. BrO, Br₂, Br and HBr are the main plume species near downwind whilst BrO and HOBr are present further downwind (where BrONO₂ and BrCl also make up a minor fraction). BrNO₂ is predicted to be only a relatively minor plume component.

The initial rise in BrO / SO₂ occurs as ozone is entrained into the plume whose reaction with Br promotes net formation of BrO. Aerosol has a modest impact on BrO / SO₂ near-downwind (< ~6 km, ~10 min) at the relatively high loadings considered. The subsequent decline in BrO / SO₂ occurs as entrainment of oxidants HO₂ and NO₂ promotes net formation of HOBr and BrONO₂, whilst the plume dispersion dilutes volcanic aerosol so slows the heterogeneous loss rates of these species. A higher volcanic aerosol loading enhances BrO / SO₂ in the (> 6 km) downwind plume.

Simulations assuming low/medium and high Etna bromine emissions scenarios show that the bromine emission has a greater influence on BrO / SO₂ further downwind and a modest impact near downwind, and show either complete or partial conversion of HBr into reactive bromine, respectively, yielding BrO contents that reach up to ~50 or ~20% of total bromine (over a timescale of a few 10 s of minutes).

Plume–air mixing non-linearly impacts the downwind BrO / SO₂, as shown by simulations with varying plume dispersion, wind speed and volcanic emission flux. Greater volcanic emission flux leads to lower BrO / SO₂ ratios near downwind, but also delays the subsequent decline in BrO / SO₂, and thus yields higher BrO / SO₂ ratios further downwind. We highlight the important role of plume chemistry models for the interpretation of observed changes in BrO / SO₂ during/prior to volcanic eruptions, as well as for quantifying volcanic plume impacts on atmospheric chemistry. Simulated plume impacts include ozone, HOₓ and NOₓ depletion, the latter converted into HNO₃. Partial recovery of ozone occurs with distance downwind, although cumulative ozone loss is ongoing over the 3 h simulations.

Published by Copernicus Publications on behalf of the European Geosciences Union.
1 Introduction

The discovery of volcanic BrO (Bobrowski et al., 2003), and its subsequent observation in many volcanic plumes globally (e.g. Oppenheimer et al., 2006; Bobrowski et al., 2007; Bobrowski and Platt, 2007; Kern et al., 2009; Bani et al., 2009; Loubian et al., 2009; Theys et al., 2009; Boichu et al., 2011; Heue et al., 2011; Bobrowski and Giuffrida, 2012; Rix et al., 2012; Hörmann et al., 2013; Kelly et al., 2013; Lübcke et al., 2014), demonstrates the reactivity of volcanic halogen emissions in the troposphere. Volcanoes release H$_2$O, CO$_2$ and SO$_2$, but also a range of hydrogen halides to the atmosphere including HF, HCl and HBr (in descending order of abundance in the emission, see e.g. Aiuppa et al., 2005). HF is too strong an acid for reactive halogen cycling, but for HBr and HCl, observational evidence shows these are not simply just washed out from the atmosphere, but can undergo transformation into reactive halogen species.

Notably, DOAS (differential optical absorption spectroscopy) measurements show that BrO forms at hundreds of pmol mol$^{-1}$ to nmol mol$^{-1}$ mixing ratios just minutes downwind, an order of magnitude higher than that found in the Arctic, where BrO episodes of up to tens of pmol mol$^{-1}$ cause significant ozone depletion and mercury deposition events (Simpson et al., 2007). Additionally, there is potential to use long-term BrO monitoring at volcanoes as an indicator of volcanic activity (Bobrowski and Giuffrida, 2012). Thus there is strong interest in developing models to simulate the formation of reactive bromine (and chlorine) in volcanic plumes, and to predict the downwind impacts from both quiescently degassing volcanoes and episodic eruptions to the troposphere. Studies to date usually use equilibrium models to predict the high-temperature chemistry of the near-vent plume, which is then used to initialise kinetic atmospheric models of the downwind reactive halogen chemistry (Bobrowski et al., 2007; Roberts et al., 2009; von Glasow, 2010; Kelly et al., 2013). See von Glasow et al. (2009) for an overview.

This study uses a purpose-built kinetic model, PlumeChem (Roberts et al., 2009), to investigate the volcanic plume reactive halogen chemistry, focusing here on bromine in a case study for Mt Etna. We include a revised methodology (Martin et al., 2009) for equilibrium calculations used to represent the near-vent high-temperature chemistry, and discuss uncertainties in the use of thermodynamic equilibrium models. Below, we outline the progression of recent research on using equilibrium models for high-temperature near-vent plume chemistry and the development of kinetic models for volcanic plume reactive halogen (BrO) chemistry. We then describe the new findings of this study specifically regarding the plume reactive bromine evolution presented by the model, and to highlight uncertainties in model high-temperature initialisation and the influence of total bromine, aerosol and plume–air mixing on the plume chemistry.

1.1 Application of the HSC equilibrium model to the near-vent plume

HSC is a commercially available model (Outotec, Finland) that predicts the thermodynamic equilibrium composition of a gas mixture at a defined temperature, pressure and chemical composition. Such models are used to represent the composition of the near-vent volcano plume (e.g. Gerlach, 2004; Martin et al., 2006), predicting a vast array (≥ 100) of chemical species. An overview of the input and outputs to HSC is provided in Table 1. The chemical composition of the mixture is determined by combining magmatic (comprising of H$_2$O, CO$_2$, CO, SO$_2$, H$_2$S, H$_2$, HF, HCl, HBr, HI, Hg, typically at around 800–1100°C) and air (N$_2$, O$_2$, Ar, typically around (20°C) components. The magmatic gas composition varies between volcanoes and may be estimated from crater-rim measurements. It is also possible to predict the abundance of gases that are missing from measurements as the magmatic gas H$_2$O—H$_2$, CO$_2$—CO, SO$_2$—H$_2$S equilibria are functions of oxygen fugacity, pressure of degassing and temperature (e.g. Giggenbach, 1987). The resulting HSC output composition depends critically on the assumed ratio of air to magmatic gases in the near-vent plume, V$_A$:V$_M$. However, this ratio is poorly defined, an issue we examine further in this study.

The HSC output is then used to initialise low-temperature kinetic models (such as PlumeChem, Roberts et al., 2009; Kelly et al., 2013; MISTRA, Bobrowski et al., 2007; von Glasow, 2010) of the volcanic plume reactive halogen chemistry including formation of BrO. These models show that elevated radicals in the HSC output accelerate the onset of autocatalytic BrO chemistry, leading to very rapid BrO formation. BrO formation occurs more slowly in kinetic models that are not initialised with high-temperature chemistry. For the interest of atmospheric modellers, we simplify the complex HSC output (≥ 100 species) in Table 1, following Roberts et al. (2009) who identified impacts of HO$_x$, NO$_x$, Br$_x$ and Cl$_x$ on the downwind plume halogen chemistry. Key species are identified to be OH, NO, Br, Cl and Cl$_2$, noting NO$_2$ ≪ NO and HO$_2$ ≪ OH, Br$_2$ ≪ Br in the HSC output. These species act to accelerate autocatalytic reactive bromine formation (see Fig. 4 of Roberts et al., 2009). High-temperature near-vent formation of SO$_3$ (a precursor to H$_2$SO$_4$) also influences the volcanic plume halogen chemistry by providing a source of aerosol surface area.

However, the thermodynamic assumption behind equilibrium models such as HSC may not always be appropriate for volcanic plume applications: Martin et al. (2009) noted that the near-complete re-equilibration (i.e. oxidation) of H$_2$S within HSC is in disagreement with the widespread observed presence of H$_2$S in volcanic plumes (exception: Erebus), and suggested a revised operation of HSC in which H$_2$S is removed prior to re-equilibration. Furthermore, recent measurements confirming volcanic H$_2$S (Aiuppa et al., 2011; Roberts et al., 2012) indicate that this argument also applies...
to H₂, as well as CO (although CO is typically present in low concentrations, with some exceptions, e.g. Mt Erebus). See e.g. Gerlach (2004) for various collated emission compositions. Uncertainties and limitations in the use of HSC to represent the near-vent plume composition are discussed further in this study in the context of downwind BrO chemistry.

### 1.2 Kinetic models of downwind volcanic plume reactive halogen chemistry

Atmospheric chemistry models have been developed in an effort to simulate the reactive halogen chemistry of volcanic plumes, explain observed BrO formation and predict impacts of reactive volcanic halogens on atmospheric chemistry. To date, two models: MISTRA (that simulates an advected column of air, Bobrowski et al., 2007) and PlumeChem (in an expanding box or multi-grid box modes, Roberts et al., 2009) have been developed for this purpose. Initialisation of these models includes the high-temperature chemistry of the near-vent plume, as represented by HSC. Calculations by Oppenheimer et al. (2006) showed BrO formation to be too slow if high-temperature near-vent radical formation is ignored. Bobrowski et al. (2007) performed the first MISTRA kinetic model simulations of volcanic plume reactive halogen chemistry, using a model initialised with HSC at Vₐ : Vₘ of 0 : 100, 15 : 85, 40 : 60, finding that the 40 : 60 simulation yielded highest downwind BrO/SO₂. Roberts et al. (2009) queried the use of such high Vₐ : Vₘ of 40 : 60 which yields rather high SO₃ : SO₂ ratios, that implies volcanic sulfate emissions would exceed volcanic SO₂. Roberts et al. (2009) presented model simulations initialised with HSC at Vₐ : Vₘ of 15 : 85 that reproduced the rapid formation of BrO/SO₂ at a range of Arc (subduction zone) volcanoes for the first time (including Etna, Soufrière Hills, Villarrica), and suggested the higher BrO/SO₂ observed in the Soufrière Hills volcano plume may be fundamentally due to higher Br/S in the emission. A model study by von Glasow (2010) with simulations initialised at Vₐ : Vₘ of 15 : 85 demonstrated good agreement to both reported column abundances of SO₂ and BrO/SO₂ ratios downwind of Etna.

All of the above-mentioned simulations (Bobrowski et al., 2007; Roberts et al., 2009; von Glasow, 2010) predict substantial in-plume depletion of oxidants, including ozone, although to varying extents, and predict contrasting plume halogen evolution. Roberts et al. (2009) also demonstrated conversion of NO₂ into nitric acid via BrONO₂, and proposed this mechanism to explain reported elevated HNO₃ in volcano plumes. Von Glasow (2010) simulated the impacts of volcanic reactive halogen chemistry on mercury speciation, predicting significant conversion to Hg²⁺ in the plume.

A number of observations of ozone abundance in volcanic plumes have recently been reported: Vance et al. (2010) observed ozone depletion in the Eyjafjallajökull plume, and at ground-level on Etna’s flanks (by passive sampling).
Schumann et al. (2011) presented multiple measurements of the downwind plume of Eyjafjallajökull that observed ozone depletion to variable degrees. There exist also observations of depleted ozone in the Mount St Helens plume (Hobbs et al., 1982) that are now believed to be likely due to BrO chemistry. Ozone depletion of up to ~35% was reported in an aircraft study of Mt Erebus plume in Antarctica (Oppenheimer et al., 2010), where BrO has also been observed (Boichu et al., 2011). A systematic instrumental aircraft investigation of ozone depletion in a volcanic plume (where emissions are also quantified) is presented by Kelly et al. (2013), and compared to PlumeChem model simulations over 2 h of plume evolution, finding good spatial agreement in the modelled and observed ozone mixing ratios. At higher altitudes, ozone depletion in a volcanic plume is reported in the UTLS (upper troposphere, lower stratosphere) region observed by Rose et al. (2006), and investigated and attributed to reactive halogen chemistry by Millard et al. (2006).

However, ozone depletion has not been universally observed: Baker et al. (2011) did not detect an ozone depletion signal relative to the (somewhat variable) background level during an aircraft transect through Eyjafjallajökull plume. An instrumented aircraft study found no evidence for O₃ depletion in the plume of Nevado del Huila (Colombia) and found ozone levels 70–80% of ambient in the plume of Tungurahua (Ecuador), which could not be conclusively attributed to BrO chemistry (Carn et al., 2011).

A number of modelling discrepancies also exist. For example, the model studies of Roberts et al. (2009), von Glasow (2010) and Kelly et al. (2013) predict contrasting Br speciation and contrast in predicted impacts on ozone and other oxidants. This may reflect differences in the model representations and modelling uncertainties or demonstrate volcano-specific differences in the plume chemistry. Navigating the vast model parameter space of volcanic plume chemistry is challenging due to the non-linear controls on the plume chemistry of multiple inter-dependent parameters including volcanic aerosol, rate of horizontal dispersion, rate of vertical dispersion, wind speed, volcanic gas flux, bromine in the emission and high-temperature radical formation. Limited observational data sets are available to compare to the models, and the available data do not fully constrain the high- and low-temperature plume chemistry. To provide further insight, this study presents new PlumeChem model simulations to compare to recently reported trends in BrO/SO₂ ratios, and illustrates several of the major controls and uncertainties in the reactive halogen chemistry of volcanic plumes.

2 Methods

2.1 HSC: equilibrium modelling of near-vent plume chemistry

The use of HSC for calculating the composition of the near-vent plume is described by Gerlach (2004) and Martin et al. (2006). This study uses HSC thermodynamic model version 7.1, and applies the modifications which were proposed by Martin et al. (2009). A simple background atmosphere of N₂ (78%), O₂ (21%) and Ar (1%) is assumed for the HSC calculations. The magmatic composition used for Etna follows that of Bagnato et al. (2007), with gas mixing ratios for H₂O, CO₂, SO₂, HCl, H₂S, H₂, CO, of 0.86, 9.6 × 10⁻², 2.9 × 10⁻², 1.4 × 10⁻², 1.5 × 10⁻⁻³, 3.7 × 10⁻⁴ and 3.5 × 10⁻⁴ respectively. Hg is excluded for the purposes of this study due to low abundance in the volcanic emission. The bromine content as HBr, was set to be either medium, high or low. “Medium” bromine (molar mixing ratio of 2.16 × 10⁻⁵, equivalent to a total bromine to SO₂ ratio (Brtot/ SO₂) in the emission of 7.4 × 10⁻⁴) corresponds to the average Br/S molar ratio at Etna north-east crater determined from filter-pack measurements over 2004, Aiuppa et al. (2005). “High” bromine (mixing ratio of 7.03 × 10⁻⁵, equivalent to Brtot/ SO₂ in the emission of 2.4 × 10⁻³) corresponds to that assumed in a previous model study of Etna (von Glasow, 2010), and is in the upper range (within one standard deviation) of the observations of Aiuppa et al. (2005). Simulations are also performed at a “lower” Brtot/ SO₂ = 4.8 × 10⁻⁴ which corresponds to a filter-pack Br/S measurement at Voragine crater reported by Oppenheimer et al. (2006). These are summarised in Table 2.

The magmatic temperature is set to 1050°C in order to match that prescribed by von Glasow (2010), although we note that Metrich and Rutherford (1998) estimated Etna magmatic temperature to be 1100°C. For the near-vent plume mixture input to HSC, ambient air temperature was set to 20°C. This is somewhat high considering Etna’s elevation (3 km), but this has a minor influence on the HSC output (especially considering 50°C difference in the magmatic temperature estimates outlined above). For the actual PlumeChem atmospheric chemistry model runs, the atmospheric temperature was a more realistic 285 K. The equilibrium composition was calculated for standard operation of HSC (in which H₂ and H₂S are allowed to re-equilibrate) and in a revised (Martin et al., 2009) operation of HSC (in which H₂ and H₂S are replaced by inert Ar such that they do not re-equilibrate). The HSC calculations were performed over 16 different V_A/V_M ranging from 0:100 to 15:85.
Table 2. Parameters varied in PlumeChem sensitivity studies.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Values</th>
</tr>
</thead>
<tbody>
<tr>
<td>HSC $V_A : V_M$</td>
<td>0 : 100</td>
</tr>
<tr>
<td></td>
<td>2 : 98</td>
</tr>
<tr>
<td></td>
<td>5 : 95</td>
</tr>
<tr>
<td></td>
<td>10 : 90</td>
</tr>
<tr>
<td></td>
<td>15 : 85</td>
</tr>
<tr>
<td>Aerosol loading:</td>
<td></td>
</tr>
<tr>
<td>$\mu m^2$ molec$^{-1}$ SO$_2$</td>
<td></td>
</tr>
<tr>
<td>High</td>
<td>$10^{-10}$</td>
</tr>
<tr>
<td>Medium</td>
<td>$10^{-11}$</td>
</tr>
<tr>
<td>Br$_{tot}$ / SO$_2$: molar ratio</td>
<td>$7.4 \times 10^{-4}$</td>
</tr>
<tr>
<td>Medium</td>
<td>$2.4 \times 10^{-3}$</td>
</tr>
<tr>
<td>High</td>
<td>$4.8 \times 10^{-4}$</td>
</tr>
<tr>
<td>Low</td>
<td></td>
</tr>
<tr>
<td>Gas Flux kg s$^{-1}$ SO$_2$</td>
<td>10, 20</td>
</tr>
<tr>
<td>(small variations)</td>
<td></td>
</tr>
<tr>
<td>(large variations)</td>
<td>10, 50, 100</td>
</tr>
<tr>
<td>Wind-speed, m s$^{-1}$</td>
<td>3, 5, 10, 15</td>
</tr>
<tr>
<td>Dispersion</td>
<td>B, C, D</td>
</tr>
<tr>
<td>Pasquill–Gifford cases</td>
<td></td>
</tr>
</tbody>
</table>

2.2 PlumeChem: kinetic model of downwind BrO chemistry

The PlumeChem model simulates the reactive halogen chemistry of volcanic plume, as described by Roberts et al. (2009). It can be run in single-box (Roberts et al., 2009) or multi-box (Kelly et al., 2013) modes. Here we used the single box that expands as a background atmosphere is entrained into it, representing dispersion of the plume as it is advected downwind. PlumeChem includes a background atmospheric chemistry scheme and bromine and chlorine reactive halogen chemistry, including photolysis, gas-phase and heterogeneous (gas–aerosol) phase reactions. Autocatalytic formation of BrO occurs through cycles involving reaction of BrO with oxidants, (HO$_2$ NO$_2$), (Reactions R1, R2), aerosol-phase heterogeneous chemistry (Reactions R3, R4) to release a halogen dimer, whose photolysis generates two halogen radicals (Reaction R5), which may react with ozone (Reaction R6) to form BrO. The heterogeneous reactive uptake of HOBr and BrONO$_2$ on volcanic aerosol are thus key drivers of reactive halogen formation. Within the volcanic aerosol, aqueous-phase equilibria (Wang et al., 1994) control the nature of the product, which is Br$_2$ for a typical volcanic plume composition, thereby enabling autocatalytic formation of reactive bromine. Once aerosol Br$_{aq}$ becomes depleted (as a consequence of the BrO formation cycles), BrCl becomes a significant product from the heterogeneous Reactions (R3, R4), leading to non-autocatalytic formation of reactive chlorine. Repeated cycling around Reactions (R1–R6) can cause substantial ozone loss (orders of magnitude greater than the BrO mixing ratio). Repeated cycling between BrO and Br (Reactions R6, R7) further enhances ozone loss in concentrated plume environments:

$$\text{BrO} + \text{HO}_2 \rightarrow \text{HOBr} + \text{O}_2$$ (R1)
$$\text{BrO} + \text{NO}_2 \rightarrow \text{BrONO}_2$$ (R2)
$$\text{HOBr} + \text{H}_2^+ + \text{Br}_{aq} \rightarrow \text{Br}_2(aq-g) + \text{H}_2\text{O}$$ (R3)
$$\text{BrONO}_2 + \text{H}_2\text{O}(l) \rightarrow \text{HOBr} + \text{HNO}_3$$ (R4)
$$\text{Br}_2 \xrightarrow{h\nu} 2\text{Br}$$ (R5)
$$\text{Br} + \text{O}_3 \rightarrow \text{BrO} + \text{O}_2$$ (R6)
$$\text{BrO} + \text{BrO} \rightarrow 2\text{Br} + \text{O}_2.$$ (R7)

The background atmosphere chemistry scheme used here is identical to that of Roberts et al. (2009), assuming a somewhat polluted atmosphere. For the model simulations initialised around midday, background ozone is $\sim 60$ nmol mol$^{-1}$, and NO$_x$ and HO$_x$ are around 0.17 and 30 pmol mol$^{-1}$ respectively, with an ambient temperature of 285 K and 60% relative humidity (RH). Plume dispersion is defined according to Pasquill–Gifford dispersion schemes (see Supplement). The base-run plume dispersion parameterisation used in this study is identical to that of Roberts et al. (2009), based on Pasquill–Gifford case D, with a SO$_2$ gas flux of 10 kg s$^{-1}$ at a wind-speed of 10 m s$^{-1}$. The influence of variations in wind speed (3–15 m s$^{-1}$), volcanic emission flux (10–20 kg s$^{-1}$), SO$_2$ gas flux of 10 kg s$^{-1}$ at a wind-speed of 10 m s$^{-1}$, and dispersion rates (Pasquill–Gifford cases B, C, D) on downwind BrO / SO$_2$ ratios are also shown, as well as simulations with much greater volcanic emission flux (5–10× the base run). Volcanic aerosol loading in the model is investigated as part of the study, and for the majority of simulations is set to $10^{-11}$ $\mu m^2$ molec$^{-1}$ SO$_2$, a factor of 10 lower than that of Roberts et al. (2009), following the PlumeChem model set-up used in Kelly et al. (2013).

The reaction of Br with BrONO$_2$ to form Br$_2$ + NO$_3$ (Orlando and Tyndall, 1996) was added to the PlumeChem model in this study. This reaction influences the overall rate of HBr conversion into reactive bromine as follows: as a sink for BrONO$_2$ it slows the conversion of HBr into reactive bromine as less BrONO$_2$ undergoes heterogeneous uptake (which converts HBr into Br$_2$ via HOBr). However, as a sink for Br it slows the conversion of reactive bromine back into HBr from the reaction Br + HCHO. Under a high volcanic aerosol loading the former dominates, whilst the latter is more important at lower aerosol loadings. It is noted that this reaction is neither included in the IUPAC Kinetics nor JPL Data evaluation databases, thus is not necessarily included “as standard” in all atmospheric models of reactive halogen chemistry.

BrNO$_2$ was suggested by von Glasow (2010) to be an important reservoir for Br in the near-downwind plume, based on the assumed formation of BrNO$_2$ from volcanic NO$_x$ and Br radicals at a rate that exceeds BrNO$_2$ loss via
photolysis. Formation of BrNO$_2$ was not included in previous PlumeChem model studies (Roberts et al., 2009; Kelly et al., 2013). Here, the fate of the products (BrNO$_2$ but also BrONO) from reaction of Br + NO$_2$ are investigated in more detail to evaluate the potential of BrNO$_2$ to influence the plume chemistry.

3 Results

3.1 Model SO$_2$ column abundance, and variability in simulated BrO/SO$_2$

The formation of volcanic BrO is typically reported relative to SO$_2$, which, given slow in-plume oxidation, acts as a plume tracer on the observation timescales (typically minutes to hours). Therefore, prior to comparing PlumeChem model output to the observed BrO/SO$_2$, a comparison is made between the simulated and reported SO$_2$ column abundances. Figure 1a shows slant SO$_2$ column abundance in Mt Etna’s plume over 2004–2005, reported from DOAS (Differential Optical Absorption Spectroscopy) observations from Oppenheimer et al. (2006) and Bobrowski et al. (2007). The data show a general decline with distance downwind, with the exception of two very near source measurements, which may have been underestimated in the very strong near-source plume, see discussion by Kern et al. (2012) and Bobrowski and Guiffrida (2012) for improved SO$_2$ evaluation. Also shown in Fig. 1a is the model downwind plume SO$_2$ column abundance calculated for the plume in the vertical. The decline in modelled SO$_2$ column abundance with distance (or time) downwind is largely due to dispersion, given the slow rate of in-plume SO$_2$ oxidation. The rate of dispersion depends on plume depth, width, gas flux and wind speed during each DOAS measurement, which are not fully constrained by available observations. Nevertheless, the broad agreement between model and observations indicates a suitable model parameterisation of plume–air mixing in the base run. This supports the use of further simulations to investigate the plume halogen chemistry using this plume–air mixing parameterisation scenario, for comparison to reported BrO/SO$_2$ observations.

Figure 1b shows formation of BrO (relative to plume tracer SO$_2$) for a range of model simulations presented later in this study, all using this same plume–air mixing parameterisation, but where the other parameters (volcanic aerosol loading, total plume bromine, initialisation using thermodynamic model output) are varied. Clearly, these variables can have a strong influence on the downwind plume halogen chemistry. Also shown are BrO/SO$_2$ ratios reported by Oppenheimer et al. (2006) and the observed trend in (mean) BrO/SO$_2$ with distance downwind reported by Bobrowski et al. (2007). Several, but not all of the model simulations in Fig. 1b conform to the BrO/SO$_2$ observations. Indeed, simulations whose initialisations assume no plume–air mixing at high temperature typically underestimate downwind BrO/SO$_2$ (see Sect. 3.3 for further discussion). The remaining model runs demonstrate broad agreement with the BrO/SO$_2$ measurements and provide an explanation for the observed rise and subsequent plateau or decline in BrO/SO$_2$ with distance downwind reported by Bobrowski and Giuffrida (2012).

In order to provide further insight into the factors controlling volcano plume reactive halogen chemistry, we investigate here the influence of the above-mentioned variables, and
particularly uncertainties regarding the initialisation by HSC. To do so, suitable values for the volcanic bromine and aerosol loading are first identified, as outlined below.

3.2 The effect of aerosol and bromine content on downwind BrO/SO₂

Highlighted in black in Fig. 1b are four model runs that assume the “medium” and “high” bromine (Br₉/SO₂) emission scenarios (see Table 2), and two contrasting aerosol surface area loadings – namely “high” aerosol estimated as ∼10⁻¹⁰ µm⁻² molec SO₂⁻¹ following Roberts et al. (2009), and the “medium” aerosol estimate, which is an order of magnitude lower, 10⁻¹¹ µm⁻² molec SO₂⁻¹ as used by Kelly et al. (2013). Both the volcanic aerosol loading and volcanic bromine content influence the downwind BrO/SO₂ evolution, as follows.

In general, a higher Br₉/SO₂ in the emission leads to greater BrO/SO₂ far downwind. This is in accordance with the proposed role of Br/S in the emission to explain order of magnitude variation in BrO/SO₂ ratios across Arc volcanoes (Roberts et al., 2009). A higher aerosol loading promotes the conversion of HBr into reactive forms, and promotes the occurrence of reactive bromine as BrO in the far-downwind plume to its role in the heterogeneous reactive uptake of HOBr and BrONO₂. Interestingly, whilst the volcanic aerosol and bromine content have a strong impact on the plateau in BrO/SO₂ far downwind (both in terms of value and when it is reached), Figure 1b indicates that aerosol and bromine content exert a much more limited impact on BrO/SO₂ in the very young plume during the first ∼8 min (∼5 km) of plume evolution, at least for the plume dispersion conditions simulated. For example, at 36 km downwind, the two contrasting aerosol loadings cause the model BrO/SO₂ to vary from 4.2 × 10⁻⁴ to 1.4 × 10⁻³ (“high” bromine scenario) and from 2 × 10⁻⁴ to 4 × 10⁻⁴ (“medium” bromine scenario), whereas at 6 km downwind all of these model runs predict BrO/SO₂ between 2.5 × 10⁻⁴ and 4 × 10⁻⁴. This near-downwind similarity in BrO/SO₂ (despite varying Br₉/SO₂ as well as aerosol loading) is related to the role of oxidants in forming BrO, and differences in the proportion of HBr converted to reactive bromine. This predicted near-downwind independence of BrO/SO₂ on aerosol loading is consistent with the observations of Bobrowski and Giuffrida (2012) at 6 km downwind that showed BrO/SO₂ was independent of relative humidity (a key control on sulfate aerosol volume hence surface area). A model explanation (see Sect. 3.4 for further discussion) is that near-downwind BrO/SO₂ ratios are primarily controlled by Br to BrO partitioning – itself a function of in-plume ozone mixing ratio – in this region where the plume is still relatively concentrated. See Sect. 3.4 for details of the plume reactive bromine speciation and Sect. 3.6 for further discussion on the plume impacts on atmospheric ozone.

Notably, the simulation with both “high” bromine and the larger aerosol estimate predicts BrO/SO₂ that exceeds reported observations far downwind, and does not reproduce the plateau in BrO/SO₂ beyond ∼5 km downwind of Etna summit recently reported by Bobrowski and Giuffrida (2012). We acknowledge the Etna bromine emission may vary with time and therefore use both medium and high bromine emission scenarios alongside a low Br scenario in the further model simulations of this study. However, for the high-bromine scenario, only the simulation with “medium” aerosol surface area loading appears consistent with the BrO/SO₂ observations by Bobrowski and Giuffrida (2012). The “medium” aerosol surface area loading is considered as likely being more representative and this estimate, ∼10⁻¹¹ µm⁻² molec SO₂⁻¹, is used in all further model simulations of this study. Further discussion on estimates of the aerosol surface area loading is provided in the Supplement.

3.3 The high-temperature near-vent plume – a source of model uncertainty

An important model parameter in the PlumeChem model initialisation is the use of output from thermodynamic model HSC to represent the composition of the high-temperature near-vent plume. Figure 2 illustrates the key species in the HSC output (Table 1) for the near-vent plume of Etna (“medium” Br scenario) using (a) the standard HSC methodology in which H₂ and H₂S re-equilibrate, and (b) the modified method (Martin et al., 2009) whereby H₂S and H₂ do not re-equilibrate. NO, OH, Cl and Br and Cl₂ gas mixing ratios are shown for Vₐ : Vₐ ranging from 0 : 100 to 15 : 85, where Vₐ : Vₐ is the ratio of air to magmatic gases in the near-vent plume (plotted as a fraction in Fig. 2), with the HSC temperature varied according to the mixture of magmatic (1050°C) and ambient (20°C) temperatures.

Of note is a step increase in radical mixing ratios in Fig. 2a (in which H₂ and H₂S re-equilibrate). This is the so-called compositional discontinuity, CD (Gerlach, 2004), which occurs at around Vₐ : Vₐ ~ 0.01 for Etna’s magmatic composition. At the CD, the reduced magmatic gases (H₂O, CO₂, etc.) are essentially fully oxidised (SO₂, H₂O, CO₂), thus addition of further oxidant (increasing Vₐ : Vₐ) yields increases in the mixing ratios of the radicals (Br, Cl, NO, OH). As Vₐ : Vₐ increases further, the greater proportion of air relative to magmatic gases yields a lower HSC temperature, leading to slight declines or a plateau in the mixing ratios of NO and OH, and altering the balance between Cl₂ and Cl radicals (Br₂ remains low over the whole Vₐ : Vₐ range). Formation of Br with increasing Vₐ : Vₐ also leads to a corresponding decrease in its “parent” or “source” species HBr (note that other “parent” species, e.g. HCl, H₂O, are in excess relative to Cl₂ and OH). However, in the revised HSC methodology the CD has shifted to low Vₐ : Vₐ, as first shown by Martin et al. (2009). Indeed, it may no longer be
Mixing ratio (10^{-6} mol/mol) of key species (NO, OH, Br, Cl, Cl\textsubscript{2}) in the HSC output as a function of \(V_a/V_M\), the assumed magmatic–atmospheric gas ratio in the near-vent plume, ranging from 0 (0.00 : 1.00) to 0.18 (0.15 : 0.85). The SO\textsubscript{3} : SO\textsubscript{2} ratios (that prescribe the volcanic sulfate / SO\textsubscript{2} emission) in the HSC output are also shown. (a) Standard operation of HSC in which volcanic H\textsubscript{2}S and H\textsubscript{2} are allowed to re-equilibrate, yielding near-zero mixing ratios of these gases in the HSC output. (b) A revised operation of HSC (Martin et al., 2009) in which volcanic H\textsubscript{2}S and H\textsubscript{2} are removed (and temporarily replaced by inert Ar) such that they do not re-equilibrate within HSC.

The fact that certain species need to be “protected” from re-equilibration within presents a major limitation to the use of thermodynamic models to represent near-vent plume, as neither the choice of \(V_a/V_M\), nor the protection of certain species (but not others) are fully justified on a physical basis. It is likely that some processes may be kinetics limited and thus poorly described by thermodynamic models. Studies suggest that this is indeed the case for the formation of NO\textsubscript{x} from background N\textsubscript{2} entrained into the plume (Martin et al., 2012), due to the high bond strength for N\textsubscript{2} (945 kJ mol). Nevertheless, there is some evidence for the high-temperature formation of radicals in the near-vent plume, for example in the presence of crater-rim sulfate at SO\textsubscript{4}\textsuperscript{2-} : SO\textsubscript{2} \sim 1 : 100 (e.g. Mather et al., 2003; Martin et al., 2008), from which near-vent SO\textsubscript{3} production might be inferred. Further, a volcanic source of HO\textsubscript{x} is suggested by plume H\textsubscript{2}O\textsubscript{2} observations of Carn et al. (2011), and a source of HO\textsubscript{x} and NO\textsubscript{x} is suggested by observations of HO\textsubscript{2}NO\textsubscript{2} at Erebus (Oppenheimer et al., 2010), and elevated NO and NO\textsubscript{2} in plumes of Masaya (Mather et al., 2004) and Mt St Helens (see Martin et al., 2012, and references therein). Given the above-mentioned kinetic limitations to near-vent NO\textsubscript{x} production from entrained background air, these results imply the need for alternative explanations for NO\textsubscript{x} at volcanoes where it has been reported, and raise the possibility that volcanic NO\textsubscript{x} emissions at other volcanoes (e.g. Etna) might be lower than predicted by HSC.

A representation of high-temperature radical formation in the near-vent plume is, however, necessary for the initialisation of atmospheric chemistry models of downwind BrO chemistry. The HSC model output is thus used for this purpose, despite the above-mentioned limitations. Figure 3 shows 1 h PlumeChem model simulations for the three bromine emission scenarios (low, medium, high), initialised using HSC operated at a range of \(V_a/V_M\) varying from 0 : 100, 2 : 98, 5 : 95, 10 : 90 to 15 : 85, compared to reported BrO / SO\textsubscript{2} ratios from Oppenheimer et al. (2006) and Bobrowski et al. (2007). Simulations initialised with \(V_a/V_M\) of 0 : 100 (i.e. with no air mixed into the near-vent plume) under-predict BrO / SO\textsubscript{2} ratios compared to the observations, as has been shown previously (e.g. Bobrowski et al., 2007; Roberts et al., 2009; von Glasow, 2010) using atmospheric chemistry models. This is due to the low radical content at \(V_a/V_M = 0:100\) as shown in Fig. 2). Previous studies therefore chose HSC initialisations using \(V_a/V_M > 0:100\); e.g. Roberts et al. (2009) suggested \(V_a/V_M = 10:90\), and von Glasow (2010) suggested \(V_a/V_M = 15:85\). Given the revised location of the compositional discontinuity outlined above in Fig. 2, even lower \(V_a/V_M\), e.g. \(V_a/V_M = 2:98\) or \(V_a/V_M = 5:95\) (shown in red), can become suitable. Further progress will require more sophisticated models to be developed, e.g. to include full kinetic representations of chemical and mixing processes.

An interesting feature of Fig. 3 is that whilst choice of HSC initialisation affects the 1 h downwind plume BrO / SO\textsubscript{2} strongly, the model runs show a degree of convergence towards the end of the model run (particularly for low/medium Br cases). Understanding the < 1 h plume chemistry is important for interpretation of flank volcano BrO / SO\textsubscript{2} observations, and is investigated further with simulations initialised using HSC with \(V_a/V_M = 5:95\).
3.4 Speciation of reactive bromine in Etna plume and implications for observations of volcanic BrO

The evolution of reactive bromine speciation is also illustrated in Fig. 3 for the three bromine emission scenarios, with simulations initialised using HSC at $V_A : V_M = 5 : 95$. A number of interesting features are identified:

- BrO / Br_{tot} rises rapidly in the first few minutes, but then stabilises or declines further downwind.

- HBr is fully converted to reactive bromine in the “medium” and “low” bromine simulations but only partially converted in the “high” bromine simulations.

- BrO is formed in the plume at up to 40–50 % (“medium” and “low” Br emission scenarios) or 10–20 % (high Br emission scenario) of total bromine. This difference is related to the extent of HBr conversion, as BrO reaches a similar maximum fraction (∼ 50 %) of reactive bromine in the three simulations.

- An increase in plume BrCl occurs when HBr becomes depleted, which is due to the aqueous-phase equilibria producing substantial BrCI in place of Br2.

- HOBr and BrONO2 are present in all simulations, and represent an increasing proportion of reactive bromine as the plume disperses downwind, whilst the proportion of BrO declines.

The observed and modelled trend in BrO / SO2 shown in Figs. 1 and 3 is thus explained as follows: HBr is converted into reactive forms by autocatalytic bromine chemistry cycles involving volcanic aerosol, entrained atmospheric oxidants and sunlight. The HBr conversion is accelerated by radical species present in the high-temperature initialisation. The initial rise in BrO / SO2 primarily reflects trends in reactive bromine speciation; entrainment of background air containing ozone into the plume, promotes greater partitioning to BrO via the reaction Br + O3. Plume–air mixing is thus an important control on BrO / SO2, because the dilution of volcanic components and entrainment of air alter the balance between Br and BrO, e.g. by reducing the rate of BrO loss by the self-reaction BrO + BrO (to form 2Br or Br2), Reaction (R7), relative to the formation of BrO by Br + O2, Reaction (R6). The subsequent decline or plateau in BrO / SO2 occurs due to net conversion of reactive bromine from BrO to HOBr and BrONO2 in the downwind plume (Reactions R1, R2). These species are formed at an accelerated rate in the downwind plume as it disperses and entrains background air containing oxidants (HO2, NO2) which react with BrO. Further, the heterogeneous loss pathways for these species are slowed in the dispersed downwind plume where volcanic aerosol is diluted. The heterogeneous reactions of HOBr and BrONO2 with aerosol present a more rapid loss pathway than photolysis in the aerosol-rich environment of a volcanic plume. As the plume disperses and dilutes further downwind, net accumulation of HOBr (and BrONO2) occurs whilst BrO declines (as a fraction of Br_{tot}), although it is emphasised that plume chemistry cycling between these species is ongoing throughout the simulation and is very rapid.

This predicted reactive bromine evolution is somewhat similar to that of Roberts et al. (2009) but contrasts to the 1 h simulations of von Glasow (2010) that did not predict the in-plume presence of HOBr and BrONO2. The higher proportion of total bromine as BrO in the “medium” and “low” Br emission scenarios (40–50 %) compared to the “high” Br emission scenario (10–20 %) is related to the extent (complete and partial, respectively) of HBr conversion into reactive bromine species. This dependence of the HBr conversion on Br_{tot} / SO2 in the emission may to some extent explain differences between the model studies of Roberts et al. (2009) and von Glasow (2010) that predicted complete and partial
in-plume conversion of HBr into reactive bromine, respectively.

Predicted BrO/SO$_2$ and BrO/Br$_{tot}$ trends for the three Br emission scenarios (initialised with $V_A : V_M = 5 : 95$) are shown in Fig. 4. The “low” Br emission scenario simulation can be compared to observations by Oppenheimer et al. (2006) who reported (using DOAS) BrO/SO$_2$ ratios reached $\sim 2 \times 10^{-4}$ within 3–4 min downwind of Etna summit, and used filter packs to quantify the emitted Br$_{tot}$ to be $4.8 \times 10^{-4}$, implying a BrO/Br$_{tot}$ of $\sim 40\%$. For the “low” Br model run initialised at Br$_{tot}$/SO$_2 = 4.8 \times 10^{-4}$, which predicts complete conversion of HBr into reactive forms over 4 min, BrO/SO$_2$ rises to $10^{-4}$ within 4 min, reaching a maximum of $2.5 \times 10^{-4}$ at about 18 min downwind (i.e. earlier than the “medium” and “high” bromine cases of this study) after which BrO/SO$_2$ declines, Fig. 4a. BrO/Br$_{tot}$ reaches 25% within 4 min, and 40% by about 8 min (Fig. 4b), thus converging towards the observations of Oppenheimer et al. (2006): the agreement is relatively good considering that the predicted BrO/Br$_{tot}$ can also be affected by other model parameters kept constant here, e.g. in HSC initialisation, rate of plume–air mixing and aerosol loading, whilst observations of BrO/Br$_{tot}$ are subject to measurement uncertainties, e.g. in filter-pack Br/S, DOAS measurement of BrO/SO$_2$.

The non-linearity of HBr conversion to BrO shown in Fig. 4 yields the following implications for volcanology: BrO/SO$_2$ ratios for these simulations (initialised at $V_A : V_M = 5 : 95$) reach maxima of $3.6 \times 10^{-4}$ and $4.6 \times 10^{-4}$ and $2.5 \times 10^{-4}$ for the medium, high and low Br scenarios respectively in the downwind plume. Thus, whilst the modelled bromine emission has varied by a factor of three between the “medium” and “high” bromine scenarios, the simulated BrO/Br$_{tot}$ ratio has varied by less than 30%. This result for small-scale bromine variations contrasts to the earlier PlumeChem simulations (Roberts et al., 2009) which suggested that order of magnitude differences in BrO/Br$_{tot}$ between Soufrière Hills volcano (BrO/Br$_{tot}$ $\sim 10^{-3}$) and other Arc volcanoes like Etna (BrO/Br$_{tot}$ $\sim 10^{-3}$) could be attributed to order of magnitude differences in the ratio of total bromine to SO$_2$ in their emissions. However, the non-linear relationship between BrO and emitted HBr as identified in Fig. 4 for small-scale bromine variations, presents a complexity to efforts to quantify volcanic bromine emissions using DOAS observations of plume BrO/Br$_{tot}$ ratios. Further understanding of the extent to which volcanic bromine is rapidly converted into reactive forms in the near-downwind plume is needed as part of efforts to evaluate global impacts from volcanic halogen degassing. Further studies of the wider model parameter space can contribute to this aim, although more observations are also needed to constrain model uncertainty. Overall, the model suggests that HBr conversion into reactive bromine depends on a balance between the autocatalytic “bromine explosion” cycles in the near-downwind plume (accelerated by radicals produced in the high-temperature near-vent plume), and the conversion of reactive bromine back into HBr (e.g. via the reaction Br + HCHO).
3.5 Low in-plume prevalence of BrNO₂

Formation of BrNO₂ from Br + NO₂ was excluded from the 1 h simulations presented in Figs. 3 and 4. However, the plume chemistry modelling study of von Glasow (2010) predicted high in-plume prevalence of BrNO₂, due to reaction of Br with NO₂, given that high Br and NO₂ mixing ratios are assumed in the (HSC) model initialisation. In the Etna simulations of von Glasow (2010), formation of BrNO₂ exceeds its photolytic loss rate in the young plume, leading to a significant partitioning (>30 %) of plume bromine as BrNO₂. To further evaluate this model difference, a similar two-reaction scheme for BrNO₂ was introduced into the PlumeChem model, with BrNO₂ the assumed (sole) product of the reaction Br + NO₂. With this two-reaction scheme, model runs for the three bromine scenarios also show rapid formation of BrNO₂ (Fig. 5a, c, d). The in-plume BrNO₂ prevalence (<30 % of plume bromine declining to just a few percent after 30 min), is still somewhat less than that of von Glasow (2010), and model differences remain in Br speciation regarding presence of HOBr and BrONO₂, potentially due to differences between the models’ aerosol loading or dispersion schemes. Figure 5a, c, d highlight that the rapid formation of BrNO₂ in these model runs causes a slight delay in the formation of BrO downwind compared to the standard model runs of Fig. 3.

However, we do not recommend use of the two-reaction BrNO₂ scheme, because the chemistry is in fact more complex. Firstly, the reaction Br + NO₂ primarily produces BrONO (~92%) rather than BrNO₂ (~8%) (Bröske and Zabel, 1998; Orlando and Burkholder, 2000). Secondly, BrONO undergoes a more rapid thermal dissociation (τ ~ 1 s at room temperature), and photolytic loss (τ ~ seconds) than BrNO₂ (Burkholder and Orlando, 2000). BrONO and BrNO₂ also react with NO (Bröske and Zabel, 1998). BrONO (and possibly also BrNO₂) also react with Br radicals. The reactions are summarised in Table 3. PlumeChem simulations using a more detailed reaction scheme for BrNO₂–BrONO–BrNO, incorporating the quantified reactions of Table 3, are illustrated in Fig. 5b, d, f. With this revised BrNO₂–BrONO–BrNO model scheme, these species account for only <12 % of reactive bromine (with BrONO and BrNO at only <1 %). The impact of this scheme on Br speciation is rather modest but some differences can be seen in comparison to the “standard” simulations of Fig. 3 – for example a slightly faster rate of HBr conversion to reactive bromine. However, this more detailed reaction scheme is itself limited in that it does not include reaction of BrNO₂ with Br (rate constant unknown), and assumes that the two possible BrONO photolysis pathways occur equally (as products are unknown). Further, the scheme does not include potential heterogeneous reactions relevant for BrNO₂. Heterogeneous reactive uptake of N₂O₄ might produce BrNO₂ or ClNO₂. However, these products might react further within the aerosol to form Br₃ or BrCl (Frenzel et al., 1998). Proper investigation of such heterogeneous chemistry on volcanic aerosol would require detailed consideration of the underlying rate constants for all the aqueous-phase reactions (e.g. in a manner similar to that recently attempted for HOBr reactive uptake, Roberts et al., 2014). In addition to uncertainty in the model chemistry, the model findings are also subject to uncertainty in the HSC initialisation (which determines the volcanic Br and NO₂ radical source), see Sect. 3.3. Nevertheless, the more detailed reaction BrNO₂–BrONO–BrNO scheme findings suggest that the influence of BrNO₂ on the plume chemistry is much lower than that proposed by von Glasow (2010). Further simulations of this study therefore do not include BrNO₂.

3.6 Influence of plume–air mixing on BrO formation and ozone depletion

Here we investigate the role of plume–air mixing on the (low-temperature) halogen chemistry evolution of the downwind plume. A first study investigates small variations as might be expected on a day-to-day basis at Etna. A second study investigates how large variations in the volcanic emission flux...
Table 3. List of gas-phase and photolytic reactions related to formation of BrNO₂, BrONO and BrNO. Reactions listed are used in the – BrONO–BrNO scheme. The two-reaction BrNO₂ scheme assumes BrNO₂ as the sole product from both Br + NO₂ reactions and photolysis of BrNO₂ as the only loss pathway. See text for discussion of possible additional heterogeneous pathways.

<table>
<thead>
<tr>
<th>Reaction</th>
<th>Rate coefficient</th>
<th>at 285 K</th>
</tr>
</thead>
<tbody>
<tr>
<td>Br + NO₂ → BrNO₂</td>
<td>~3.8 × 10⁻³ cm³ molec⁻¹ s⁻¹</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td>Br + NO₂ → BrONO</td>
<td>~4.8 × 10⁻⁻² cm³ molec⁻¹ s⁻¹</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td>BrONO + Br → Br₂ + NO₂</td>
<td>2.4 × 10⁻¹¹ cm³ molec⁻¹ s⁻¹</td>
<td>Mellouki et al. (1989)</td>
</tr>
<tr>
<td>BrONO + NO₂ → BrNO₂ + NO₂</td>
<td>~2 × 10⁻¹⁶ cm³ molec⁻¹ s⁻¹ (uncertain)</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td>BrONO → BrNO₂</td>
<td>~1.2 s⁻¹ (at 298 K, 1 atm)</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td></td>
<td>τ &lt; 1 s at 298 K</td>
<td>Orlando and Burkholder (2000)</td>
</tr>
<tr>
<td>BrONO → BrNO₂</td>
<td>unknown</td>
<td>–</td>
</tr>
<tr>
<td>BrNO₂ + Br → Br₂ + NO₂</td>
<td>unknown</td>
<td>–</td>
</tr>
<tr>
<td>BrNO₂ + NO → BrNO₂ + NO₂</td>
<td>2.3 × 10⁻¹² Exp[-17.8/RT] cm³ molec⁻¹ s⁻¹</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td>BrONO + NO → BrNO₂ + NO₂</td>
<td>unknown, larger than BrNO₂ equivalent</td>
<td>–</td>
</tr>
<tr>
<td>BrNO + Br → Br₂ + NO</td>
<td>3.7 × 10⁻¹⁰ cm³ molec⁻¹ s⁻¹ or: 5.2 × 10⁻¹² cm³ molec⁻¹ s⁻¹</td>
<td>Grimley and Houston (1980)</td>
</tr>
<tr>
<td>BrNO₂ → BrNO₂</td>
<td>≤ 4.0 × 10⁻⁴ s⁻¹</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td>2BrNO₂ → Br₂ + 2NO₂</td>
<td>~6.4 × 10⁻⁵ s⁻¹</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td>2BrNO₂ + hν → Br₂ + 2NO₂</td>
<td>Unknown (slow)</td>
<td>Brökse and Zabel (1998)</td>
</tr>
<tr>
<td>BrONO + hν → Br₂ + NO</td>
<td>τ ~ s (products unknown)</td>
<td>Burkholder and Orlando (2000)</td>
</tr>
<tr>
<td>BrONO + hν → BrO + NO</td>
<td>or: τ ~ s (products unknown)</td>
<td>–</td>
</tr>
<tr>
<td>BrNO₂ + hν → Br + NO₂</td>
<td>τ ~ min</td>
<td>Scheffler et al. (1997)</td>
</tr>
</tbody>
</table>

(e.g. due to an eruption) influence the plume chemistry, albeit within the limitations of an idealised model scenario.

3.6.1 Influence of plume dispersion parameters, volcanic emission flux and wind speed on BrO/SO₂

As already discussed in Sect. 3.4, BrO formation is initially promoted by the entrainment of background air (containing ozone, HO₅ and NO₂), due to the balance between the reaction Br + O₃ (Reaction R6) and the self-reaction of BrO (Reaction R7), but as the plume becomes more diluted the entrainment of air acts to reduce BrO/SO₂ due to the reaction of BrO with HO₂ and NO₂ (Reactions R1, R2). Thus, the proportion of background air that has been entrained into the plume acts as a key control on BrO/SO₂. In the single-box Gaussian plume dispersion model used here, the extent of mixing of the background air into the plume is controlled by the Pasquill–Gifford dispersion parameters as a function of distance downwind, whose choice depends on atmospheric turbulence (a function of wind speed and atmospheric stability). Further, for a given dispersion parameterisation, the extent of mixing depends inversely on the volcanic emission flux, and also depends on wind speed (through dilution along the plume). Here the effects of these three variables are illustrated for a range of plausible volcanic and meteorological conditions at Etna in Fig. 6.

For the base-run simulations (also shown in Figs. 3 and 4), a Pasquill–Gifford (PG) dispersion case D was used, that is for a relatively neutral atmosphere, with a wind speed of 10 m s⁻¹ and volcanic gas flux of 10 kg s⁻¹ SO₂ (with the emission of all other volcanic gas and aerosol components scaled accordingly). This SO₂ flux estimate is close to the ~ 13 kg s⁻¹ reported by McGonigle et al. (2005) for 30 July 2004. The model 10 kg s⁻¹ SO₂ flux is, however, a somewhat low representation for Mt Etna during 2004–2005 in general. Aiuppa et al. (2005) report gas flux data that show summertime variations between 800 and 2000, equivalent to 9–23 kg s⁻¹ SO₂, with even greater SO₂ flux during eruption periods. Burton et al. (2005) report 7-day average SO₂ fluxes of 1000–2500 t d⁻¹ (12–25 kg s⁻¹). To illustrate the influence of variation gas flux and plume dispersion, simulations were also performed at 20 kg s⁻¹ SO₂ flux, and for a range of dispersion and wind-speed cases. Cases C and B are introduced for more unstable atmospheric conditions involving enhanced plume–air mixing, which occur more readily at lower wind speed (< 6 m s⁻¹) (see Supplement).

Simulations performed at wind speeds of 10 m s⁻¹ (case D and C), 15 m s⁻¹ (case D and C), 5 m s⁻¹ (case D and C), and 3 m s⁻¹ (case C and B) are shown in Fig. 6 (a “medium” bromine scenario is assumed for all these simulations, with V₄₆/V₅₄ = 5.95 in the initialisation). The model runs illustrate how plume–air mixing may cause variation in the downwind BrO/SO₂. The variation is of a similar magnitude to that identified in the model runs with the three bromine scenarios, Fig. 4 (which themselves encompass only a portion of the reported variability in Br/S in the emission, see Aiuppa et al., 2005). The model runs suggest that a combination of...
variations in plume–air mixing and bromine emission could provide – at least theoretically – a variability in BrO/SO$_2$ similar to the observed variability in BrO/SO$_2$ (5. × $10^{-3}$ – $3.9 \times 10^{-4}$) reported by Bobrowski and Guiffrida (2012) at 6 km downwind. Variability in the volcanic aerosol emission could potentially add further to this.

Plume dispersion causes a transition between the two chemical regimes outlined above and an intermediate maximum in BrO/SO$_2$. The magnitude and location of the downwind maximum in BrO/SO$_2$ depends on the extent of plume–air mixing, as determined by the gas flux, rate of dispersion and wind speed, as well as on the volcanic aerosol loading and bromine content, and the HSC initialisation. Variations in background atmospheric composition (e.g. ozone, HO$_x$, NO$_x$, aerosol) could further modify the results. Finally, if applying these results to volcanoes elsewhere, the summit altitude is also a relevant consideration, as the greater atmospheric density at lower altitude will yield a higher in-plume ratio of background oxidants to bromine, for a given volcanic SO$_2$ flux.

Nevertheless, large increases in the volcanic emission flux tend to maintain for longer the more “concentrated” regime where BrO/SO$_2$ is limited by the balance between Reactions (R6) and (R7), as discussed further below.

3.6.2 Effect of a large increase in volcanic flux on BrO/SO$_2$

The sensitivity study is continued for high-emission scenarios by keeping the plume dispersion case and bromine emission scenario constant (case D, “medium” Br$_{tot}$/SO$_2$), but increasing the volcanic gas and aerosol emission (by a factor of × 5 and × 10 of the base-run 10 kg s$^{-1}$ SO$_2$ flux). Such an increased volcanic emission maintains higher concentrations of volcanic gases and thus reduces the extent of plume–air mixing, and hence entrainment of background oxidants into the plume. We caution that in a real volcanic environment, such a large change in degassing rate may also be accompanied by a change in composition of the volcanic emission (including halogen content or aerosol loading) or act to alter the plume dimensions somewhat (e.g. by the dynamics of explosive eruptions). The model results here focus solely on the effect of (substantially) enhanced gas flux with all other variables held constant.

Simulations of 3 h duration (equivalent to 108 km downwind plume propagation assuming 10 m s$^{-1}$ wind speed) with volcanic emission flux increased from the base run to × 5 and × 10 are shown in Fig. 7, for both the “medium” and “high” bromine emission scenarios (initialised with HSC using $V_A$: $V_M$ = 5:95). The enhanced volcanic emission flux linearly enhances in-plume SO$_2$ abundance, as expected, but exerts a non-linear effect on the plume chemistry and impacts.

In particular, the greater volcanic emission (lower plume–air mixing) leads to a slower rise, and a later onset and slower decline in BrO/SO$_2$. At distances far downwind (> 2 h for the specific simulation conditions), high BrO/SO$_2$ is sustained for longer in plumes with high gas flux. Conversely, in the near downwind (several 10s of min), plumes with lowest gas flux exhibit the fastest initial rise and highest BrO/SO$_2$ ratios. As described above, these model findings are readily explained by the model chemistry that partitions reactive bromine between Br and BrO (during the initial rise), and BrO, and HOBr, BrONO$_2$ (during the subsequent decline) as the plume disperses. The onset and magnitude of the decline is greatest for low-flux plumes that are more dilute and where a higher proportion of background air has been mixed into the plume. Conversely, high-flux and thus more concentrated plumes have a slower initial increase in BrO/SO$_2$, with a delayed maximum. In the relatively near-downwind plume (0–30 min), the model predicts lower BrO/SO$_2$ at greater volcanic gas fluxes, as shown by the arrows in Fig. 6. Implications for the interpretation of volcano plume observations are discussed in Sect. 3.7.

3.6.3 Atmospheric impacts of volcanic reactive halogen chemistry

BrO chemistry causes ozone, HO$_x$ and NO$_x$ to become depleted in the downwind plume, Fig. 8. For HO$_x$ and NO$_x$
Figure 7. Simulated plume SO$_2$, BrO and BrO/SO$_2$ over 3 h for the medium- and high-bromine emission scenarios, and with varying volcanic emission flux (baseline run, and with volcanic gas + aerosol emissions flux $\times$ 5 and $\times$ 10, shown by full-, long-dashed and short-dashed lines, respectively), whilst keeping the same plume dispersion parameterisation, wind speed and initialisation (see text for model details and interpretation). Arrows highlight the reduction in near-downwind BrO/SO$_2$ predicted at greater volcanic emission flux.

The near-downwind plume abundances are initially elevated as the HSC initialisations used assumed a volcanic source of these species (Fig. 2), but become depleted within a few to 10s of minutes downwind. The maximum depletion reaches is near 100 and $>70\%$ depletion relative to background values of around 30 and 0.17 ppbv for HO$_x$ and NO$_x$ respectively. HO$_x$ is converted into H$_2$O$_{2(g)}$ via HOB$_r$ chemistry (Reactions R1, R3). HO$_x$ abundances are also reduced by the gas-phase reaction of OH with SO$_2$, and by ozone depletion in the plume (see below). The volcanic NO$_x$ source is converted into HNO$_3$ by BrONO$_2$ chemistry (Reactions R2, R4), causing a rapid increase of in-plume HNO$_3$, particularly in the concentrated near-downwind plume, where HNO$_3$ reaches up to 60 nmol mol$^{-1}$ (exceeding the background NO$_x$ of $\sim 6$ nmol mol$^{-1}$). This mechanism was proposed by Roberts et al. (2009) as an explanation for observations of high HNO$_3$ in volcanic plumes. See collated observations by Martin et al. (2012) reporting plume HNO$_3$ / SO$_2$ that can reach up to $10^{-1}$. For Etna in particular, reported crater-rim HNO$_3$ / SO$_2$ ratios are somewhat inconsistent and show large variability ($-2.3 \times 10^{-4}$, $7.8 \times 10^{-6}$, $4.2 \times 10^{-3}$), which in itself might be indicative of a role of plume chemistry processing. Recently Voigt et al. (2014) also observed elevated HNO$_3$ in the downwind Etna plume, with HNO$_3$ as the dominant form of NO$_x$. Importantly, elevated “volcanic” HNO$_3$ produced by the BrONO$_2$ mechanism can originate from both NO$_x$ of volcanic origin, and from NO$_x$ from background air entrained into the plume. As a consequence, the in-plume NO$_x$ declines from initially elevated abundance (due to the assumed high-temperature volcanic NO$_x$ source) to become depleted relative to the background abundance downwind. Finally, it is noted that simple acidification of nitrate aerosol from background air entrained into the plume could also lead to gas partitioning and therefore enhance the “volcanic” HNO$_3(g)$ signature. Such acid displacement of HNO$_3(g)$ by H$_2$SO$_4(aq)$ has been observed by Satsumabayashi et al. (2004). The observations of volcanic HNO$_3$ collated by Martin et al. (2012) and Voigt et al. (2014) thus require consideration in the context of these two mechanisms.

Ozone is also depleted in the plume and reaches a maximum depletion (up to 100 %) around 10 min downwind, coincident with the highest in situ BrO abundances that reach $\sim 1$ nmol mol$^{-1}$ (Fig. 7). For the base run, the maximum plume ozone depletion is 30 or 45 nmol mol$^{-1}$ for the medium- and high-bromine emission scenarios respectively. Greater in-plume ozone loss occurs at higher emissions flux (lower relative plume–air mixing). However, for these runs the maximum ozone loss is constrained by the fact that it cannot exceed $\sim 60$ nmol mol$^{-1}$ (the background ozone mixing ratio). Thereafter ozone begins to recover as the plume disperses (Fig. 8), entraining background air, and BrO declines (Fig. 7), albeit at a slower rate than the SO$_2$ plume tracer. Ozone recovery is greater for the base run than the higher volcanic flux cases due to both physical and chemical consequences of enhanced plume–air mixing. Thus the presence of a detectable ozone depletion signature at distances far downwind depends on the emission flux and plume dispersion. Further, the single-box simulations presented here that predict the downwind trend do not simulate the ozone distribution across the plume cross section. Ozone loss is typically greater in the plume centre than near the edges, as demonstrated by both observations and spatially resolved PlumeChem simulations of Redoubt volcano plume (Kelly et al., 2013). The single-box simulations should be interpreted in this context, e.g. a predicted loss of 45 nmol mol$^{-1}$ implies greater loss at the plume centre (likely close to 60 nmol mol$^{-1}$ or 100 %) declining to near-ambient ozone at the plume edges. With distance (time) downwind, the ozone mixing ratio starts to increase when the entrainment of ambient air containing O$_3$ is faster than...
the local O₃ destruction (Fig. 8). Nevertheless, ongoing occurrence of ozone-depleting BrO chemistry is demonstrated by the continuing negative trend in the cumulative ozone loss: the ozone difference (plume – background) integrated across the plume cross-sectional area declines along the 3 h simulations to reach ∼1, 4, and 7 g cm⁻¹ for the three flux scenarios (SO₂ flux = 10, 50, 100 kg s⁻¹) respectively with greater ozone loss for the high Br compared to the medium Br scenario, as expected. These Lagrangian simulations of plume “puff” ozone evolution over 3 h can also be viewed in an Eulerian context: the 3 h impact of continuous volcano emissions is calculated by integrating the cross-sectional impact (g cm⁻¹) over the distance downwind. This yields ozone losses of 35 × 10³ (38 × 10³), 23 × 10³ (26 × 10³), and 4 × 10³ (6 × 10³) kg for the × 10 flux, × 5 flux and base-run (10 kg s⁻¹ SO₂ flux) scenarios respectively, assuming the medium Br scenario (numbers in brackets refer to high Br scenario). Whilst there is some linearity in ozone loss per Br emitted (e.g. in comparing the base run to × 5 flux cases), the constraint that ozone loss cannot exceed 100 % of the background abundance introduces some non-linearity for the × 10 flux case, thereby reducing its overall ozone loss. Note that the plume cross-sectional area after 3 h is π × √2 × σ_h × √2 × σ_z = 2 × π × 4470 × 485 = 1.4 × 10⁵ m². The volcanic plume cone thus resides within a cylinder of volume 1.4 × 10⁵ × 108 × 10³ = 1.5 × 10¹² m³, containing approx. 10⁵ kg ozone.

Figure 8 indicates that the plume atmospheric impacts extend beyond the 1 to 3 h simulations presented in this study. Simulations over the lifetime of volcanic plumes under different volcanological and meteorological conditions are therefore required to quantify the global tropospheric impact from volcanic halogen emissions.

3.7 Implications for modelling and observations of volcanic BrO

The parameter space governing volcanic plume reactive halogen chemistry is vast, and is not fully constrained by available observations. Of particular importance in controlling the reactive bromine formation and downwind plume bromine speciation are: Br₉/ SO₂ in the emission, the volcanic aerosol loading, and the extent of background air mixing into the plume (itself a function of the plume dispersion parameterisation, volcanic emission flux and wind speed). These factors exert non-linear influences on the conversion of emitted HBr into plume reactive bromine, and its speciation through interconversion of BrO, Br₂, BrCl, HOBr, BrONO₂. The onset of the autocatalytic reactive bromine formation is also accelerated in the model by radicals in the high-temperature model initialisation (Br, Cl, NOₓ, HOₓ). A major area of uncertainty is, however, the representation of this high-temperature near-vent plume environment using thermodynamic models such as HSC. Development of
high-temperature kinetic models of the near-vent plume is encouraged for progress in this area.

Further uncertainty to the downwind plume chemistry is contributed by uncertainty in the volcanic bromine emission, and in aerosol surface area, that sustains halogen cycling downwind. Crater-rim filter-pack measurements (e.g. Aiuppa et al., 2005) provide estimates of volcanic Br/S emissions for model initialisation (see Table 2) but also highlight temporal variability in this parameter. The volcanic aerosol emission is poorly constrained by observations at Etna, and from volcanoes globally. A surface area loading of \( \sim 10^{-11} \text{m}^2 / \text{molec SO}_2^+ \), i.e. an order of magnitude lower than that used by Roberts et al. (2009), yields simulated (0–20 km) downwind BrO/\text{SO}_2 more consistent with that observed in the Etna plume. Volcanic aerosol has a small influence on BrO/\text{SO}_2 ratio near source, but is an important control in the more dispersed plume downwind. Uncertainties in the volcanic aerosol emission magnitude, and its size distribution (which for sulfate varies as a function of temperature and humidity) thus contribute to uncertainties in models of the plume halogen chemistry. Plume aerosol may be augmented by in-plume oxidation of volcanic SO\(_2\) to \( \text{H}_2\text{SO}_4\), and the entrainment and acidification of background aerosol may also promote halogen cycling. Future model evaluation of volcanic reactive halogen impacts in the wider troposphere will require the development of regional and global models, with detailed treatment of aerosol processes as well as plume dispersion (shown to be a key control on the downwind chemistry). An improved quantification of the kinetics of HOB\(_2\) reactive uptake on volcanic aerosol is also needed according to Roberts et al. (2014). Global models may need to include a representation of the subgrid-scale volcanic plume processes, particularly as this study has highlighted how the proportion of emitted HBr converted into reactive forms is non-linearly dependent on the degassing scenario.

We emphasise the complex role of plume chemistry in the interpretation of volcanic flank DOAS measurements of BrO/\text{SO}_2. Bobrowski and Giuffrida (2012) recently reported variation in BrO/\text{SO}_2 ratios at Etna related to the onset of eruption activity, for example with increasing BrO/\text{SO}_2 shortly prior to an eruptive event, and lower BrO/\text{SO}_2 during the eruption event, according to DOAS measurements 6 km downwind from the summit. These observations have been interpreted in the context of variable bromine and SO\(_2\) emissions, related to subsurface magmatic processes. Lübcke et al. (2014) identified a decrease in BrO/\text{SO}_2 observed using a DOAS instrument prior to an eruption event at Nevado del Ruiz, Colombia (in a period whilst SO\(_2\) emissions were increasing). However, we emphasise that a variation in plume BrO/\text{SO}_2 can also result from differences in the plume chemistry for varying volcanic emission flux magnitudes. Figure 6 shows that changes in volcanic gas flux (for a fixed plume dimension) can yield substantial changes in plume BrO/\text{SO}_2 ratio, even for a fixed Br/\text{SO}_2 ratio in the emission. In the near-downwind plume, a key control on BrO formation is the entrainment of oxidants. A substantial increase in volcanic emission flux leads to greater plume strength and reduced ratio of background oxidants to bromine in the model. Thus, on the < 60 min timescale of volcano flank DOAS observations, a substantially enhanced rate of volcanic degassing generally leads to lower plume BrO/\text{SO}_2 ratios in more concentrated plumes. Potentially, the variations in BrO/\text{SO}_2 identified by Bobrowski and Giuffrida (2012), and Lübcke et al. (2014) may result from a combination of volcanological and plume chemistry factors. This example highlights the complexity surrounding the interpretation of volcanic BrO and shows the role of plume chemistry modelling in the effort to use volcanic BrO observations to monitor and predict volcanic activity.

We also highlight that the plume chemical evolution causes a decline in BrO/\text{SO}_2 ratios in the dispersed plume further downwind through net conversion of BrO into reservoirs such as HOB\(_2\) and BrONO\(_2\). This plume chemical evolution acts to reduce the BrO column abundance, contributing additional limitations to its possible detection in dispersed plumes, and is the model explanation for the plateau in BrO/\text{SO}_2 downwind of Etna reported by Bobrowski and Giuffrida (2012). Detection of volcanic BrO by satellite is primarily constrained to large volcanic emissions (Theys et al., 2009; Rix et al., 2012; Hörmann et al., 2013). Smaller volcanic emissions that generate high but localised BrO at lower altitudes are less readily detected, particularly due to dilution effects across the satellite measurement pixel (Afe et al., 2004). The modelled plume chemical evolution adds to this limitation for satellite detection of BrO in dispersed volcanic plumes (even at higher resolution). Importantly, however, the model Br speciation shows that a declining trend in BrO abundance as the volcanic plume disperses does not preclude the occurrence of continued in-plume reactive bromine chemistry as predicted by the model.

4 Conclusion

We present a PlumeChem model study of the reactive halogen chemistry of Mt Etna volcano plume that reproduces the recently reported trends in BrO/\text{SO}_2; namely a rapid increase in the near-downwind followed by stability or decline in the far-downwind plume. A new in-plume evolution of Br speciation is predicted: BrO Br\(_2\), Br and HBr are the main plume species in the near-downwind plume whilst BrO, HOB\(_2\) (and BrONO\(_2\), BrCl) are present in significant quantities further downwind. An evaluation of the (quantifiable) chemistry surrounding BrNO\(_2\) suggests a rather low prevalence in volcanic plumes, although uncertainties in model chemistry and initialisation are highlighted.

Emitted volcanic HBr is converted into reactive bromine by autocatalytic bromine chemistry cycles whose onset is accelerated by the model high-temperature initialisation. The initial rise in BrO/\text{SO}_2 is primarily due to entrainment of
ozone through plume dispersion that promotes BrO formation from Br radicals. A subsequent decline or plateau in BrO/SO2 occurs upon plume dispersion, which both dilutes the volcanic aerosol (slowing HOBr and BrONO2 heterogeneous loss rates) and entrains HO2 and NO2 from the background atmosphere (promoting HOBr and BrONO2 formation from BrO). This promotes net accumulation of reservoirs HOBr and BrONO2 and a reduction in BrO in the dispersed downwind plume. Thus the model can explain the reported BrO/SO2 trend at Etna. We demonstrate the role of plume chemistry models to interpret volcanic BrO/SO2 observations as well as quantify atmospheric impacts on HOx, NOx, HNO3 and ozone. A number of volcanological and meteorological factors can influence plume BrO/SO2 ratios, and we illustrate simulations with contrasting total bromine content and volcanic aerosol loading. The influence of plume–air mixing is shown by simulations with varying dispersion rate, as well as wind speed and volcanic gas flux.

BrO contents reach up to 20 and ~50% of total bromine (over a timescale of a few 10s of minutes), for the high- and medium-/low-bromine emission scenarios, respectively. The latter agrees well with observations that report BrO (at 3–5 min downwind) can reach up to 40% of the total bromine emission at Etna (Oppenheimer et al., 2006). Partial (up to ~50%) or complete (100%) conversion of HBr to reactive forms is predicted over the 1 h simulations, depending on bromine content (high, medium or low, respectively) as well as other the plume conditions (e.g. aerosol, dispersion, HSC initialisation). Simulations using the two volcanic aerosol loadings significantly differ in the downwind plume chemistry but result in a similar initial rise in BrO/SO2 near-downwind (up to 6 km), a finding that is in agreement with the reported low relative humidity dependence of BrO/SO2 (Bobrowski and Giuffrida, 2012).

Simulations with a fixed dispersion rate but enhanced volcanic emission flux are presented. For higher emission fluxes, the stronger plume and reduced ratio of background oxidants to bromine causes a slower rise in BrO/SO2 in the near-downwind plume (<40 min) and a slower and delayed onset of the decrease in BrO/SO2 in the far-downwind plume (>2 h, for the volcanic conditions simulated). This simulated dependence of BrO/SO2 on volcanic emission flux (albeit in an idealised model scenario) is particularly relevant for the interpretation of changes in BrO/SO2 during/prior to eruptive events (e.g. Bobrowski and Giuffrida, 2012; Lübcke et al., 2014).

Impacts of the plume halogen chemistry include downwind depletion of HOx, NOx and ozone, and formation of HNO3. Partial recovery of ozone is predicted, particularly for low gas flux emissions. However, cumulative impacts on ozone are ongoing over the 3 h simulations.

Author contributions. T. J. Roberts, designed and performed the PlumeChem model experiments and HSC calculations and wrote the manuscript. R. S. Martin advised on HSC methodology and contributed to manuscript writing. L. Jourdain advised on scientific scope and contributed to manuscript writing.

Acknowledgements. This study was financed by LABEX VOLTAIRE (VOLatils-Terre Atmosphère Interactions – Ressources et Environnement) ANR-10-LABX-100-01 (2011–20) and an NSINK career development allowance that enabled HSC hardware purchase. R. S. Martin acknowledges Christ’s College, Cambridge for a research fellowship. We are grateful to R. Sander, C. Kern and an anonymous reviewer, whose feedback helped to improve the manuscript.

Edited by: W. Birmili

References


T. J. Roberts et al.: Reactive bromine chemistry in Mount Etna’s volcanic plume


