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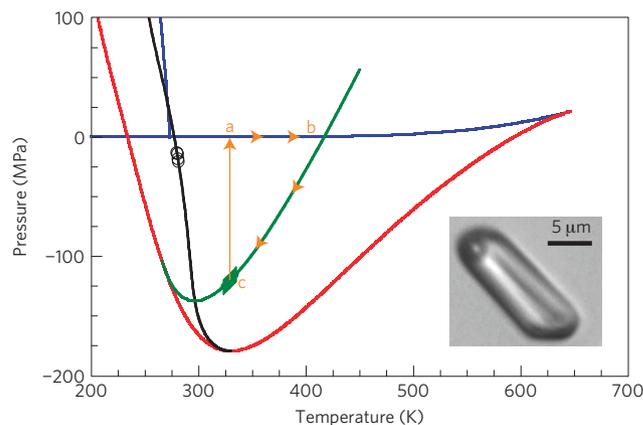
# A coherent picture of water at extreme negative pressure

Mouna El Mekki Azouzi<sup>1</sup>, Claire Ramboz<sup>2</sup>, Jean-François Lenain<sup>3</sup> and Frédéric Caupin<sup>1\*</sup>

Liquid water at atmospheric pressure can be supercooled to  $-41^\circ\text{C}$  (ref. 1) and superheated to  $+302^\circ\text{C}$  (ref. 2). Experiments involving fluid inclusions of water in quartz suggest that water is capable of sustaining pressures as low as  $-140\text{ MPa}$  before it breaks by cavitation<sup>3</sup>. Other techniques, for which cavitation occurs consistently at around  $-30\text{ MPa}$  (ref. 4), produce results that cast doubt on this claim. Here we reproduce the fluid-inclusion experiment, performing repeated measurements on a single sample—a method used in meteorology<sup>5</sup>, bioprotection<sup>6</sup> and protein crystallization<sup>7</sup>, but not yet in liquid water under large mechanical tension. The resulting cavitation statistics are characteristic of a thermally activated process, and both the free energy and the volume of the critical bubble are well described by classical nucleation theory when the surface tension is reduced by less than 10%, consistent with homogeneous cavitation. The line of density maxima of water at negative pressure is found to reach  $922.8\text{ kg m}^{-3}$  at around  $300\text{ K}$ , which further constrains its contested phase diagram.

Owing to the strong cohesion of water, which manifests in its large surface tension, the liquid is expected to withstand pressures in excess of  $-100\text{ MPa}$  (ref. 8), as predicted for instance by classical nucleation theory<sup>9</sup> (CNT). The only experimental technique for which such large tensions have been reported uses water inclusions in quartz<sup>3</sup>. The magnitude of the tension was also independently confirmed by light scattering<sup>10</sup>. Several other independent experimental techniques report cavitation at much lower tensions<sup>4</sup>, prompting the proposal of a new kind of heterogeneous cavitation involving stabilization of water in an inclusion against nucleation by impurities<sup>11</sup> or surfaces<sup>12</sup>. Agreement between cavitation-pressure measurements and CNT would thus be advantageous.

In the work that pioneered the fluid-inclusions technique<sup>3</sup>, a large number of inclusions, covering a wide density range, were studied. However, the cavitation event in each individual inclusion was usually observed only once. One inclusion was observed in repeated runs “to nucleate randomly in the range  $40\text{--}47^\circ\text{C}$  and occasionally not at all”<sup>3</sup>, which was qualitatively interpreted as evidence for the crossing of the line of density maxima (LDM) of water. In the present work, we chose to focus on only one inclusion, and to perform a large number of cavitation experiments to obtain the statistics of nucleation. We confirm the previous data in terms of cavitation threshold, but in addition, by probing nucleation rates over three decades and using the nucleation theorem<sup>13</sup>, we gain insight into the nanoscopic mechanism underlying bubble nucleation. The results are consistent with CNT with a slightly



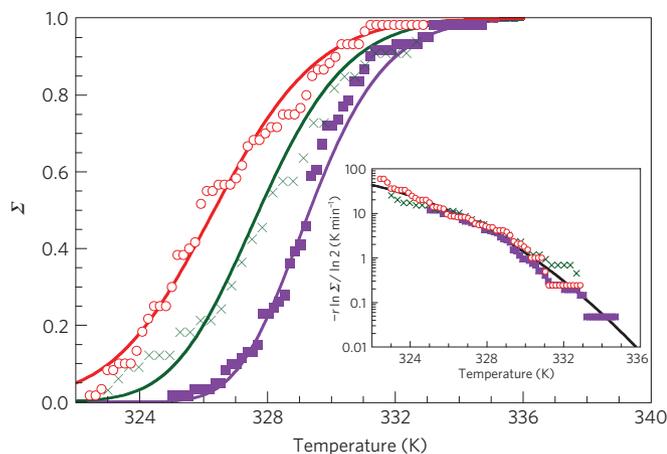
**Figure 1** | Schematic path followed by water in a quartz inclusion.

Equilibrium transitions are given by blue curves. The extrapolation of the EOS measured at positive pressure<sup>21</sup> is used to plot the liquid–vapour spinodal (red), the LDM (black) and the isochore at  $\rho = 922.8\text{ kg m}^{-3}$  (green) corresponding to the inclusion studied (see Methods for details). Here the spinodal is re-entrant<sup>14</sup>, but other scenarios suggest that it remains monotonic<sup>9,15</sup>. The LDM has been measured only down to  $-20.3\text{ MPa}$  at  $280.84\text{ K}$  (ref. 16; black circles). The inclusion studied here (volume  $V = 570\text{ }\mu\text{m}^3$ ) is shown as an inset, with a bubble in the top left corner. The thick green section shows the series of 154  $T_{\text{cav}}$  values measured in the present study; they correspond to a cavitation pressure  $P_{\text{cav}}$  around  $-120\text{ MPa}$ .

reduced surface tension, related to the small size of the critical bubble. Moreover, the analysis allows us to derive quantitative information on the LDM, which is of paramount importance in the debate about the origin of water anomalies. Indeed, if the LDM sustains a negative slope (Fig. 1), it will intersect the liquid–vapour spinodal and bend it to lower tension at lower temperature<sup>14</sup>; whereas if the LDM changes slope, bending to lower temperature at higher tension, the spinodal remains monotonic<sup>15</sup>. Up to now, the LDM had been measured only down to  $-20.3\text{ MPa}$  at  $280.84\text{ K}$  as shown in Fig. 1 (ref. 16). Our result supports a negatively sloped LDM much further in the metastable region.

Tension in a single inclusion was generated by isochoric cooling (see Fig. 1 and Methods) at a constant rate, with three values of the cooling rate  $r$ . A total of 154 values of the cavitation temperature  $T_{\text{cav}}$  were measured and sorted (see Methods) to generate the survival curves  $\Sigma(T)$  (probability that nucleation did

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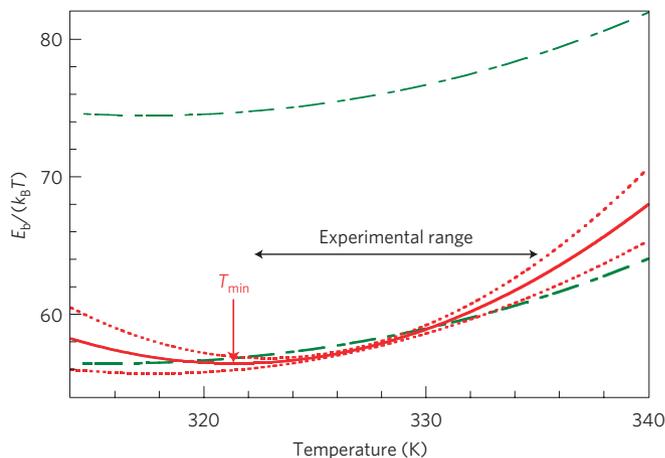
**Figure 2 | Cavitation survival probability  $\Sigma$  as a function of temperature.** Three cooling rates were investigated:  $r = 2$  (purple squares), 5 (green crosses) and  $10 \text{ K min}^{-1}$  (red circles). The data can be collapsed together (inset) by plotting the quantity  $-r \ln \Sigma / \ln 2$ . The solid black curve in the inset is a fit based on a quadratic expansion of  $E_b / (k_B T)$  (see Methods). The fit parameters are given in Supplementary Table S1. The calculated  $\Sigma$  for each  $r$  based on these values (main graph, solid curves) is in good agreement with each separate data set.

not occur during cooling down to  $T$ ) shown in Fig. 2. We find that the faster the cooling rate, the higher  $\Sigma$  at a given  $T$ . This feature and the shape of the survival curves can be explained as follows. The probability of nucleation is described by a thermally activated process that occurs at a rate  $\Gamma = \Gamma_0 \exp(-E_b / (k_B T))$ , where  $\Gamma_0$  is a kinetic prefactor (see Methods) and  $E_b$  is the free-energy barrier that has to be overcome. Integration from the temperature  $T_h$  at which the cooling starts (Fig. 1, point b) yields (see Supplementary Information):

$$\ln \Sigma = -\frac{1}{r} \int_T^{T_h} \Gamma_0(T') V \exp\left(-\frac{E_b(T')}{k_B T'}\right) dT' \quad (1)$$

As  $E_b / (k_B T)$  appears in an exponential, we can safely replace  $T_h$  by infinity, and  $\Gamma_0(T')$  by its value at a reference temperature. Equation (1) explains why  $\Sigma$  increases with  $r$ . Moreover, it predicts a data collapse when plotting  $-r \ln \Sigma$  versus  $T$ : this is borne out by the experiment (Fig. 2, inset). The small data spread is due to statistical fluctuations, as we have checked by simulations (see Supplementary Information). This proves that all data are consistently described by the same function for  $E_b$  over more than three decades of nucleation rates. For a quantitative analysis, we use a second-order temperature expansion of  $E_b / (k_B T)$  (equation (2), Methods). A first-order expansion would give a straight line in Fig. 2, inset. As the data exhibit a curvature, the experiment is able to capture the second-order term in the expansion. Including this term and integrating equation (1) yields an equation for  $-r \ln \Sigma$  involving the error function and three adjustable parameters (equation (3), Methods), which was used to fit the collapsed data for  $\ln(-r \ln \Sigma)$  simultaneously (Fig. 2, inset). The main plot in Fig. 2 also shows that the global fit parameters give  $\Sigma$  in good agreement with each data set taken separately. Uncertainties on the parameters were calculated using simulations of the whole experimental procedure (see Supplementary Information).

This statistical measurement contains a wealth of physical information. Figure 3 compares the experimental result for  $E_b$  with CNT (ref. 9). In CNT, the free-energy barrier  $E_b$  is reached for a critical bubble that is a sphere of radius  $R_c$  with an infinitely thin boundary. Using the tabulated surface tension  $\sigma$  (ref. 17),



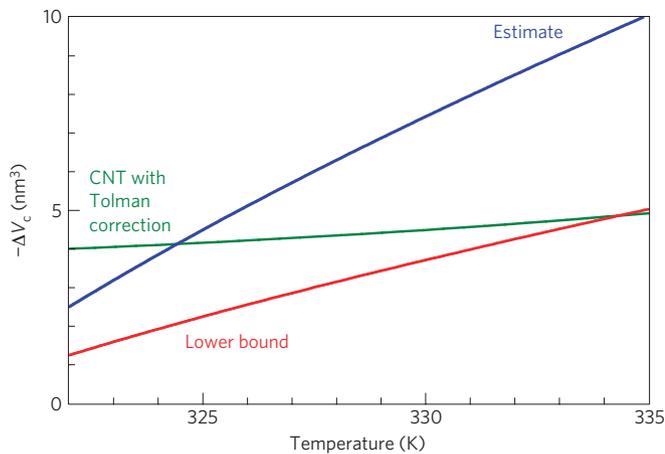
**Figure 3 | Reduced free-energy barrier  $E_b / (k_B T)$  as a function of temperature.** CNT predictions (dash-dotted green curves) without (top, thin) and with (bottom, thick) a Tolman length  $\delta = -0.047 \text{ nm}$  compared with the experimental result (equation (2), Methods, with parameters from Supplementary Table S1, solid red curve). The red dotted curves show the effect of changing all fitted parameters by plus or minus their error bars.  $E_b / (k_B T)$  reaches a minimum at  $T_{\min}$  (red arrow).

we find  $E_b$  larger than the experimental values. One possible explanation for this discrepancy involves heterogeneous nucleation on the walls of the inclusion: a finite contact angle  $\theta_c$  of the liquid/vapour interface on the substrate would reduce  $E_b$  by a geometrical factor<sup>9</sup>. However, this factor varies slowly with  $\theta_c$ , and the observed reduction would require  $\theta_c \simeq 60^\circ$ , whereas water wets silica perfectly ( $\theta_c = 0^\circ$ ). Although an impurity inducing a large  $\theta_c$  cannot be totally excluded, we note that the “dramatic character” of the cavitation process observed in a previous study<sup>18</sup> was invoked as an argument in favour of homogeneous nucleation in smooth-walled inclusions such as the one studied here. The reduction of  $E_b$  is more likely due to the simplifying approximations involved in CNT: with  $R_c$  being in the nanometre range, the surface tension for such a small bubble and curved interface might differ from the bulk one; and the liquid/vapour interface thickness, which is also around 1 nm (ref. 8), cannot be neglected. In general, these effects lead to a lower free-energy barrier. In fact, a small relative decrease in  $\sigma$ , between 7.4 and 9% across the experimental temperature range, is sufficient to bring CNT in agreement with the experiment. Such a modification is usually ascribed to the Tolman length  $\delta$  (ref. 19), which is a microscopic length that accounts for the change in surface tension with the radius of curvature  $R$  of the interface:

$$\sigma(R) = \frac{\sigma_{\text{bulk}}}{1 - \frac{2\delta}{R}}$$

For simplicity we assume a constant  $\delta$ . Remarkably, this is sufficient to obtain a good fit (Fig. 3), choosing  $\delta = -0.047 \text{ nm}$  (see Supplementary Information). This is plausible, being negative and less than the size of a water molecule (0.3 nm), similar to results obtained from simulations of the Lennard-Jones fluid<sup>20</sup>. An alternative view regarding the origin of the correction to CNT is given by density functional theory. Instead of a sharp liquid/vapour interface, it considers smoothly varying profiles and has the spinodal curve built in. For cavitation in water, this gives a correction at 330 K equivalent to a decrease in surface tension by 13 and 23.8% for two choices for the equation of state<sup>8</sup> (EOS), close to the present value.

From the curvature in  $E_b / (k_B T)$ , we deduce that it reaches a minimum at  $T_{\min} = 321.4 \pm 4.3 \text{ K}$  (red arrow in Fig. 3), very close



**Figure 4 | Volume of the critical bubble  $-\Delta V_c$  as a function of temperature.** The temperature range is that covered by the experiments. The CNT prediction with a Tolman length  $\delta = -0.047$  nm (green) is compared with the experimental lower bound (red) and estimate (blue).

to the minimum  $T_{\text{cav}}$  in the experiment, 322.35 K, which gives us confidence in the short extrapolation involved. At  $T_{\text{min}}$ , the nucleation rate reaches a maximum, and so does the metastability. This is related to the shape of the isochore (Fig. 1), as already proposed in ref. 3: when it intersects the LDM of water at  $T_{\text{LDM}}$ , the system reaches a maximum tension. The extrapolation of the EOS measured at positive pressure<sup>21</sup> gives  $T_{\text{LDM}} = 296.4$  K. With this EOS and after accounting for the temperature variation of  $\sigma^3/T$  (see Supplementary Information), CNT predicts  $T_{\text{min}} = 317.6$  K. The agreement with the experimental  $T_{\text{min}}$  is remarkable, because the analysis of the experimental data to find  $T_{\text{min}}$  does not rely on any specific EOS. We thus find support for the extrapolation of the EOS measured at positive pressure<sup>21</sup>, which gives a negative slope in the  $P$ - $T$  plane for the LDM, far inside the negative pressure region. We emphasize that the previous measurement of the LDM (ref. 16) reached a maximum tension of  $-20.3$  MPa. Even if we add only one point, it lies at a much larger tension, around  $-120$  MPa, a region highly relevant in the debate about the scenarios proposed to explain water anomalies. Our result stands in contrast with some of the simulation results: for instance, for the ST2 potential<sup>15</sup>, the LDM reaches a maximum temperature at  $P = -25$  MPa, and exhibits a positive slope at larger tension. We conclude that either the LDM is monotonic and the spinodal is re-entrant, or the turning point of the LDM lies at a density lower than  $922.8$  kg  $\text{m}^{-3}$ .

The results for  $E_b/(k_B T)$  can also be analysed using the nucleation theorem<sup>13</sup>. Without assuming any specific microscopic model for nucleation, it relates the properties of the critical nucleus to the variation of the free-energy barrier. In the present case of an isochoric experiment, it directly provides a lower bound on  $-\Delta V_c$  (see Supplementary Information), where  $\Delta V_c < 0$  is the excess volume of the critical nucleus. Using an analogy with CNT, we estimate that  $-\Delta V_c$  is around twice this lower bound (see Supplementary Information). The results are shown in Fig. 4. The small difference between the data and the corrected CNT for  $E_b/(k_B T)$  (Fig. 3) is amplified by taking the derivative, but the agreement remains satisfactory. For comparison, in the acoustic cavitation experiment<sup>11</sup> at the same temperature, the cavitation pressure was around five times less negative but  $-\Delta V_c$  was only slightly higher, around  $10$  nm<sup>3</sup>: this suggests that acoustic cavitation occurs heterogeneously on a ubiquitous impurity<sup>11</sup>, which helps to keep a small critical nucleus with a free-energy barrier that can be overcome by thermal fluctuations. The nature of this ubiquitous impurity remains an outstanding question and calls for more experiments and simulations.

## Methods

The quartz sample was synthesized in a previous study<sup>22</sup> following the thermal crack sealing method<sup>23</sup> using an internally heated apparatus at 750 MPa and 530 °C for 13 days. The inclusion studied was chosen for its extreme metastability, around  $-120$  MPa, close to the largest tension obtained previously with a higher density inclusion<sup>3</sup>. It is a 20- $\mu\text{m}$ -long and 4- $\mu\text{m}$ -wide tube, with a volume of 570  $\mu\text{m}^3$  (Fig. 1). Among the samples available to us<sup>22</sup>, those with lower densities nucleate at higher temperatures and lower tensions, and samples with higher densities nucleate very rarely, which precludes the statistical approach. A previous statistical study on the lifetime of the metastable state in one inclusion has been reported<sup>24</sup>, but the largest tension reached was  $-16.7$  MPa and only a limited analysis was performed.

The quartz fragment was placed on a heating-cooling stage (Linkam THMS 600) mounted on a microscope (Olympus BHS). Phase changes in the inclusions were observed with a  $\times 50$  long-working-distance objective (Olympus). The temperature cycles of the stage are controlled using a computer program. Phase changes are recorded using a black and white camera with a complementary metal-oxide-semiconductor 2/3 inch sensor, 1280  $\times$  1024 pixels<sup>2</sup> (Marlin) at a rate of 13 fps. The generation of negative pressure can be understood from Fig. 1. We start with a bi-phasic fluid inclusion, containing liquid and vapour (a in Fig. 1). On warming along the liquid-vapour equilibrium (a-b), the bubble disappears at the homogenization temperature  $T_h$  (b), from which the liquid density  $\rho$  is obtained using the known EOS (ref. 21). For the inclusion studied,  $T_h = 143.7$  °C and  $\rho = 922.8$  kg  $\text{m}^{-3}$ . Then, on isochoric cooling, the bubble does not reappear and the liquid is stretched (b-c), until cavitation occurs at  $T_{\text{cav}}$  (c), bringing the system back to equilibrium (a).

We have studied a single inclusion with ramp-like cooling cycles: the temperature decreases linearly with time. We used three cooling rates,  $r = 2, 5$  and  $10$  K  $\text{min}^{-1}$ , measuring 61, 33 and 60 values of  $T_{\text{cav}}$ , respectively. For a given  $r$ , the  $N_r$  values of  $T_{\text{cav}}$  were sorted in ascending order to give a list  $(T_i)_{1 \leq i \leq N_r}$ . The survival probability was deduced as  $\Sigma(T_i) = (i - 1/2)/N_r$ .

The nucleation rate is  $\Gamma = \Gamma_0 \exp[-E_b/(k_B T)]$ , where  $\Gamma_0$  is a prefactor and  $E_b$  is the free-energy barrier that has to be overcome. We take  $\Gamma_0 = N_A \rho \sqrt{2\sigma/(\pi m)}$  (refs 9,25), where  $N_A$  is Avogadro's constant,  $m$  is the mass of one molecule and  $\sigma$  is the surface tension of water. There is some uncertainty in the value of the prefactor  $\Gamma_0$ , but, because of the exponential in the rate, changes by several orders of magnitude affect the results only marginally. In the present experiment, the rate  $\Gamma$  ranges from  $2.5 \times 10^{11}$  to  $3 \times 10^{14}$   $\text{m}^{-3} \text{s}^{-1}$ .

The free-energy barrier to nucleation is calculated from equation (1) (see Supplementary Information for details). We use a second-order temperature expansion:

$$\frac{E_b(T)}{k_B T} = \frac{E_b(T^*)}{k_B T^*} + \xi \left( \frac{T}{T^*} - 1 \right) + \frac{1}{2} \kappa \left( \frac{T}{T^*} - 1 \right)^2 \quad (2)$$

where  $\xi$  and  $\kappa$  are non-dimensional parameters, and  $T^* = 328$  K is a reference temperature in the middle of the experimental range of  $T_{\text{cav}}$ . The choice of  $T^*$  does not affect the conclusions of the study. Equation (1) becomes:

$$-r \ln \Sigma = \Gamma_0(T^*) V \exp\left(-\frac{E_b(T^*)}{k_B T^*}\right) T^* \times \sqrt{\frac{\pi}{2\kappa}} e^{\frac{\xi^2}{2\kappa}} \left\{ 1 - \text{erf}\left[\frac{\xi}{\sqrt{2\kappa}} + \sqrt{\frac{\kappa}{2}} \left(\frac{T}{T^*} - 1\right)\right] \right\} \quad (3)$$

where erf stands for the error function. Equation (3) was used to fit the data for  $\ln(-r \ln \Sigma)$  (Fig. 2, inset).

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### Author contributions

F.C. designed the study, M.E.M.A. selected the inclusion, M.E.M.A. and C.R. performed the experiments, F.C. analysed the data with contributions from J-F.L. and M.E.M.A., and F.C. wrote the paper.

### Additional information

Supplementary information is available in the online version of the paper. Reprints and permissions information is available online at [www.nature.com/reprints](http://www.nature.com/reprints). Correspondence and requests for materials should be addressed to F.C.

### Competing financial interests

The authors declare no competing financial interests.